Localization of the continuous Anderson Hamiltonian in 1-d and its transition towards delocalization

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Works in collaboration with Cyril Labbé

## Schrödinger operators

Schrödinger operator in 1-d For  $u : [0, L] \rightarrow \mathbb{R}$ 

$$u\mapsto -u''+\mathbf{V}\cdot u$$
.

 $V : [0, L] \rightarrow \mathbb{R}$ : potential, self-adjoint operator with Dirichlet boundary conditions.

Models disordered solids in physics where disorder = V.

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and the multiplication by the potential V:

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Discrete analog: tridiagonal matrix

$$\begin{pmatrix} V_1 & 1 & & \\ 1 & V_2 & \ddots & \\ & \ddots & \ddots & 1 \\ & & 1 & V_N \end{pmatrix}$$

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And the associated eigenvectors are:

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Eigenvectors are completely delocalized!

## Continuous Anderson Hamiltonian in 1-d

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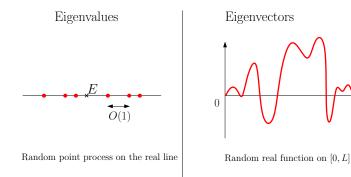
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Fukushima, Nakao ('77) proved:

- Well-defined self-adjoint operator,
- ▶ discrete simple spectrum bounded from below:  $\lambda_1 < \lambda_2 < \cdots$ ,
- ► associated eigenvectors (\(\varphi\_k\))\_k\) form an orthonormal basis of L<sup>2</sup>([0, L]) and are C<sup>3/2-</sup>.

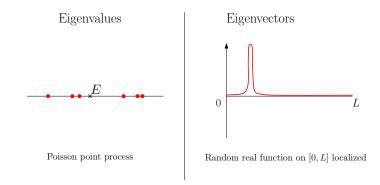
## Goal

#### Study the **spectrum** of this operator when $L \rightarrow \infty$ .



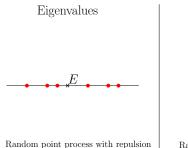
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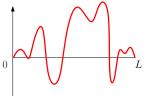


Usually for random operators, there is a dichotomy:

 Localization of the eigenvectors and Poisson distribution of eigenvalues. Study the **spectrum** of this operator when  $L \rightarrow \infty$ .



Eigenvectors



Random real function on [0, L] delocalized

Usually for random operators, there is a dichotomy:

Delocalization of the eigenvectors and repulsion of the eigenvalues.

## Previous results on $\mathcal{H}_L$

## Density of states

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For the Laplacian, its eigenvalues are:

$$\lambda_k = (\pi k/L)^2.$$

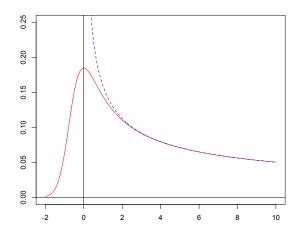
 $\rightarrow$  Density of states:

$$E \in \mathbb{R}_+ \mapsto rac{1}{2\pi\sqrt{E}}.$$

#### Density of states for $\mathcal{H}_L$

Frisch and Lloyd ('60), Halperin ('65) and then Fukushima, Nakao ('77): Explicit integral formula for the density of states of  $H_L$ :

$$n(E) = \frac{d}{dE} \left( \sqrt{2\pi} \int_0^\infty u^{-1/2} e^{-\frac{1}{6}u^3 - 2Eu} du \right)^{-1}$$

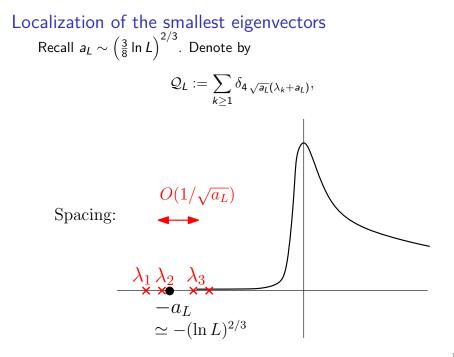


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McKean ('94) : Convergence of the smallest eigenvalue  $\lambda_1$  (recentred and rescaled) for Dirichlet, Neumann and periodic b.c.:

$$-4\sqrt{a_L}(\lambda_1 + a_L) \Rightarrow_{L \to \infty} e^{-e^{-x}} dx,$$
  
where  $a_L \sim \left(\frac{3}{8} \ln L\right)^{2/3}$ 

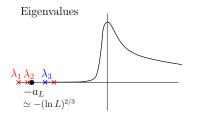
# Our results on $\mathcal{H}_L$

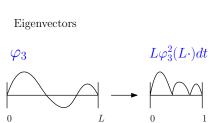


## Localization of the smallest eigenvectors

- /-

Recall 
$$a_L \sim \left(\frac{3}{8} \ln L\right)^{2/3}$$
. Denote by  
 $Q_L := \sum_{k \ge 1} \delta_4 \sqrt{a_L} (\lambda_k + a_L), \qquad m_L(dt) := (L \varphi_k (L t)^2 dt)_{k \ge 1}.$ 



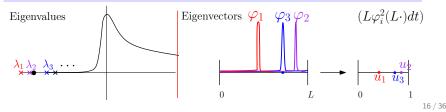


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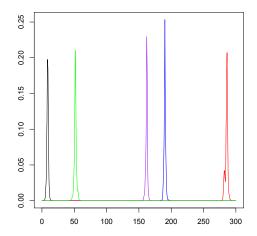
## Theorem (D., Labbé ('17))

 $(\mathcal{Q}_L, m_{L,k}(dt))$  converges in distribution towards  $(\mathcal{Q}_\infty, m_\infty)$  where:

Q<sub>∞</sub>: Poisson point process of intensity e<sup>x</sup> dx,
 m<sub>∞</sub> = (δ<sub>uk</sub>)<sub>k≥1</sub> : (u<sub>k</sub>)<sub>k≥1</sub> i.i.d, uniform on [0, 1], independent of Q<sub>∞</sub>.

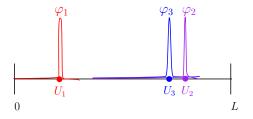


## Simulation of the first eigenvectors



The first 5 eigenvectors  $\varphi_k^2$  in order: black, blue, purple, red, green (L = 300).

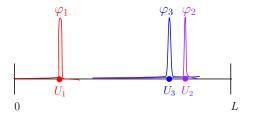
## Shape of the first eigenvectors



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For all fixed k,  $\varphi_k$  decays exponentially at rate  $\sqrt{a_L}$ .

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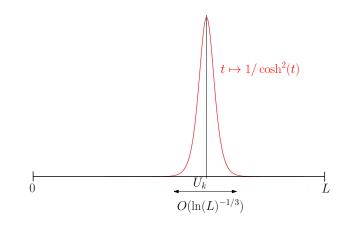
#### Theorem (D., Labbé ('17))

- For all fixed k,  $\varphi_k$  decays exponentially at rate  $\sqrt{a_L}$ .
- Let  $U_k$  be the point where  $\varphi_k$  reaches its maximum.

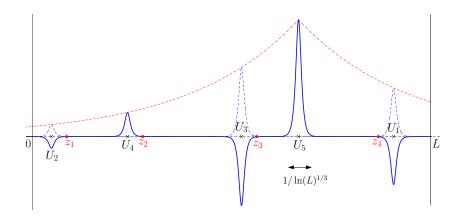
$$h_k(t) := \sqrt{a_L} \ arphi_k^2 (U_k + \sqrt{a_L} \ t) 
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uniformly over compact subsets of  $\mathbb{R}$ .

## Zoom around the maximum of $\varphi_k^2$

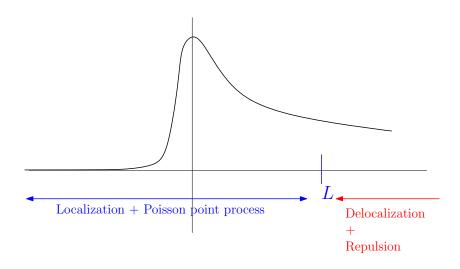


## Schematic shape of the fifth eigenvector

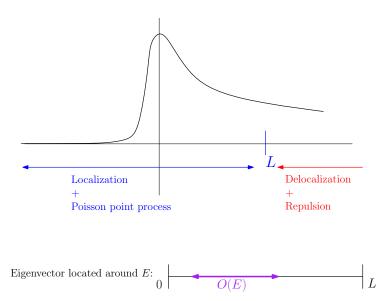


Note that we know for example precisely the position of the k-1 zeros of  $\varphi_k$ 

