# Quantitative inequalities and convergence of thresholding schemes in optimal control theory

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# Objective of the talk

- Develop quantitative inequalities for spectral optimisation problems or more general optimal control problems.
- Apply them to the study of numerical schemes.
- Main reference: Chambolle, M., Privat, Math. Ann., 2025.

## Reference spectral optimisation problem

 $\Omega$ : bounded, smooth domain.  $V \in L^{\infty}(\Omega)$ .

$$\lambda(V) := \min_{u \in W_0^{1,2}(\Omega), \int_{\Omega} u^2 = 1} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} V u^2 \leadsto \begin{cases} -\Delta u = \lambda u + V u & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega. \end{cases}$$

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$$\min_{0\leq V\leq 1 \text{ a.e., } \int_{\Omega}V=V_0}\lambda(V).$$

Old and well understood problem:

- Composite membrane problem (Cox, Lipton, MacLaughlin...)
- Applications to mathematical biology (Cantrell, Cosner, Berestycki, Hamel, Roques, Kao, Lou, Yanagida...)

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## What do we want to do?

$$egin{aligned} \min_{0 \leq V \leq 1 \text{ a.e., } \int_{\Omega} V = V_0} \lambda(V). \ \mathscr{V}^* := & & rgmin \ 0 \leq V \leq 1 \text{ a.e., } \int_{\Omega} V = V_0 \end{aligned}$$

Fix some  $V^* \in \mathscr{V}^*$ .

$$\lambda(V) - \lambda(V^*) \ge C \operatorname{dist}(V, \mathscr{V}^*)^{\alpha}$$

for:

- Some distance,
- **2** Some exponent  $\alpha$ ,
- And if both could be optimal...

## What do we want to do?

$$\begin{split} \underset{0 \leq V \leq 1}{\text{min}} & \underset{\text{a.e., } \int_{\Omega} V = V_0}{\text{min}} \lambda(V). \\ \mathscr{V}^* := & \underset{0 \leq V \leq 1}{\operatorname{argmin}} & \lambda(V) \end{split}$$

Fix some  $V^* \in \mathscr{V}^*$ .

$$\lambda(V) - \lambda(V^*) \ge C \operatorname{dist}(V, \mathscr{V}^*)^{\alpha}$$

for:

- Some distance,
- **2** Some exponent  $\alpha$ ,
- 4 And if both could be optimal...

Spoiler:

$$\lambda(V) - \lambda(V^*) \ge C \operatorname{dist}_{L^1}(V, \mathcal{V}^*)^2.$$

# Why do we want to do it?

#### Several applications:

- Application to numerical schemes: allows to study the ubiquitous thresholding scheme for the numerical optimisation of potentials, see later. [Chambolle, M., Privat, Math. Ann., 2025].
- Allows to simplify qualitative questions in optimal control: turnpike in bilinear control [M., Ruiz-Balet, SIMA, 2021], optimal placement of captors [M., Privat, Trélat, AIHP-C, 2025].
- In general: allows to analyse perturbation problems.

## Plan of the talk

- Basic facts about the underlying optimisation problem.
- A (short) review of the bibliography & a discussion of the coercivity norm.
- A discussion of a part of the proof.
- (Briefly) Application to numerical schemes.

# Basic facts I: bang-bang property

#### Recall:

$$\min_{0\leq V\leq 1 \text{ a.e., } \int_{\Omega}V=V_0}\left(\lambda(V):=\min_{u\in W_0^{1,2}(\Omega),\int_{\Omega}u^2=1}\int_{\Omega}|\nabla u|^2-\int_{\Omega}Vu^2\right).$$

- **1** Inf of linear functionals:  $\lambda$  is (strictly) concave.
- ② Thus: any optimal  $V^*$  is an extreme point of  $\{0 \le V \le 1, \int_{\Omega} V = V_0\}$ .
- **3** So that: any optimal  $V^*$  is a characteristic function:  $V^* = \mathbb{1}_{E^*}$ .

