

Relative entropy and particle systems

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Mardi 3 juin 2025

1 Relative entropy

2 Particle system

Definition

Let μ and ν be two probability measures with densities f and g with respect to Lebesgue measure on $\mathbb{R}$. The relative entropy of ν with respect to μ is

$$H(\nu|\mu) = \int_{\mathbb{R}} f \log \left(\frac{f}{g} \right) dx$$

We can also write it, if $h = f/g$ is the density of ν with respect to μ :

$$H(\nu|\mu) = \int h \log(h) d\mu$$

Proposition

Due to Jensen's inequality, the relative entropy is nonnegative.

Proposition

$$H(\nu|\mu) = \sup_{\varphi} \left\{ \int \varphi d\nu - \log \left(\int e^{\varphi} d\mu \right) \right\}$$

Proof : take a function φ . Consider the probability measure $\tilde{\mu}$ defined as

$$d\tilde{\mu}(x) = \frac{e^{\varphi(x)}}{\int e^{\varphi(y)} d\mu(y)} d\mu(x) = \frac{e^{\varphi(x)} g(x)}{\int e^{\varphi(y)} g(y) dy} dx$$

Variational formulation

$$\begin{aligned} H(\nu|\tilde{\mu}) - H(\nu|\mu) &= \int f(x) \left(\log \left(\frac{f(x) \int e^{\varphi(y)} g(y) dy}{e^{\varphi(x)} g(x)} \right) - \log \left(\frac{f(x)}{g(x)} \right) \right) dx \\ &= \int f(x) \log \left(\frac{\int e^{\varphi(y)} g(y) dy}{e^{\varphi(x)}} \right) dx \\ &= - \int \varphi(x) f(x) dx + \log \left(\int e^{\varphi(x)} g(x) dx \right) \end{aligned}$$

so

$$H(\nu|\mu) = H(\nu|\tilde{\mu}) + \int \varphi(x) d\nu - \log \left(\int e^{\varphi} d\mu \right)$$

and $H(\nu|\tilde{\mu}) \geq 0$ with equality if $\varphi = \log \left(\frac{f}{g} \right)$. It proves the proposition.

An important consequence of the variational formulation of the relative entropy is the following theorem, called "entropic inequality".

Theorem

For all functions φ such that the quantities are well-defined, if X is a random variable with law ν , we have

$$\forall \alpha > 0, \quad \mathbb{E}[\varphi(X)] \leq \frac{1}{\alpha} \left(H(\nu|\mu) + \log \left(\int e^{\alpha \varphi} d\mu \right) \right)$$

Plan

1 Relative entropy

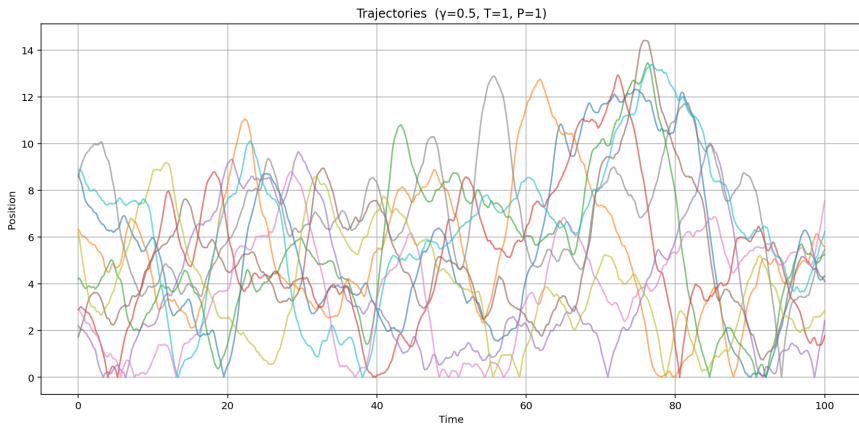
2 Particle system

We consider n particles moving on a line. q_i is the position of the particle i and p_i is its velocity. The equations of the movement are

$$\begin{aligned}dq_i(t) &= p_i(t)dt \\dp_i(t) &= \left(-2\delta_0(q_i(t))p_i(t^-) - P \left(\prod_{j=1}^n \mathbb{1}_{\{q_j(t) \leq q_i(t)\}} \right) \right) dt \\&\quad - \gamma p_i(t)dt + \sqrt{2\gamma T}dW_t^i\end{aligned}$$

where the W^i are independant Brownian motions.