## Localization and transition towards delocalization



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Size of localization

Localization for  $E \ll L$ 

Let E = E(L) be the re-centering of the eigenvalues and define

 $\mathcal{Q}_E(dx) := \sum_{i \ge 1} \delta_{Ln(E)(\lambda_i - E)}(dx) \;, \quad \text{where } n(E) = \text{density of states}.$ 

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- Bulk regime: E fixed (independent of L)
   Q<sub>E</sub> converges to a Poisson point process of intensity dx.
   Eigenvectors: exponentially decreasing at constant speed c(E) (not depending on L).
- ► Crossover regime: 1 ≪ E ≪ L Q<sub>E</sub> converges to a Poisson point process of intensity dx. Eigenvectors: exponentially decreasing at speed c(E) = O(1/E). Moreover, a "typical" eigenvector chosen w.r.t. the "spectral measure" looks like the exponential of a Brownian motion plus a drift on a region of size E.

## Transition towards delocalization

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If E ~ αL with α ∈ (0,∞) Q<sub>E</sub> - ((2√αL<sup>3/2</sup>) mod 2π) converges towards a point process with repulsion between the points. It corresponds to the point process Sch<sub>1/α</sub>. Eigenvectors: exponential of a Brownian motion plus a drift (as conjectured by Rifking and Virág).

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- If E ≫ L then Q<sub>E</sub> converges to the deterministic process of eigenvalues of −d<sup>2</sup>/dx<sup>2</sup>.

## Limiting operator

For this slide, let us define  $\mathcal{H}_L$  on  $L^2[-L, L]$  instead of  $L^2[0, L]$ , so that [-L, L] converges to the whole line  $\mathbb{R}$ . We denote  $\varphi_i^L$  and  $\lambda_i^L$  its eigenvectors and eigenvalues.

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Define  $\mathcal{H}f := -f'' + \xi f$  on a domain  $\mathcal{D} := \{f \in L^2(\mathbb{R}), f \text{ AC}, f' - Bf \text{ AC}, \mathcal{H}f \in L^2(\mathbb{R})\} \text{ of } L^2(\mathbb{R}).$  It is a self-adjoint operator, which is limit point at both sides.

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#### Theorem (D., Labbé ('20+))

The spectral measures  $\sum_{i} (\varphi_{i}^{L}(0)^{2} + \varphi_{i}^{\prime L}(0)^{2}) \delta_{\lambda_{i}^{L}}$  converge a.s. (for the topology of the vague convergence) towards the spectral measure of  $\mathcal{H}$ . This spectral measure is pure point: the operator  $\mathcal{H}$  is a **pure point operator**.

## Some ideas for the proofs

Eigenvalue equation for  $\mathcal{H}_L$  defined on [0, L]:

$$-\varphi'' + \mathbf{\xi} \cdot \varphi = \lambda \varphi$$

with  $\varphi(0) = 0$  (without any condition on  $\varphi(L)$ ).

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 $\varphi_{\lambda}(L)=0.$ 

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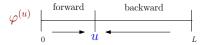
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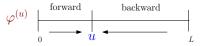
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**Key idea:** Use forward solution  $\varphi_{\lambda}$  on the time-interval [0, u] and then backward solution  $\hat{\varphi}_{\lambda}$  on [u, L] for some well-chosen u.  $\rightarrow$  Concatenation is  $\varphi^{(u)}$ .

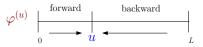


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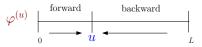
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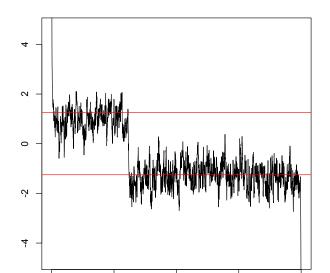


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It helps A LOT because it is **much easier** to analyze the forward or backward solution of the ODE than the eigenvalue equation (when  $\lambda$  eigenvalue,  $\lambda$  is random and depends on the whole potential  $\xi$ !).