# Basic facts II: Free boundary problem

Double minimisation procedure:

$$\min_{0 \le V \le 1, \int_{\Omega} V = V_0} \lambda(V) = \min_{u \in W_0^{1,2}(\Omega), \int_{\Omega} u^2 = 1} \min_{0 \le V \le 1, \int_{\Omega} V = V_0} \int_{\Omega} |\nabla u|^2 - \int_{\Omega} V u^2.$$

**1** Bathtub principle (simplified and slightly wrong):  $f: \Omega \to \mathbb{R}$ .

$$\max_{0 \le V \le 1, \int_{\Omega} V = V_0} \int_{\Omega} fV = \int_{\Omega} f \mathbf{1}_{\{f > c\}}$$

where c is chosen so that

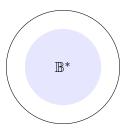
$$|\{f>c\}|=V_0.$$

Thus: if V\* is optimal,

$$V^* = \mathbb{1}_{\{u_{V^*} > c^*\}} \leadsto -\Delta u_{V^*} = \lambda (V^*) u_{V^*} + u_{V^*} \mathbb{1}_{\{u_{V^*} > c^*\}}.$$

# Basic fact III: rearrangement inequalities

When  $\Omega$  is a ball the situation is remarkably simple:  $V^* = \mathbb{1}_{\mathbb{B}^*}$ .



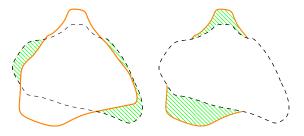
In general, the symmetries of the domain are not preserved but, if  $\Omega$  is convex, so is  $E^*$  etc. We refer to [Imai, Grieser, Kurata].

### Illustration

Assume: unique minimiser  $V^* = \mathbb{1}_{E^*}$ . We claimed:

$$\lambda(V) - \lambda(V^*) \geq C \|V - V^*\|_{L^1(\Omega)}^2.$$

For instance, if  $V = \mathbb{1}_E$ ,



In orange,  $E^*$ , in dashed, E.

$$\lambda(\mathbb{1}_E) - \lambda(\mathbb{1}_{E^*}) \ge C ^2$$
 (1)

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## Quantitative inequalities in "true" shape optimisation

Shape optimisation problem with a certain constraint:

$$\inf_{\Omega, |\Omega| = V_0 \text{ or } Per(\Omega) = P_0} \mathcal{J}(\Omega) = \mathcal{J}(\Omega^*). \tag{2}$$

A quantitative inequality writes

$$\mathcal{J}(\Omega) - \mathcal{J}(\Omega^*) \ge C \inf_{\Omega^* \text{ optimal}} \text{Vol}(\Omega \Delta \Omega^*)^2,$$
 (3)

where  $\Delta$  stands for the symmetric difference of sets. If  $\Omega^* = \mathbb{B}$ : Fraenkel asymmetry.

## Some examples

- Isoperimetric inequality:  $\mathcal{J} = \text{Per}$ , with a volume constraint, and balls are the only solutions.
  - Fusco, Maggi, Pratelli, Annals of Mathematics, 2008,
  - Figalli, Maggi, Pratelli, Inventiones Mathematicae, 2010.
- Faber-Krahn inequality:  $\mathcal{J} = \lambda_D$  (first eigenvalue of the Dirichlet laplacian) with a volume constraint, and balls are the only solutions.
  - Brasco, De Philippis, Velichkov, Duke Mathematical Journal, 2015,
  - Marpukhin, Nahon, Polterovich, Stern. 2021.

Thresholding schemes

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### Two related contributions

To come back to our problem:

$$\min_{0\leq V\leq 1,\int_{\Omega}V=V_0}\lambda(V).$$

Related contributions for the optimisation of the Dirichlet energy/of eigenvalues w.r.t a potential with  $L^p$  ( $p < \infty$ ) constraints.

- 1 Brasco, Buttazzo, Calculus of Variations and Partial Differential Equations, 2014,
- 2 Carlen, Frank, Lieb, Geometric and Functional Analysis, 2014.

The norm is different and reflects the different types of constraints. The methods break down for  $L^{\infty}$  constraints.

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## Towards the optimal norm

Why choose the  $L^1$  norm squared?

Derivative of the eigenvalue:

$$\lim_{\varepsilon \to 0} \frac{\lambda(V + \varepsilon h) - \lambda(V)}{\varepsilon} = \lambda'(V)[h] = -\int_{\Omega} h u_V^2.$$

At an optimal potential, with  $h = V' - V^*$ , the concavity of  $\lambda$  implies

$$\lambda(V) - \lambda(V^*) \ge \lambda'(V^*)[h] = -\int_{\Omega} (V - V^*)u_V^2.$$

This is another way to get  $V^* = \mathbb{1}_{\{u_{V^*} > c^*\}}$ .