Simulation



Invariant measure

We say that a measure μ is an invariant measure of the system if, when the initial configuration q_0, p_0 is distributed with this measure, then it is the case at each time t .

Theorem

Let

$$g(p, q) = \frac{(\beta P)^n}{n!} \sum_{k=1}^n e^{-\beta P q_k} \prod_{\substack{i=1 \\ i \neq k}}^n \mathbb{1}_{\{q_i \leq q_k\}} \prod_{j=1}^n \sqrt{\frac{\beta}{2\pi}} e^{-\frac{\beta p_j^2}{2}}$$

The measure $\mu_{\beta, P}^n$ such that

$$d\mu_{\beta, P} = g d\lambda_{\Gamma_n}$$

(where λ_{Γ_n} is the Lebesgue measure on $\Gamma_n = (\mathbb{R}_+)^n \times \mathbb{R}^n$) is the unique invariant measure of the system.

Generator

Let $\mathcal{L} = \mathcal{A} + \gamma \mathcal{S}$ be an operator defined on

$$\mathcal{D}_n(\mathcal{L}) = \left\{ F \in \mathcal{C}^2(\mathbb{R}_+^n \times \mathbb{R}^n), \forall j \in \llbracket 1, n \rrbracket, \right. \\ \left. F(q, p_1, \dots, p_j, \dots, p_n) \big|_{q_j=0} = F(q, p_1, \dots, -p_j, \dots, p_n) \big|_{q_j=0} \right\}$$

and such that

$$\mathcal{A} = \sum_{j=1}^n p_j \partial_{q_j} - P \left(\prod_{i \neq j} \mathbb{1}_{\{q_j \geq q_i\}} \right) \partial_{p_j}$$

$$\mathcal{S} = \sum_{j=1}^n \left(T \partial_{p_j}^2 - p_j \partial_{p_j} \right)$$

Proposition

$\mathcal{S}$ is symmetric and $\mathcal{A}$ antisymmetric with respect to $\mu_{\beta, P}^n$

Proposition

The operator $\mathcal{L}$ is the generator of the system, ie

$$\forall F \in \mathcal{D}_n(\mathcal{L}), \quad \forall t \geq 0,$$

$$\mathbb{E}[F(q(t), p(t))] = \mathbb{E}[F(q_0, p_0)] + \int_0^t \mathbb{E}[\mathcal{L}F(q(s), p(s))] ds$$

Application : if the law ν_t^n of $(q(t), p(t))$ is absolutely continuous with respect to $\mu_{\beta, p}^n$ with density $f_n(t, \cdot, \cdot)$ for all $t \geq 0$, then f_n satisfies the PDE

$$\partial_t f_n = \mathcal{L}^* f_n$$

(with $\mathcal{L} = -\mathcal{A} + \gamma\mathcal{S}$)

Important hypothesis

In the following we admit that such a density exists.

We define

$$H_n(t) = H(\nu_t^n | \mu_{\beta, P}^n) = \int_{\Gamma_n} f_n(t, q, p) \log(f_n(t, q, p)) d\mu_{\beta, P}^n.$$

Proposition

H_n is a non-increasing function.

