# Localization of the continuous Anderson Hamiltonian in 1-d and its transition towards delocalization 

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Works in collaboration with Cyril Labbé

## SCHRÖDINGER OPERATORS

## Schrödinger operator in 1-d

For $u:[0, L] \rightarrow \mathbb{R}$

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u \mapsto-u^{\prime \prime}+V \cdot u
$$

$V:[0, L] \rightarrow \mathbb{R}$ : potential, self-adjoint operator with Dirichlet boundary conditions.
Models disordered solids in physics where disorder $=V$.

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Discrete analog: tridiagonal matrix

$$
\left(\begin{array}{cccc}
V_{1} & 1 & & \\
1 & V_{2} & \ddots & \\
& \ddots & \ddots & 1 \\
& & 1 & V_{N}
\end{array}\right)
$$

## Case $V=0$ : Laplacian on $[0, L]$ with Dirichlet b.c.

$$
\begin{aligned}
& -u^{\prime \prime}(x)=\lambda u(x) \\
& u(0)=u(L)=0
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Eigenvalues $\lambda_{1}<\lambda_{2}<\cdots$ satisfy:

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\lambda_{k}=(\pi k / L)^{2} .
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And the associated eigenvectors are:

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x \in[0, L] \mapsto \sin \left(\frac{\pi k}{L} x\right) .
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Eigenvectors are completely delocalized!

## Continuous Anderson Hamiltonian in 1-d

We choose $V=\xi$ : white noise. For $u:[0, L] \rightarrow \mathbb{R}$,

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Be careful: Multiplication by the white noise does not make sense!
Fukushima, Nakao ('77) proved:

- Well-defined self-adjoint operator,
- discrete simple spectrum bounded from below: $\lambda_{1}<\lambda_{2}<\cdots$,
- associated eigenvectors $\left(\varphi_{k}\right)_{k}$ form an orthonormal basis of $L^{2}([0, L])$ and are $C^{3 / 2-}$.


## Goal

Study the spectrum of this operator when $L \rightarrow \infty$.


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- Localization of the eigenvectors and Poisson distribution of eigenvalues.

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- Delocalization of the eigenvectors and repulsion of the eigenvalues.

Previous results on $\mathcal{H}_{L}$

## Density of states

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For the Laplacian, its eigenvalues are:

$$
\lambda_{k}=(\pi k / L)^{2}
$$

$\rightarrow$ Density of states:

$$
E \in \mathbb{R}_{+} \mapsto \frac{1}{2 \pi \sqrt{E}}
$$

## Density of states for $\mathcal{H}_{L}$

Frisch and Lloyd ('60), Halperin ('65) and then Fukushima, Nakao ('77): Explicit integral formula for the density of states of $\mathcal{H}_{L}$ :

$$
n(E)=\frac{d}{d E}\left(\sqrt{2 \pi} \int_{0}^{\infty} u^{-1 / 2} e^{-\frac{1}{6} u^{3}-2 E u} d u\right)^{-1}
$$



## First eigenvalue

McKean ('94): Convergence of the smallest eigenvalue $\lambda_{1}$ (recentred and rescaled) for Dirichlet, Neumann and periodic b.c.:

$$
-4 \sqrt{a_{L}}\left(\lambda_{1}+a_{L}\right) \Rightarrow_{L \rightarrow \infty} e^{-e^{-x}} d x
$$

where $a_{L} \sim\left(\frac{3}{8} \ln L\right)^{2 / 3}$

## Our results on $\mathcal{H}_{L}$

## Localization of the smallest eigenvectors

Recall $a_{L} \sim\left(\frac{3}{8} \ln L\right)^{2 / 3}$. Denote by

$$
\mathcal{Q}_{L}:=\sum_{k \geq 1} \delta_{4 \sqrt{a_{L}}\left(\lambda_{k}+a_{L}\right)}
$$

Spacing:
$O\left(1 / \sqrt{a_{L}}\right)$

$$
\begin{aligned}
& \lambda_{1} \lambda_{2} \quad \lambda_{3} \\
& \times \times \times \times \sim \times \\
& \times a_{L} \\
& \simeq-(\ln L)^{2 / 3}
\end{aligned}
$$

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\mathcal{Q}_{L}:=\sum_{k \geq 1} \delta_{4 \sqrt{a_{L}}\left(\lambda_{k}+a_{L}\right)}, \quad m_{L}(d t):=\left(L \varphi_{k}(L t)^{2} d t\right)_{k \geq 1} .
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## Theorem (D., Labbé ('17))

$\left(\mathcal{Q}_{L}, m_{L, k}(d t)\right)$ converges in distribution towards $\left(\mathcal{Q}_{\infty}, m_{\infty}\right)$ where:

- $\mathcal{Q}_{\infty}$ : Poisson point process of intensity $e^{x} d x$,
- $m_{\infty}=\left(\delta_{u_{k}}\right)_{k \geq 1}:\left(u_{k}\right)_{k \geq 1}$ i.i.d, uniform on $[0,1]$, independent of $\mathcal{Q}_{\infty}$.