**Natural question:** Can we quantify the bathtub principle and obtain something like

$$\int_{\Omega} V u_{V^*}^2 \leq \int_{\Omega} V^* u_{V^*}^2 - C \mathrm{dist}(V,V^*)^{\alpha}?$$

# The quantitative bathtub principle

Replace  $u_{V^*}^2$  by  $\varphi$ .

$$-\int_{\Omega}\varphi\mathbb{1}_{\{\varphi>c\}}+\int_{\Omega}V\varphi.$$

**1** If  $|\{\varphi = c\}| = 0$  then

$$V^* = \mathbb{1}_{\{\varphi > c\}}$$
 is the unique solution of  $\max_{0 \leq V \leq 1, \int_{\Omega} V = V_0} T(V) = \int \varphi V$ .

What we have is a quantitative inequality for a linear problem:

$$T(V) - T(V^*) \le -C||V - V^*||_{L^1}^2.$$

To obtain it: we consider

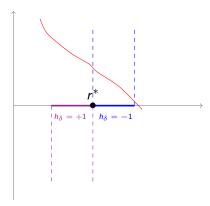
$$\max_{V,\|V-V^*\|_{I^1}=\delta} T(V) = T(V^*_\delta).$$

# The quantitative bathtub principle

If we consider the one dimensional case (or radially symmetric) we have

$$\mathcal{T}(V^*) - \mathcal{T}(V^*_\delta) = \int h_\delta arphi$$

with



# The quantitative bathtub principle

And we can compute explicitly

$$T(V_{\delta}^{*}) - T(V^{*}) = \int_{\mathbb{B}} h_{\delta} \varphi$$

$$= -\int_{r_{\delta}^{-}}^{r^{*}} \varphi \, d\mathbf{r} + \int_{r^{*}}^{r_{\delta}^{+}} \varphi \, d\mathbf{r}$$

$$\approx \varphi'(r^{*}) \left( \int_{r_{\delta}^{-}}^{r^{*}} |r - r^{*}| \, d\mathbf{r} + \int_{r^{*}}^{r_{\delta}^{+}} |r - r^{*}| \, d\mathbf{r} \right)$$

$$\sim C \varphi'(r^{*}) \delta^{2}.$$

This gives the expected norm.

#### Some references

The quantitative bathtub principle is found under a variety of guises:

- Cianchi, Ferone, quantitative Hardy-Littlewood inequality,
- Lemou, for stability issues in mathematical physics,
- Wachsmuth, for the numerical analysis of optimal control problems,
- and probably many other places.

## Some results

## Theorem (M., JDE, 2020, Chambolle, M., Privat, Math. Ann., 2025)

Let  $\mathcal{V}^*$  be the set of minimisers of  $\lambda$ .

- **1** When  $\Omega$  is a ball
- ② Or when  $\Omega$  is smooth and the volume constraint  $V_0$  is big enough there exists C>0 such that

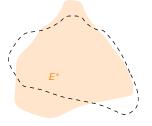
$$\lambda(V) - \lambda(V^*) \geq C \mathrm{dist}_{L^1(\Omega)}(V, \mathcal{V}^*)^2.$$

# Local stability of optimal shapes

Goal:

## Local stability of critical or optimal shapes

• We fix a critical/optimal set  $E^*$  (i.e.  $V^* = \mathbb{1}_{E^*}$  is optimal or critical in the sense that  $E^* = \{u_{V^*} > c^*\}$ ).



Dashed:  $(\operatorname{Id} + \Phi)(E^*) = E_{\Phi}^*$ . We prove that if  $\|\Phi\|_{W^{2,p}}$  is small enough (independently of  $E^*$ ) then

$$\lambda(\mathbb{1}_{E_{\Phi}^*}) - \lambda(\mathbb{1}_{E^*}) \ge C |E_{\Phi}^* \triangle E^*|^2 \tag{4}$$

If this inequality is satisfied critical/optimal sets are isolated.

## How to show that?

- First and second-order shape derivatives  $\leadsto$  assumes regularity of  $E^*$   $\leadsto$  uses the free boundary theory as developed by [Chanillo, Kenig & To, JEMS, 2008]. First place where  $V_0$  large enough is used.
- ② In [M., JDE, 2020] in the case of the ball, explicit computations.
- $\odot$  In [Chambolle, M., Privat, Math. Ann., 2025]: indirect (non computational) argument to obtain a good enough control of the second order shape derivative.  $V_0$  large enough also used.