Proof : Using the fact that $\partial_t f_n = \mathcal{L}^* f_n$, we have

$$\frac{dH_n}{dt}(t) = \int_{\Gamma_n} \mathcal{L}^* f_n \log(f_n) d\mu_{\beta, P}^n + \int_{\Gamma_n} \partial_t f_n d\mu_{\beta, P}^n$$

$$\int_{\Gamma_n} \partial_t f_n d\mu_{\beta,P}^n = \partial_t \int_{\Gamma_n} f_n d\mu_{\beta,P}^n = \partial_t 1 = 0$$

$$\int_{\Gamma_n} f_n \mathcal{A} \log(f_n) d\mu_{\beta,P}^n = \int_{\Gamma_n} \mathcal{A} f_n d\mu_{\beta,P}^n = 0$$

Then

$$\begin{aligned} \frac{dH_n}{dt}(t) &= \gamma \int (\mathcal{S} f_n) \log(f_n) d\mu_{\beta,P} \\ &= \gamma \sum_{j=1}^n T \int \partial_{p_j}^2 f_n \log(f_n) d\mu_{\beta,P} - \gamma \sum_{j=1}^n \int p_j \partial_{p_j} f_n \log(f_n) d\mu_{\beta,P} \end{aligned}$$

Finally, by an integration by part :

$$\begin{aligned}\int p_j \partial_{p_j} f_n \log(f_n) d\mu_{\beta, P}^n &= T \int \partial_{p_j} (\partial_{p_j} f_n \log(f_n)) d\mu_{\beta, P}^n \\ &= T \int \left(\partial_{p_j}^2 f_n \log(f_n) + \frac{(\partial_{p_j} f_n)^2}{f_n} \right) d\mu_{\beta, P}^n\end{aligned}$$

So

$$\frac{dH_n}{dt}(t) = -T \gamma \sum_{j=1}^n \frac{(\partial_{p_j} f_n)^2}{f_n} d\mu_{\beta, P}^n \leq 0$$

For a function φ "regular enough", we set

$$\pi_n(t, \varphi) = \frac{1}{n} \sum_{i=1}^n \varphi \left(t, \frac{q_i(n^2 t)}{n} \right)$$

The goal is to compute the "limit" of π_n ; it is called the "hydrodynamic limit" of the particle system.

Conjecture

Let $\rho_0 : [0, L_0] \rightarrow \mathbb{R}_+$ be a density function. Assume that the initial configuration ν_0^n is such that

- $\forall \varphi \in \mathcal{C}_c^2(\mathbb{R}_+), \quad \pi_n(0, \varphi) \xrightarrow[n \rightarrow +\infty]{\mathbb{P}} \int_0^{L_0} \rho_0(x) \varphi(x) dx$
- $H_n = H(\nu_0^n | \mu_{\beta, P}^n) \underset{n \rightarrow +\infty}{=} O(n).$

Then, for all $\varphi \in \mathcal{C}^{1,2}(\mathbb{R}_+ \times \mathbb{R}_+)$ compactly supported in space,

$$\pi_n(t, \varphi) \xrightarrow[n \rightarrow +\infty]{\mathbb{P}} \int_0^{L(t)} \rho(t, x) \varphi(t, x) dx$$

where (ρ, L) is the unique weak solution of the following free boundary heat equation.

Free boundary heat equation

$$\begin{aligned}\partial_t \rho &= \frac{T}{\gamma} \partial_x^2 \rho & t \in \mathbb{R}_+, x \in [0, L(t)] \\ L'(t) &= -\frac{T^2}{\gamma P} \partial_x \rho(t, L(t)) & t \geq 0 \\ \partial_x \rho(t, 0) = 0 & \quad \rho(t, L(t)) = \frac{P}{T} & \rho(0, x) = \rho_0(x)\end{aligned}$$

Definition

(ρ, L) is a weak solution of this equation if

- $L \in \mathcal{C}(\mathbb{R}_+, \mathbb{R}_+)$, $\rho \in \mathcal{C}(\{(t, x) \in (\mathbb{R}_+)^2, x < L(t)\}, \mathbb{R}_+)$
- For all $t \geq 0$ and all test function $\varphi \in \mathcal{C}^{1,2}(\mathbb{R}_+ \times \mathbb{R}_+)$ compactly supported in space such that $\partial_x \varphi(\cdot, 0) = \varphi(t, \cdot) = 0$,

$$\begin{aligned}- \int_0^t \int_0^{L(s)} \rho(s, x) \partial_t \varphi(s, x) dx ds - \frac{T}{\gamma} \int_0^t \int_0^{L(s)} \rho(s, x) \partial_x^2 \varphi(s, x) dx ds \\ = -\frac{P}{\gamma} \int_0^t \partial_x \varphi(s, L(s)) + \int_0^{L_0} \rho_0(x) \varphi(0, x) dx.\end{aligned}$$

Free boundary heat equation

Theorem

The equation admits at most one weak solution.

