

# Mean-field limits for Coulomb-type dynamics via the modulated energy method

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# The discrete coupled ODE system

Energy

$$H_N(x_1, \dots, x_N) = \frac{1}{2} \sum_{1 \leq i \neq j \leq N} g(x_i - x_j), \quad x_i \in \mathbb{R}^d$$

$$\begin{cases} g(x) = \frac{1}{s|x|^s} & s < d & \text{Riesz case} \\ g(x) = -\log|x| & s = 0 & \text{log case} \end{cases}$$

Model case

$$s = d - 2 \quad \text{Coulomb}$$

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Evolution equation

$$\dot{x}_i = -\frac{1}{N} \nabla_i H_N(x_1, \dots, x_N) \quad \text{gradient flow}$$

$$\dot{x}_i = -\frac{1}{N} \mathbb{J} \nabla_i H_N(x_1, \dots, x_N) \quad \text{conservative flow} \quad (\mathbb{J}^T = -\mathbb{J})$$

$$\ddot{x}_i = -\frac{1}{N} \nabla_i H_N(x_1, \dots, x_N) \quad \text{Newton's law}$$

possibly with added noise  $\sqrt{\theta} dW_i^t$ ,  $N$  independent Brownian motions,  $\theta$ =temperature

# Questions

For a general system

$$\dot{x}_i = \frac{1}{N} \sum_{j \neq i} K(x_i - x_j) + \sqrt{\theta} dW_i^t$$

- What is the limit of the **empirical measure**? Is there  $\mu^t$  such that for each  $t$

$$\frac{1}{N} \sum_{i=1}^N \delta_{x_i^t} \rightharpoonup \mu^t$$

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- if  $f_N^0(x_1, \dots, x_N)$  is the probability density of position of the system at time 0, what is the limit behavior of  $f_N^t$ ?
- **propagation of chaos** (Boltzmann, Kac, Dobrushin): if  $f_N^0(x_1, \dots, x_N) \simeq \mu^0(x_1) \dots \mu^0(x_N)$  is it true that

$$f_N^t(x_1, \dots, x_N) \simeq \mu^t(x_1) \dots \mu^t(x_N)?$$

in the sense of convergence of the  $k$ -point marginal  $f_{N,k}$ .

## Formal limit

We *formally* expect  $\mu_N^t := \frac{1}{N} \sum_{i=1}^N \delta_{x_i^t} \rightharpoonup \mu^t$  where  $\mu^t$  solves the mean-field equation

$$\partial_t \mu = \operatorname{div} ((K * \mu) \mu) + \frac{1}{2} \theta \Delta \mu \quad (\text{MF})$$

So for us

$$\partial_t \mu + \operatorname{div} ((\nabla g * \mu) \mu) = \frac{1}{2} \theta \Delta \mu$$

## How to prove propagation of chaos?

- ▶ Classical method [Mc Kean, Sznitman] for **Lipschitz** interaction kernel.
- ▶ convergence in a good metric, like Wasserstein [Braun-Hepp, Dobrushin, Neunzert-Wick, Hauray '09, Carrillo-Choi-Hauray '14, Carrillo-Ferreira-Precioso '12, Berman-Onnheim '15]



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- ▶ **relative entropy method**: show a Gronwall relation for

$$0 \leq H_N(f_N | \mu^{\otimes N}) := \frac{1}{N} \int_{(\mathbb{R}^d)^N} f_N \log \frac{f_N}{\mu^{\otimes N}} dx_1 \dots dx_N.$$

[Jabin-Wang '16] for  $\theta > 0$ , kernel not too irregular.

# The modulated energy method

[S, Duerinckx] Introduced for Ginzburg-Landau vortex dynamics.