# Strategy of proof

We fix a minimiser  $V^*$  and V such that

$$||V-V^*||_{L^1(\Omega)}=\delta\ll 1.$$

To further simplify we assume that

$$V_{\delta} = \mathbb{1}_{E_{\delta}}$$

and we let  $u_\delta$  be the associated eigenfunction. To improve the eigenvalue:

$$V_\delta'$$
 solution of  $\sup_{0 \leq W \leq 1, \int_\Omega W = V_0} \int_\Omega W u_\delta^2 \geq \int_\Omega V_\delta u_\delta^2.$ 

Indeed

$$\lambda(V) = \inf_{u \in W_0^{1,2}, \int_{\Omega} u^2 = 1} \left( \int_{\Omega} |\nabla u|^2 - \int_{\Omega} V_{\delta} u^2 \right) = \int_{\Omega} |\nabla u_{\delta}|^2 - \int_{\Omega} V_{\delta} u_{\delta}^2$$

$$\geq \int_{\Omega} |\nabla u_{\delta}|^2 - \int_{\Omega} V_{\delta}' u_{\delta}^2$$

$$\geq \lambda(V_{\delta}').$$

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Indeed

$$\lambda(V) = \inf_{u \in W_0^{1,2}, \int_{\Omega} u^2 = 1} \left( \int_{\Omega} |\nabla u|^2 - \int_{\Omega} V_{\delta} u^2 \right) = \int_{\Omega} |\nabla u_{\delta}|^2 - \int_{\Omega} V_{\delta} u_{\delta}^2$$

$$\geq \int_{\Omega} |\nabla u_{\delta}|^2 - \int_{\Omega} V_{\delta}' u_{\delta}^2 + remainder$$

$$\geq \lambda(V_{\delta}') + remainder.$$

## The remainder

For the remainder, recall the quantitative bathtub principle; this gives

$$V_{\delta}' = \{u_{\delta} > c_{\delta}\}$$

and

$$\int_{\Omega} V_{\delta}' u_{\delta}^2 \geq \int_{\Omega} V_{\delta} u_{\delta}^2 + C \|V_{\delta} - V_{\delta}'\|_{L^1(\Omega)}^2$$

and, in turn,

$$\lambda(V_{\delta}) \geq \lambda(V_{\delta}') + C\|V_{\delta}' - V_{\delta}\|_{L^{1}(\Omega)}^{2}.$$

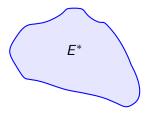


Figure: Depiction of the optimal set  $E^* = \{u_{V^*} > c^*\}$ .

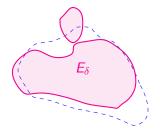


Figure: Depiction of the set  $E_{\delta}$  s.t.  $V_{\delta}=\mathbb{1}_{E_{\delta}}$ ;  $E^*$  is depicted in dashed blue.

We replace  $E_{\delta}$  with  $E'_{\delta} = \{u_{E_{\delta}} > c_{\delta}\}$ . We have

$$\lambda(V_{\delta}) - \lambda(V_{\delta}') \ge C \|\mathbb{1}_{E_{\delta}} - \mathbb{1}_{E_{\delta}'}\|_{L^{1}(\Omega)}^{2}.$$
 (5)

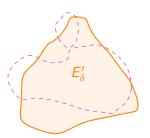


Figure: Level set  $E'_{\delta}$  of the eigenfunction  $u_{\delta}$  associated with  $E_{\delta}$  and satisfying the volume constraint;  $E_{\delta}$  is depicted in dashed magenta.

But now, by elliptic regularity,  $u_{\delta} \approx u_{V^*}$  in  $W^{2,p}$ .

We use the quantitative inequality for deformation of sets:

$$\lambda(V_{\delta}') - \lambda(V_{\delta}) \ge C \|\mathbb{1}_{E_{\delta}'} - \mathbb{1}_{E^*}\|_{L^1(\Omega)}^2. \tag{5}$$

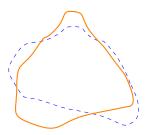


Figure: Comparison of  $E_{\delta}$ " (in orange) and of  $E^*$  (in dashed blue).