Sketch of the proof

Fix $t \geq 0$ and $\omega \in \Omega$. We can remark that $\pi_n(t, \cdot)(\omega)$ characterizes the probability measure on $\mathbb{R}_+$

$$\nu_{n,t,\omega}(A) = \frac{1}{n} \sum_{i=1}^n \mathbb{1}_A \left(\frac{q_i(n^2 t)(\omega)}{n} \right)$$

(Indeed, $\pi_n(t, \varphi) = \int_{\mathbb{R}_+} \varphi d\nu_{n,t,\omega}$).

Then, we can show that $\pi_n(\omega)$ defines a continuous function from $\mathbb{R}_+$ to $\mathcal{M}_1(\mathbb{R}_+)$. Denote Q_n the probability distribution of π_n on $\mathcal{C}([0, t_1], \mathcal{M}_1(\mathbb{R}_+))$. Q_n is a sequence in $\mathcal{M}_1(\mathcal{C}([0, t_1], \mathcal{M}_1(\mathbb{R}_+)))$

Sketch of the proof

There are two main steps in the proof :

- Prove that the sequence $(Q_n)_{n \in \mathbb{N}}$ is tight on $\mathcal{M}_1(\mathcal{C}([0, t_1], \mathcal{M}_1(\mathbb{R}_+)))$.
- Identify the limit point of any convergent subsequence of $(Q_n)_n$ as the weak solution of the equation.

Conclude by using the Prokhorov theorem.

Identify the limit

Suppose that for all functions $\varphi \in \mathcal{C}^{1,2}(\mathbb{R}_+ \times \mathbb{R}_+)$ compactly supported in space,

$$\pi_n(t, \varphi) \xrightarrow[n \rightarrow +\infty]{\mathbb{P}} \int_0^{L(t)} \rho(t, x) \varphi(t, x) dx$$

Set

$$\mathcal{M}_n(t, \partial_x \varphi) = \sqrt{2T} \sum_{i=1}^n \int_0^{n^2 t} \partial_x \varphi \left(s, \frac{q_i(s)}{n} \right) dW_s^i$$

By a computation, we can show :

Identify the limit

$$\begin{aligned}\pi_n(t, \varphi) - \pi_n(0, \varphi) &= \int_0^t \pi_n(s, \partial_t \varphi) ds \\ &+ \frac{1}{\gamma n^2} \left(\sum_{i=1}^n p_i(0) \partial_x \varphi \left(0, \frac{q_i(0)}{n} \right) \right. \\ &\quad \left. - \sum_{i=1}^n p_i(n^2 t) \partial_x \varphi \left(t, \frac{q_i(n^2 t)}{n} \right) \right) \\ &+ \frac{1}{\gamma n^2} \sum_{i=1}^n \int_0^t p_i(n^2 s) \partial_{xt} \varphi \left(s, \frac{q_i(n^2 s)}{n} \right) ds \\ &+ \frac{1}{\gamma n} \int_0^t \sum_{i=1}^n p_i^2(n^2 s) \partial_x^2 \varphi \left(s, \frac{q_i(n^2 s)}{n} \right) ds \\ &- \frac{P}{\gamma} \int_0^t \partial_x \varphi \left(s, \frac{q_{\max}(n^2 s)}{n} \right) ds + \frac{1}{\sqrt{\gamma} n^2} \mathcal{M}_n(t, \partial_x \varphi)\end{aligned}$$

Identify the limit

Some terms are easy to identify : by hypothesis, if $\varphi(t, \cdot) = 0$, then

$$\pi_n(t, \varphi) - \pi_n(0, \varphi) \xrightarrow{n \rightarrow +\infty} - \int_0^{L_0} \rho_0(x) \varphi(0, x)$$

$$\int_0^t \pi_n(s, \partial_t \varphi) ds \xrightarrow{n \rightarrow +\infty} \int_0^t \int_0^{L(s)} \rho(s, x) \partial_t \varphi(s, x) dx ds$$