Use **interaction-based metric**:

$$\|\mu - \nu\|^2 = \iint_{\mathbb{R}^d \times \mathbb{R}^d} g(x - y) d(\mu - \nu)(x) d(\mu - \nu)(y).$$

Observe weak-strong uniqueness property of the solutions to (MF) for  $\|\cdot\|$ :

$$\|\mu_1^t - \mu_2^t\|^2 \leq e^{Ct} \|\mu_1^0 - \mu_2^0\|^2 \quad C = C(\|\nabla^2(g * \mu_2)\|_{L^\infty})$$

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In the discrete case, let  $X_N$  denote  $(x_1, \dots, x_N)$  and take for **modulated energy**,

$$F_N(X_N, \mu) = \iint_{\mathbb{R}^d \times \mathbb{R}^d \setminus \Delta} g(x - y) d\left(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu\right)(x) d\left(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu\right)(y)$$

where  $\Delta$  denotes the diagonal in  $\mathbb{R}^d \times \mathbb{R}^d$ , and  $\mu = \mu^t$  solves mean-field equation.

Analogy with “relative entropy” and “modulated entropy” methods [Dafermos '79] [DiPerna '79] [Yau '91] [Brenier '00]....

## The commutator estimate / main functional inequality

Computing, we find that along solutions  $X_N^t$  of the ODE system with  $\theta = 0$ ,

$$\frac{d}{dt} F_N(X_N^t, \mu^t) \leq \iint_{\mathbb{R}^d \times \mathbb{R}^d \setminus \Delta} (v(x) - v(y)) \cdot \nabla g(x - y) d\left(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu\right)^{\otimes 2}(x, y)$$

with  $v = \nabla g * \mu^t$ .

# The commutator estimate / main functional inequality

Theorem (S, Nguyen-Rosenzweig-S, Rosenzweig-S, Hess-Childs-Rosenzweig-S)

*All Riesz cases  $s < d$ ,  $v$  Lipschitz vector field.*

$$\begin{aligned} \iint_{\mathbb{R}^d \times \mathbb{R}^d \setminus \Delta} (v(x) - v(y)) \cdot \nabla g(x - y) d\left(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu\right)^{\otimes 2}(x, y) \\ \leq C \|Dv\|_{L^\infty} (F_N(X_N, \mu) + N^{-1+\frac{s}{d}}). \end{aligned}$$

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Proof by electric formulation + stress-energy tensor structure or by commutator estimates

Why commutator ? Let  $f = \frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu$

$$\int v \cdot \nabla (g * f) - g * (\nabla \cdot (vf)) = \langle f, \left[ v, \frac{\nabla}{(-\Delta)^{\frac{d-s}{2}}} \right] f \rangle_{L^2}$$

Estimate used to treat the *quantum* Coulomb mean-field limit [Golse-Paul, Rosenzweig, Ben Porat], *quasi-neutral* limits [Iacobelli-Han Kwan, Rosenzweig, RS]

Theorem (S '18, H. Q. Nguyen-Rosenzweig-S '21, HCRS '25)

Case  $\theta = 0$ . Assume (MF) admits a regular enough (in particular  $\mu^t \in L^\infty([0, T], L^\infty(\mathbb{R}^d))$ ) solution. There exist constants  $C_1, C_2$  depending on the norms of  $\mu^t$  and  $\gamma > 0$  depending on  $d, s$ , s.t.  $\forall t \in [0, T]$

$$|F_N(X_N^t, \mu^t)| \leq \left( |F_N(X_N^0, \mu^0)| + C_1 N^{-1+\frac{s}{d}} \right) e^{C_2 t}.$$

In particular, if  $\mu_N^0 \rightharpoonup \mu^0$  and is such that

$$\lim_{N \rightarrow \infty} F_N(X_N^0, \mu^0) = 0,$$

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# The electric rewriting of the energy

Set  $h^f = g * f$ . In the Coulomb case

$$-\Delta h^f = c_d f$$

We have by IBP

$$\iint_{\mathbb{R}^d \times \mathbb{R}^d} g(x-y) df(x) df(y) = \int_{\mathbb{R}^d} h^f df = -\frac{1}{c_d} \int_{\mathbb{R}^d} h^f \Delta h^f = \frac{1}{c_d} \int_{\mathbb{R}^d} |\nabla h^f|^2.$$

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Positivity of  $F_N$  not clear! Use suitable **truncations** obtained by replacing  $\delta_{x_i}$  by  $\delta_{x_i}^{(r_i)}$  with  $r_i$  = nearest neighbor distance. Use almost-**monotonicity** with respect to truncation parameter.