Combining these two steps we obtain

$$\lambda(V_{\delta}) - \lambda(V^*) \ge C \left( \|\mathbb{1}_{E_{\delta}'} - \mathbb{1}_{E^*}\|_{L^1(\Omega)}^2 + \|\mathbb{1}_{E_{\delta}} - \mathbb{1}_{E_{\delta}'}\|_{L^1(\Omega)}^2 \right). \tag{5}$$

Since

$$\|\mathbb{1}_{E^*} - \mathbb{1}_{E_{\delta}}\|_{L^1(\Omega)} = \delta \tag{6}$$

we have either (up to a subsequence)

$$\|1_{E^*} - 1_{E'_{\delta}}\|_{L^1(\Omega)} \ge c_0 \delta$$

or

$$\|\mathbb{1}_{E_{\delta}}-\mathbb{1}_{E_{\delta}'}\|_{L^{1}(\Omega)}\geq c_{1}\delta$$

so the contradiction follows and the inequality is proved.

# Summarising

Recall that at an optimiser

$$V^*=\mathbb{1}_{\{u_{V^*}>c^*\}}.$$

We say that a set E is critical (or that  $V = \mathbb{1}_E$  is critical) if

$$E=\{u_{\mathbb{1}_E}>c_E\}.$$

In fact we can show that if  $V_0$  is large enough (or in the ball) then any critical set E is locally stable: in a small  $L^1$  ball,

$$\lambda(V) - \lambda(\mathbb{1}_E) \geq C \|V - \mathbb{1}_E\|_{L^1(\Omega)}^2.$$

# How to numerically approximate the optimal potentials?

$$\min_{0 \leq V \leq 1, \int V = V_0} \lambda(V) \text{ first eigenvalue of } \begin{cases} -\Delta u_V = \lambda(V) u_V + V u_V \,, \\ u_V \in W_0^{1,2}(\Omega). \end{cases}$$

Recall that at a perturbation h

$$\lambda'(V)[h] = -\int_{\Omega} h u_V^2.$$

This suggests a gradient descent/fixed-point algorithm

1: Initialisation at  $V_0$ 

2:  $k \leftarrow 0$ 

3: Compute  $u_{V_k}$ 

4: Compute  $c_k$  such that  $\operatorname{Vol}(\{u_{V_k} > c_k\}) = V_0$ .

5:  $V_k \leftarrow \mathbb{1}_{\{u_{V_k} > c_k\}}$ 

6:  $k \leftarrow k + 1$ .

#### Does this algorithm converge?

## The thresholding algorithm: is it successful?

## The answer is **yes**. Remarkably efficient:

- First derived by Céa, Gioan and Michel for some shape optimisation problems.
- Optimisation of the Dirichlet energy, of eigenvalues etc...
  - 1 Kao, Lou, Yanagida, Mathematical biosciences and engineering, 2008
  - 2 Hintermüller, Kao, Laurain, Applied Mathematics & Optimization, 2012
  - 3 Lamboley, Laurain, Nadin, Privat, Calculus of Variations & PDEs, 2016.
- Generalises to large classes of optimal control problems
  - Ding, Finotti, Lenhart, Lou, Ye, Nonlinear analysis: Real world applications, 2010,
  - 2 M-F, Nadin, Privat, Journal de Mathématiques Pures et Appliquées, 2020,
  - 3 M-F, Ruiz-Balet, SIAM Journal on Applied Mathematics, 2021,
  - Mao, Mohammadi, Journal of Mathematical Biology, 2022.
- Topology optimisation:
  - 1 Amstutz, Optimization Methods and Software, 2011
  - 2 Amstutz, Dapogny, Ferrer, Numerische Mathematik, 2018

### Main difficulties

We want to obtain convergence to a local minimiser. Let's list the difficulties:

**Order of the algorithm:** First-order algorithm. The most optimistic outcome is that we find a critical point. Here that means finding  $V=\mathbbm{1}_F$  with

$$E = \{u_V > c_E\}.$$

② Degenerate minimisers and regularity properties: No unambiguous notion of "non-degenerate critical point". Quantitative inequalities will play the role of non-degeneracy conditions. They require some regularity of minimisers  $E^*$ .