Using some properties of the stochastic integral, we can show that

$$\frac{1}{\sqrt{\gamma} n^2} \mathcal{M}_n(t, \partial_x \varphi) \xrightarrow{n \rightarrow +\infty} 0$$

Identify the limit

$$\begin{aligned}
 \cancel{\pi_n(t, \varphi)} - \pi_n(0, \varphi) &= \int_0^t \cancel{\pi_n(s, \partial_t \varphi)} ds \\
 &+ \frac{1}{\gamma n^2} \left(\sum_{i=1}^n p_i(0) \partial_x \varphi \left(0, \frac{q_i(0)}{n} \right) \right. \\
 &\quad \left. - \sum_{i=1}^n p_i(n^2 t) \partial_x \varphi \left(t, \frac{q_i(n^2 t)}{n} \right) \right) \\
 &+ \frac{1}{\gamma n^2} \sum_{i=1}^n \int_0^t p_i(n^2 s) \partial_{xt} \varphi \left(s, \frac{q_i(n^2 s)}{n} \right) ds \\
 &+ \frac{1}{\gamma n} \int_0^t \sum_{i=1}^n p_i^2(n^2 s) \partial_x^2 \varphi \left(s, \frac{q_i(n^2 s)}{n} \right) ds \\
 &- \frac{P}{\gamma} \int_0^t \partial_x \varphi \left(s, \frac{q_{\max}(n^2 s)}{n} \right) ds + \frac{1}{\cancel{\sqrt{\gamma} n^2}} \cancel{\mathcal{M}_n(t, \partial_x \varphi)}
 \end{aligned}$$

Identify the limit

For the term

$$\frac{1}{\gamma n^2} \sum_{i=1}^n p_i(n^2 t) \partial_x \varphi \left(t, \frac{q_i(n^2 t)}{n} \right),$$

we will use the entropy inequality.

Let M be a constant such that $|\partial_x \varphi| \leq M$. We have

$$\begin{aligned} \mathbb{E} \left[\left| \frac{1}{\gamma n^2} \sum_{i=1}^n p_i(n^2 t) \partial_x \varphi \left(t, \frac{q_i(n^2 t)}{n} \right) \right|^2 \right] &\leq \frac{M^2}{\gamma^2 n^4} \mathbb{E} \left[\left(\sum_{i=1}^n p_i(n^2 t) \right)^2 \right] \\ &\leq \frac{M^2}{\gamma^2} \sum_{i=1}^n \mathbb{E} \left[\frac{p_i^2(n^2 t)}{n^3} \right] \end{aligned}$$

Identify the limit

By the entropy inequality, for all $\alpha > 0$ we have

$$\mathbb{E} \left[\frac{p_i^2(n^2 t)}{n^3} \right] \leq \frac{H_n(n^2 t)}{\alpha} + \frac{1}{\alpha} \log \left(\int e^{\alpha \frac{p_i^2}{n^3}} d\mu_{\beta, P}^n \right)$$

so

$$\sum_{i=1}^n \mathbb{E} \left[\frac{p_i^2(n^2 t)}{n^3} \right] \leq n \frac{H_n(n^2 t)}{\alpha} + \frac{n}{\alpha} \log \left(\int e^{\alpha \frac{p_1^2}{n^3}} d\mu_{\beta, P}^n \right)$$

We know that the entropy is nonincreasing and $H_n(0) \underset{n \rightarrow +\infty}{=} O(n)$ so $nH_n(n^2 t) \underset{n \rightarrow +\infty}{=} O(n^2)$. Then $nH_n(n^2 t)/\alpha \xrightarrow{n \rightarrow +\infty} 0$ if $\alpha = n^\delta$ with $\delta > 2$.