## Simulation of the first eigenvectors



The first 5 eigenvectors $\varphi_{k}^{2}$ in order: black, blue, purple, red, green ( $L=300$ ).

## Shape of the first eigenvectors



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- For all fixed $k, \varphi_{k}$ decays exponentially at rate $\sqrt{a_{L}}$.


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Theorem (D., Labbé ('17))

- For all fixed $k, \varphi_{k}$ decays exponentially at rate $\sqrt{a_{L}}$.
- Let $U_{k}$ be the point where $\varphi_{k}$ reaches its maximum.

$$
h_{k}(t):=\sqrt{a_{L}} \varphi_{k}^{2}\left(U_{k}+\sqrt{a_{L}} t\right) \rightarrow_{L \rightarrow \infty} 1 / \cosh (t)^{2}
$$

uniformly over compact subsets of $\mathbb{R}$.

## Zoom around the maximum of $\varphi_{k}^{2}$



## Schematic shape of the fifth eigenvector



Note that we know for example precisely the position of the $k-1$ zeros of $\varphi_{k}$

## Localization and transition towards delocalization



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Eigenvector located around $E$ :


## Localization for $E \ll L$

Let $E=E(L)$ be the re-centering of the eigenvalues and define $\mathcal{Q}_{E}(d x):=\sum_{i \geq 1} \delta_{\operatorname{Ln}(E)\left(\lambda_{i}-E\right)}(d x), \quad$ where $n(E)=$ density of states.

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- Crossover regime: $1 \ll E \ll L$
$\mathcal{Q}_{E}$ converges to a Poisson point process of intensity $d x$.
Eigenvectors: exponentially decreasing at speed $c(E)=O(1 / E)$. Moreover, a "typical" eigenvector chosen w.r.t. the "spectral measure" looks like the exponential of a Brownian motion plus a drift on a region of size $E$.


## Transition towards delocalization

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- If $E \sim \alpha L$ with $\alpha \in(0, \infty) \mathcal{Q}_{E}-\left(\left(2 \sqrt{\alpha} L^{3 / 2}\right) \bmod 2 \pi\right)$ converges towards a point process with repulsion between the points. It corresponds to the point process Sch $h_{1 / \alpha}$.
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Eigenvectors: exponential of a Brownian motion plus a drift (as conjectured by Rifking and Virág).
- If $E \gg L$ then $\mathcal{Q}_{E}$ converges to the deterministic process of eigenvalues of $-d^{2} / d x^{2}$.


## Limiting operator

For this slide, let us define $\mathcal{H}_{L}$ on $L^{2}[-L, L]$ instead of $L^{2}[0, L]$, so that $[-L, L]$ converges to the whole line $\mathbb{R}$. We denote $\varphi_{i}^{L}$ and $\lambda_{i}^{L}$ its eigenvectors and eigenvalues.

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Define $\mathcal{H} f:=-f^{\prime \prime}+\xi f$ on a domain
$\mathcal{D}:=\left\{f \in L^{2}(\mathbb{R}), f \mathrm{AC}, f^{\prime}-\mathrm{Bf} \mathrm{AC}, \mathcal{H} f \in L^{2}(\mathbb{R})\right\}$ of $L^{2}(\mathbb{R})$. It is a self-adjoint operator, which is limit point at both sides.

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## Theorem (D., Labbé ('20+))

The spectral measures $\sum_{i}\left(\varphi_{i}^{L}(0)^{2}+\varphi_{i}^{\prime L}(0)^{2}\right) \delta_{\lambda_{i}^{L}}$ converge a.s. (for the topology of the vague convergence) towards the spectral measure of $\mathcal{H}$.
This spectral measure is pure point: the operator $\mathcal{H}$ is a pure point operator.