### What we show

#### With A. Chambolle, Y. Privat:

- **①** Convergence of the algorithm for large volume constraints  $(V_0 \sim |\Omega|)$
- 4 Holds for the Dirichlet energy, for eigenvalue optimisation, for some classes of non-energetic optimal control problems.
- First "general" convergence result.
  - Kao, Mohammadi, Osting, Journal of Scientific Computing, 2021: linear convergence of rearrangement algorithms in the one-dimensional case, based on explicit computations.

#### A related scheme

$$V_k \Rightarrow -\Delta u_k = \Psi(V_k, u_{V_k}) \Rightarrow V_{k+1} = \mathbb{1}_{\{u_k > c_k\}}$$
  
Or, equivalently,  $u_k = G_k \star \Psi(u_k, V_k), V_{k+1} = \mathbb{1}_{\{u_k > c_k\}}$ 

with  $G_k$  the Green kernel of a certain operator.

Falls in the category of thresholding schemes, the main one being the Bence-Merriman-Osher approximation of the mean curvature flow.

- 1 Merriman, Bence, Osher. Diffusion generated motion by mean curvature, 1992,
- Bellettini, Caselles, Chambolle, Novaga, Journal de Mathématiques Pures et Appliquées, 2009,
- 3 Esedoglu, Otto, Communications on Pure and Applied Mathematics, 2015,
- 4 Laux, Otto, Calculus of Variations and Partial Differential Equations, 2016.

Main problem here: the kernel can depend on the iteration and presence of boundary conditions.

### A rough idea of the proof

Recall:  $\lambda$  is concave. In particular,

$$\lambda(V_{k+1}) - \lambda(V_k) \leq \lambda'(V_k)[V_{k+1} - V_k] = \int_{\Omega} V_k u_k^2 - \int_{\Omega} V_{k+1} u_k^2.$$

But we can use the quantitative bathtub principle (again...)

$$\lambda(V_{k+1}) - \lambda(V_k) \le -C \|V_{k+1} - V_k\|_{L^1(\Omega)}^2.$$

Summing:

$$\sum_{k=0}^{\infty} \|V_{k+1} - V_k\|_{L^1(\Omega)}^2 < \infty.$$

# One or infinitely many closure points

$$\sum_{k=0}^{\infty} \|V_{k+1} - V_k\|_{L^1(\Omega)}^2 < \infty \Rightarrow \|V_{k+1} - V_k\|_{L^1(\Omega)} \underset{k \to \infty}{\to} 0.$$

**Conclusion:** the sequence  $\{V_k\}_{k\in\mathbb{N}}$  has exactly one or infinitely many closure points.

# One or infinitely many closure points

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**Conclusion:** the sequence  $\{V_k\}_{k\in\mathbb{N}}$  has exactly one or infinitely many closure points.

**But:** If  $V_{\infty}$  is a closure point,  $V_{\infty}$  is critical.

### One closure point

But we saw that critical points are strongly isolated in  $L^1$ . In particular, the sequence can only have one closure point, and so the algorithm converges to a stable local minimiser.

### Generalisation to optimal control problems

$$\mathcal{L}\mathbf{u} = g(\mathbf{u}) + \Phi(\mathbf{u}, V).$$

- Ω: bounded smooth domain
- 2 L: differential operator
- u: state
- g: semilinearity
- Φ: state/control coupling.

### Generalisation to optimal control problems

$$\mathcal{L}\mathbf{u} = g(\mathbf{u}) + \Phi(\mathbf{u}, V).$$

- Ω: bounded smooth domain
- ② L: differential operator
- u: state
- g: semilinearity
- V: control
- Φ: state/control coupling.

$$\max / \min_{f ext{ admissible control}} J(V) = \begin{cases} ext{Energy of the PDE,} \\ \int_{\Omega} j(x, u_V), \\ ext{Eigenvalue.} \end{cases}$$

#### Constraints on the controls

**1** Pointwise  $(L^{\infty})$  constraints:

$$0 \le V \le 1$$
 a.e..