## Proof in the Coulomb case

**Stress-energy tensor**

$$[\nabla h^f]_{ij} = 2\partial_i h^f \partial_j h^f - |\nabla h^f|^2 \delta_{ij}.$$

For regular  $f$ ,

$$\operatorname{div} [\nabla h^f] = 2\Delta h^f \nabla h^f = -\frac{2}{c_d} f \nabla h^f.$$

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$$\begin{aligned} & \iint_{\mathbb{R}^d \times \mathbb{R}^d} (v(x) - v(y)) \cdot \nabla g(x - y) df(x) df(y) \\ &= 2 \int_{\mathbb{R}^d} v(x) \cdot \nabla h^f(x) df(x) = -c_d \int_{\mathbb{R}^d} v \cdot \operatorname{div} [\nabla h^f] \\ &= c_d \int_{\mathbb{R}^d} Dv \cdot [\nabla h^f] \leq c_d \|Dv\|_{L^\infty} \iint g(x - y) df(x) df(y) \end{aligned}$$

Then needs to be **renormalized** by truncation procedure.

# The truncation method

[Hess-Childs - Rosenzweig - S '25]

Use truncation scale  $\eta = N^{-1/d}$ . Find “truncation”  $g_\eta$  satisfying

- ▶  $g_\eta \leq g$
- ▶  $\hat{g}_\eta \geq 0$
- ▶  $|g - g_\eta| \leq C_\gamma \frac{\eta^\gamma}{|x|^{s+\gamma}}$  for any  $\gamma > s$ .

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Set

$$g_\eta(x) = c_{\phi,d,s} \int_\eta^\infty t^{d-s-1} \phi_t(x) dt$$

with  $\phi_t = t^{-d} \phi(\frac{\cdot}{t})$  and  $\phi$  is a Bessel potential, defined by  $\hat{\phi}(\xi) = (1 + 4\pi^2 |\xi|^2)^{-s/2}$ ,  
hence fundamental solution to  $(-\Delta + I)^{s/2}$

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Handle  $\iint \nabla g_\eta(x-y) \cdot (v(x) - v(y)) d(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu)^{\otimes 2}(x, y)$  by stress-tensor structure + Kato-Ponce commutator estimates, using Bessel nature.



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Handle  $\iint \nabla g_\eta(x-y) \cdot (v(x) - v(y)) d(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu)^{\otimes 2}(x, y)$  by stress-tensor structure + Kato-Ponce commutator estimates, using Bessel nature.

Control  $\iint \nabla (g - g_\eta)(x-y) \cdot (v(x) - v(y)) d(\frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu)^{\otimes 2}(x, y)$  more brutally by small-scale interaction control obtained by decomposing  $g$  as  $g_\eta + (g - g_\eta)$  in definition of  $F_N$ .

# Modulated free energy

In the cases of gradient flow **with noise**

**Modulated free energy** method [Bresch-Jabin-Wang]

$$\mathcal{F}_N^\theta(f_N, \mu) := \theta H_N(f_N | \mu^{\otimes N}) + \int F(X_N, \mu) df_N(X_N).$$

Introduce **modulated Gibbs measure**

$$\mathbb{Q}_{N,\theta}(\mu) = \frac{1}{K_{N,\theta}(\mu)} e^{-\frac{1}{\theta} N F_N(X_N, \mu)} d\mu(x_1) \dots d\mu(x_N)$$

Then (remark in [Rosenzweig-S])

$$\mathcal{F}_N^\theta(f_N, \mu) = \theta H_N(f_N | \mathbb{Q}_{N,\theta}(\mu)) - \underbrace{\frac{1}{N} \log K_{N,\theta}(\mu)}_{o(1) \text{ constant}}$$