## Some ideas for the proofs

## Eigenvalue equation

Eigenvalue equation for $\mathcal{H}_{L}$ defined on $[0, L]$ :

$$
-\varphi^{\prime \prime}+\xi \cdot \varphi=\lambda \varphi
$$

with $\varphi(0)=0$ (without any condition on $\varphi(L)$ ).

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For all $\lambda \in \mathbb{R}$, there is an unique solution $\varphi_{\lambda}$ (up to a scaling).
The couple $\left(\lambda, \varphi_{\lambda}\right)$ is an eigenvalue/eigenvector when

$$
\varphi_{\lambda}(L)=0
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## Eigenvalue equation

One can also impose first $\hat{\varphi}(L)=0$ and solve

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$$
\hat{\varphi}_{\lambda}(0)=0 .
$$

## Concatenation forward/backward

Key idea: Use forward solution $\varphi_{\lambda}$ on the time-interval $[0, u]$ and then backward solution $\hat{\varphi}_{\lambda}$ on $[u, L]$ for some well-chosen $u$. $\rightarrow$ Concatenation is $\varphi^{(u)}$.


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- FACT: If $\lambda$ close to an eigenvalue $\rightarrow$ close to eigenvector if $u=$ argmax of eigenvector.

It helps A LOT because it is much easier to analyze the forward or backward solution of the ODE than the eigenvalue equation (when $\lambda$ eigenvalue, $\lambda$ is random and depends on the whole potential $\xi!$ ).

## Localization of the first eigenvector

 Simulation of $\varphi_{\lambda_{1}}^{\prime} / \varphi_{\lambda_{1}}$ :

## Localization of the first eigenvector

One can approximate $\varphi_{\lambda_{1}}^{\prime} / \varphi_{\lambda_{1}}$ by $\varphi_{\lambda}^{\prime} / \varphi_{\lambda}$ on $[0, u]$ and then by $\hat{\varphi}_{\lambda}^{\prime} / \hat{\varphi}_{\lambda}$ on $[u, L]$ for $\lambda$ close to $\lambda_{1}$


$$
\frac{\varphi_{\lambda_{1}}(t)}{\varphi_{\lambda_{1}}\left(t_{0}\right)}=\exp \left(\int_{t_{0}}^{t} \frac{\varphi_{\lambda_{1}}^{\prime}(s)}{\varphi_{\lambda_{1}}(s)} d s\right)
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## Localization in the bulk: A key formula

## Proposition (Goldsheid Molchanov Pastur formula)

For all continuous and bounded $G$ :

$$
\mathbb{E}\left[\sum_{\lambda \text { eigenvalue }} G\left(\lambda, \varphi_{\lambda}\right)\right]
$$

$$
=\int_{\lambda \in \mathbb{R}} \int_{0}^{L} \int_{0}^{\pi} \sin ^{2}(\theta) p_{\lambda}(\theta) p_{\lambda}(\pi-\theta) \mathbb{E}\left[G\left(\lambda, \frac{\varphi_{\lambda}^{(u)}}{\left\|\varphi_{\lambda}^{(u)}\right\|_{2}}\right)\right] d \lambda d u d \theta
$$

where

- $\varphi^{(u)}$ is the concatenation of the forward process and backward process at time $u$.
- $p_{\lambda}(\theta)$ transition probability of $\theta_{\lambda}$ "phase function" (argument of $\varphi_{\lambda}^{\prime}+i \varphi_{\lambda}$ ).

Strategy to prove convergence towards a Poisson point process when you know localization

Let $\Delta=[E-h /(\operatorname{Ln}(E)), E+h /(\operatorname{Ln}(E))](E$ fixed $)$ and denote $N(\Delta)=\#$ eigenvalues in $\Delta$.

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Divide $[0, L]$ into small boxes $B_{i}, i=1, \cdots, N$ of same length.


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Eigenvectors

(A) $N(\Delta) \simeq \sum_{i} N_{i}(\Delta)$ where $N_{i}(\Delta)$ is the number of eigenvalues in $\Delta$ of $\mathcal{H}_{B_{i}}:=\left(-d^{2} / d x^{2}+B^{\prime}(x)\right)_{\mid B_{i}}$.

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(B) $\mathbb{E}[N(\Delta)] \sim 2 h \rightarrow$ much stronger than the density of states!
(C) $\sum_{i} \mathbb{P}\left[N_{i}(\Delta) \geq 2\right] \rightarrow 0$.

## THANK YOU!