**Q** Global  $(L^1)$  constraints:

$$\int_{\Omega}V=V_{0}.$$

**3** We let  $\mathcal{V}$  be the set of admissible controls.

$$\min_{0 \le V \le 1, \int V = V_0} J(V) = \begin{cases} \text{Energy of the PDE,} \\ \int_{\Omega} j(x, u_V), \\ \text{Eigenvalue,} \end{cases} \begin{cases} \mathcal{L}u = \Phi(u, V), \\ u_V \in W_0^{1,2}(\Omega). \end{cases}$$

We want to **compute good approximations of**  $V \leadsto \text{fixed-point}$  on first order optimality conditions/gradient descent.

$$\min_{0 \le V \le 1, \int V = V_0} J(V) = \begin{cases} \text{Energy of the PDE,} \\ \int_{\Omega} j(x, u_V), \\ \text{Eigenvalue,} \end{cases} \begin{cases} \mathcal{L}u = \Phi(u, V), \\ u_V \in W_0^{1,2}(\Omega). \end{cases}$$

**Adjoint state:** V admissible control. If J is smooth enough there exists  $p_V$  (the adjoint state) such that for any h such that  $V + \varepsilon h \in \mathcal{F}$ ,

$$J'(f)[h] = -\int_{\Omega} \rho_V h.$$

**Iterative scheme:**  $V_k$  given,  $h = V - V_k \rightsquigarrow V_{k+1}$  chosen as a solution of

$$\min_{V \in \mathcal{V}} - \int_{\Omega} p_{V_k} (V - V_k) \Leftrightarrow \min_{V \in \mathcal{V}} - \int_{\Omega} p_{V_k} V.$$

Conclusion: by the bathtub principle

$$V_{k+1} = \mathbb{1}_{\{p_{V_k} > c_k\}}$$

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$$\min_{0 \le V \le 1, \int V = V_0} J(V) = \begin{cases} \text{Energy of the PDE,} \\ \int_{\Omega} j(x, u_V), \\ \text{Eigenvalue,} \end{cases} \begin{cases} \mathcal{L}u = \Phi(u, V), \\ u_V \in W_0^{1,2}(\Omega). \end{cases}$$

$$J'(V)[h] = -\int_{\Omega} \frac{p_V}{h}.$$

- 1: Initialisation at  $V_0 \in \mathcal{V}$
- 2:  $k \leftarrow 0$
- 3: Compute  $p_{V_{\nu}}$
- 4: Compute  $c_k$  such that  $Vol(\{p_{V_k} > c_k\}) = V_0$ .
- 5:  $V_k \leftarrow \mathbb{1}_{\{p_{V_k} > c_k\}}$
- 6:  $k \leftarrow k + 1$ .

$$\min_{0 \le V \le 1, \int V = V_0} J(V) = \begin{cases} \text{Energy of the PDE,} \\ \int_{\Omega} j(x, u_V), \\ \text{Eigenvalue,} \end{cases} \begin{cases} \mathcal{L}u = \Phi(u, V), \\ u_V \in W_0^{1,2}(\Omega). \end{cases}$$

$$J'(V)[h] = -\int_{\Omega} \mathbf{p}_V h.$$

- 1: Initialisation at  $V_0 \in \mathcal{V}$
- 2:  $k \leftarrow 0$
- 3: Compute  $p_{V_k}$  as the solution of a PDE
- 4: Compute  $c_k$  such that  $Vol(\{p_{V_k} > c_k\}) = V_0$ .
- 5:  $V_k \leftarrow \mathbb{1}_{\{p_{V_k} > c_k\}}$
- 6:  $k \leftarrow k + 1$ .

Usually,  $p_V$  solves a PDE.

#### What we show

Essentially: if J is concave and  $V_0$  is large enough, we have convergence to stable minimisers.

#### Perspectives

- Lifting the large volume constraints for the study of regularity of optimisers?
- General regularity theory?
- Frank-Wolfe type methods?
- Explicit rates of convergence?

Thank you!

#### The shape Lagrangian

- ① Critical set: Lagrange multiplier c s.t.  $E^* = \{u_{1_{E^*}} > c\}$ .
- 2 Shape Lagrangian:

$$L_{E^*}(F) = \mathscr{E}(F) + Vol(F).$$

4 Hadamard formula.

#### Computation of the first order shape derivative

- **1** :  $E^*$  fixed critical set,  $E^* = \{u_{E^*} > c\}$ .  $\Phi$ : smooth enough vector field.  $E_t := (Id + t\Phi)E^*$ .
- ② Shape derivative  $u'_{\Phi}$ : derivative at t = 0 of  $t \mapsto u_{E_t}$ :

$$\int_{\Omega} \langle \nabla u, \nabla v \rangle = \int_{E_t} v \Rightarrow \int_{\Omega} \langle \nabla u', \nabla v \rangle = \int_{\partial E^*} v \langle \Phi, \nu \rangle.$$