## Localization of the first eigenvector

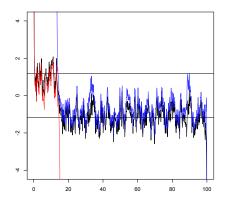
Simulation of  $\varphi'_{\lambda_1}/\varphi_{\lambda_1}$ :



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#### Localization of the first eigenvector

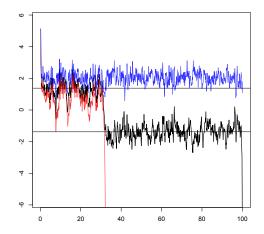
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$$\frac{\varphi_{\lambda_1}(t)}{\varphi_{\lambda_1}(t_0)} = \exp\Big(\int_{t_0}^t \frac{\varphi_{\lambda_1}'(s)}{\varphi_{\lambda_1}(s)} ds\Big)$$

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### Localization in the bulk: A key formula

#### Proposition (Goldsheid Molchanov Pastur formula)

For all continuous and bounded G:

$$\begin{split} & \mathbb{E}\Big[\sum_{\lambda \text{ eigenvalue}} G(\lambda,\varphi_{\lambda})\Big] \\ &= \int_{\lambda \in \mathbb{R}} \int_{0}^{L} \int_{0}^{\pi} \sin^{2}(\theta) p_{\lambda}(\theta) p_{\lambda}(\pi-\theta) \mathbb{E}\Big[G\Big(\lambda,\frac{\varphi_{\lambda}^{(u)}}{||\varphi_{\lambda}^{(u)}||_{2}}\Big)\Big] d\lambda du d\theta, \end{split}$$

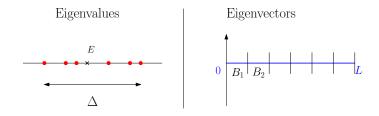
where

- φ<sup>(u)</sup> is the concatenation of the forward process and backward process at time u.
- ▶  $p_{\lambda}(\theta)$  transition probability of  $\theta_{\lambda}$  "phase function" (argument of  $\varphi'_{\lambda} + i\varphi_{\lambda}$ ).

Let  $\Delta = [E - h/(Ln(E)), E + h/(Ln(E))]$  (*E* fixed) and denote  $N(\Delta) = \#$  eigenvalues in  $\Delta$ .

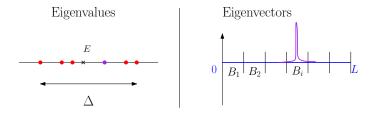
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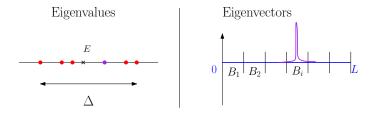
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(A)  $N(\Delta) \simeq \sum_i N_i(\Delta)$  where  $N_i(\Delta)$  is the number of eigenvalues in  $\Delta$  of  $\mathcal{H}_{B_i} := (-d^2/dx^2 + B'(x))_{|B_i}$ .

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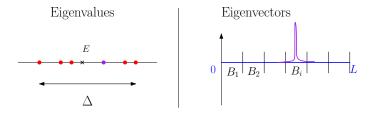


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(B)  $\mathbb{E}[N(\Delta)] \sim 2h$ 

Let  $\Delta = [E - h/(Ln(E)), E + h/(Ln(E))]$  (*E* fixed) and denote  $N(\Delta) = \#$  eigenvalues in  $\Delta$ .

Divide [0, L] into small boxes  $B_i$ ,  $i = 1, \dots N$  of same length.

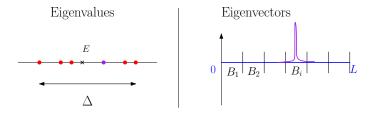


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(B)  $\mathbb{E}[N(\Delta)] \sim 2h \rightarrow$  much stronger than the density of states! (C)  $\sum_{i} \mathbb{P}[N_i(\Delta) \ge 2] \rightarrow 0.$ 

# THANK YOU!