On the other hand,

$$\int e^{\alpha \frac{p_i^2}{n^3}} d\mu_{\beta, P}^n = \int_{\mathbb{R}} \frac{e^{\alpha \frac{p^2}{n^3} - \frac{p^2}{2T}}}{\sqrt{2\pi T}} dp \xrightarrow{n \rightarrow +\infty} 1$$

if $\alpha = n^\delta$ with $\delta < 3$.

Identify the limit

Take $\alpha = n^{2,5}$. We have finally

$$\sum_{i=1}^n \mathbb{E} \left[\frac{p_i^2(n^2 t)}{n^3} \right] \leq \frac{H_n(n^2 t)}{n^{1,5}} + \frac{1}{n^{1,5}} \log \left(\int e^{\frac{p_i^2}{\sqrt{n}}} d\mu_{\beta, P}^n \right) \xrightarrow{n \rightarrow +\infty} 0$$

Conclusion :

$$\mathbb{E} \left[\left| \frac{1}{\gamma n^2} \sum_{i=1}^n p_i(n^2 t) \partial_x \varphi \left(t, \frac{q_i(n^2 t)}{n} \right) \right|^2 \right] \xrightarrow{n \rightarrow +\infty} 0$$

Identify the limit

$$\begin{aligned}
 \cancel{\pi_n(t, \varphi) - \pi_n(0, \varphi)} &= \cancel{\int_0^t \pi_n(s, \partial_t \varphi) ds} \\
 &+ \cancel{\frac{1}{\gamma n^2} \left(\sum_{i=1}^n p_i(0) \partial_x \varphi \left(0, \frac{q_i(0)}{n} \right) \right.} \\
 &\quad \left. - \sum_{i=1}^n p_i(n^2 t) \partial_x \varphi \left(t, \frac{q_i(n^2 t)}{n} \right) \right) \\
 &+ \cancel{\frac{1}{\gamma n^2} \sum_{i=1}^n \int_0^t p_i(n^2 s) \partial_{xt} \varphi \left(s, \frac{q_i(n^2 s)}{n} \right) ds} \\
 &+ \frac{1}{\gamma n} \int_0^t \sum_{i=1}^n p_i^2(n^2 s) \partial_x^2 \varphi \left(s, \frac{q_i(n^2 s)}{n} \right) ds \\
 &- \frac{P}{\gamma} \int_0^t \partial_x \varphi \left(s, \frac{q_{\max}(n^2 s)}{n} \right) ds + \frac{1}{\sqrt{\gamma} n^2} \mathcal{M}_n(t, \partial_x \varphi)
 \end{aligned}$$

Identify the limit

We can use again the entropy inequality to show

$$\mathbb{E} \left[\left| \frac{1}{n} \int_0^t \sum_{i=1}^n (p_i^2(n^2 s) - T) \partial_x^2 \varphi \left(t, \frac{q_i(n^2 s)}{n} \right) ds \right| \right] \xrightarrow{n \rightarrow +\infty} 0$$

ie

$$\begin{aligned} \frac{\gamma}{n} \int_0^t \sum_{i=1}^n p_i^2(n^2 s) \partial_x^2 \varphi \left(t, \frac{q_i(n^2 s)}{n} \right) ds &\sim \frac{T}{\gamma} \int_0^t \pi_n(s, \partial_x^2 \varphi) ds \\ &\xrightarrow{n \rightarrow +\infty} \frac{T}{\gamma} \int_0^t \int_0^{L(s)} \rho(s, x) \partial_x^2 \varphi(s, x) dx ds \end{aligned}$$

Identify the limit

It remains a last term :

$$\frac{P}{\gamma} \int_0^t \partial_x \varphi \left(s, \frac{q_{\max}(n^2 s)}{n} \right) ds$$

It should give, in the weak formulation :

$$\frac{P}{\gamma} \int_0^t \partial_x \varphi(s, L(s)) ds$$

We are currently trying to show that

$$\frac{\max(q_1, \dots, q_n)(n^2 s)}{n} \xrightarrow[n \rightarrow +\infty]{} L(s)$$

but we have not finished yet.

Thank you for your attention !

All is in the title of the slide.