Rewrites as

$$\begin{cases} -\Delta u' = 0 \text{ in } E^* \cup (E^*)^c, \\ [\partial_{\nu} u'] = -\langle \Phi, \nu \rangle \end{cases}$$

First order derivative of the Lagrangian:

$$\begin{split} L_{E^*}(E^*)' &= \int \langle \nabla u, \nabla u' \rangle - \int_{E^*} u' - \int_{\partial E^*} u \langle \Phi, \nu \rangle + c^* \int_{\partial E^*} \langle \Phi, \nu \rangle \\ &= 0 \text{ since } u_{\mathbb{1}_{E^*}} = c^* \text{ on the boundary.} \end{split}$$

#### Computation of the second order shape derivative

- **1** :  $E^*$  fixed critical set,  $E^* = \{u_{E^*} > c\}$ .  $\Phi$ : smooth enough vector field.  $E_t := (Id + t\Phi)E^*$ .
- 2

$$\begin{cases} -\Delta u' = 0 \text{ in } E^* \cup (E^*)^c, \\ [\partial_{\nu} u'] = -\langle \Phi, \nu \rangle \end{cases}$$

 $\bigcirc$  For a general E,

$$L'_{E^*}(E) = \int_{\partial E} (u_E - c^*) \langle \Phi, \nu \rangle.$$

Second order derivative of the Lagrangian (Hadamard second formula):

$$L_{E^*}(E^*)'' = \int_{\partial E^*} u' \langle \Phi, \nu \rangle + \int_{\partial E^*} \frac{\partial u_{E^*}}{\partial \nu} \langle \Phi, \nu \rangle^2.$$

**3** We want to find a diagonalisation basis for the shape Hessian: the expression above is not explicit enough when  $\Omega$ ,  $E^*$  are not balls.

#### Diagonalising the shape hessian

$$L_{E^*}(E^*)'' = \int_{\partial E^*} u' \langle \Phi, \nu \rangle + \int_{\partial E^*} \frac{\partial u_{E^*}}{\partial \nu} \langle \Phi, \nu \rangle^2 \text{ with } \begin{cases} -\Delta u' = 0 \text{ in } E^* \cup (E^*)^c, \\ [\![ \partial_\nu u' ]\!] = -\langle \Phi, \nu \rangle \end{cases}$$

Diagonalisation basis:

$$\begin{cases} -\Delta \psi_k = 0 & \text{in } E^* \cup (E^*)^c, \\ [\![ \partial_\nu \psi_k ]\!] = -\sigma_k \frac{1}{|\partial_\nu u|} \psi_k \end{cases}$$

- Pound by looking for a suitable basis.
- In this basis we obtain

$$L_{E^*}^{"} = \sum \alpha_k^2 \left( 1 - \frac{1}{\sigma_k} \right).$$

- Two consequences:
  - **1** Best coercivity norm:  $L^2(\partial E^*)$
  - The coercivity is equivalent to requiring that

 $\sigma_1 > 1$ .

#### Consequence of large volume constraints II: coercivity of shape hessians

To obtain coercivity it suffices that

$$\sigma_1 > 1 \text{ with } \sigma_1 = \inf_{v \,,\, \int_{\partial E^*} v^2 > 0} \frac{\int_{\Omega} |\nabla v|^2}{\int_{\partial E^*} \frac{1}{|u_{\nu}|} v^2}.$$

2 By regularity and  $u_{\nu} \neq 0$  this is implied by

$$\mu_1\gg 1$$
 with  $\mu_1=\inf_{v\,,\int_{\partial E^*}v^2>0}rac{\int_\Omega |
abla v|^2}{\int_{\partial E^*}v^2}$ 

**3** However, we have, for  $v_1$  the eigenfunction:  $\int_{\partial E^*} v^2 = 1$ ,  $\partial E^*$  close to  $\partial \Omega$  (large volume),  $v_1 = 0$  on  $\partial \Omega$ . We can then show that

$$\mu_1 \to \infty$$
 as  $V_0 \to |\Omega|$ .

Once we have coercivity, we can adapt the tools of Dambrine & Lamboley 2019 → critical points are isolated.