# Evolution of modulated free energy

When  $\mu^t$  solves (MF)

$$\begin{aligned} \frac{d}{dt} \mathcal{F}_N^\theta(f_N^t, \mu^t) = & \underbrace{-\frac{\theta^2}{N} \int \left| \nabla \sqrt{\frac{f_N^t}{\mathbb{Q}_{N,\theta}(\mu^t)}} \right|^2 d\mathbb{Q}_{N,\theta}(\mu^t)}_{\text{relative Fisher information}} \\ & + \underbrace{\frac{1}{2} \int df_N^t \iint_{\Delta^c} (v^t(x) - v^t(y)) \cdot \nabla g(x - y) d \left( \frac{1}{N} \sum_{i=1}^N \delta_{x_i} - \mu^t \right)^{\otimes 2}}_{\text{commutator term}}(x, y) dx dy \end{aligned}$$

where  $v^t = \nabla g * \mu^t + \theta \nabla \log \mu^t$

## Global in time convergence?

**Optimized** version of the functional inequality using that  $v^t = \theta \nabla \log \mu^t + \nabla g * \mu^t$

$$\frac{d}{dt} \mathcal{F}_N^\theta(X_N^t, \mu^t) \leq C \|\nabla v^t\|_{L^\infty} \left( \mathcal{F}_N^\theta(X_N^t, \mu^t) + N^{\frac{s}{d}-1} \|\mu^t\|_{L^\infty}^{s/d} \right),$$

$\rightsquigarrow$  prove and exploit the **decay** rate of  $\nabla v^t$  as  $t \rightarrow \infty$  ?

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$\rightsquigarrow$  prove and exploit the **decay** rate of  $\nabla v^t$  as  $t \rightarrow \infty$  ? Works in the torus setting with exponential decay rate

Theorem (Chodron de Courcel - Rosenzweig - S '23)

*Riesz case  $s \in [d-2, d)$ , gradient flow with additive noise on the torus. We have global in time convergence:*

$$\mathcal{F}_N^\theta(f_N^t, \mu^t) \leq C \left( \mathcal{F}_N^\theta(f_N^0, \mu^0) + N^{\frac{s}{d}-1} \right).$$

In whole space, problem because  $\|\mu\|_{L^\infty} \rightarrow 0$  due to lack of confinement.  
Replace with use of **self-similar transformation** [Rosenzweig-S '24]

$$\xi := \frac{x}{\sqrt{t+1}} \quad \tau := \log(t+1)$$

transforms the equation into

$$\partial_\tau \mu + \operatorname{div} \left( \left( -\nabla g * \mu - \nabla \left( \frac{1}{4} |\xi|^2 \right) \right) \mu \right) = \frac{1}{2} \theta \Delta \mu$$

i.e. new added quadratic confining potential, equilibrium = Gaussian.  
Modulated free energy transforms well too.

## Approach by functional inequalities

- ▶ [Guillin-Le Bris-Monmarché] for relative entropy method : works for 2D point vortex system on the torus, [Gong-Wang-Xie] same on the whole plane
- ▶ [Rosenzweig-S '23] for modulated free energy method  
Exploit Fisher information term to obtain global exponential convergence under some **modulated Logarithmic Sobolev Inequality** assumption:

$$\underbrace{N H_N(f_N^t | \mathbb{Q}_{N,\theta}(\mu^t))}_{\text{relative entropy}} \leq C_{LS} \underbrace{\int \left| \nabla \sqrt{\frac{f_N^t}{\mathbb{Q}_{N,\theta}(\mu^t)}} \right|^2 d\mathbb{Q}_{N,\theta}(\mu^t)}_{\text{relative Fisher}}$$

then this + commutator estimate gives

$$\frac{d}{dt} \mathcal{F}_N^\theta \leq -C \mathcal{F}_N^\theta + o(1)$$

exponential convergence to tensorized state (**generation of chaos**)

THANK YOU FOR YOUR ATTENTION!