

# Hydrodynamic limits of particle systems

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1 A first example

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# Simple symmetric Random Walk on $\mathbb{Z}$

## Definition

We consider a sequence  $(z_i)_{i \in \mathbb{N}^*}$  of independent random variables such that

$$\mathbb{P}(z_i = -1) = \mathbb{P}(z_i = 1) = \frac{1}{2}.$$

We set, for  $n \geq 1$ ,

$$X_n = \sum_{i=1}^n z_i.$$

The sequence  $(X_n)_n$  is called a simple symmetric random walk on  $\mathbb{Z}$ .

## Theorem

We set

$$Y_n(t) = \frac{X_{[n^2 t]}}{n}.$$

Then, the process  $(Y_n)$  converges in distribution to the standard brownian motion.

Rescaling : the space becomes continuous ( $1/n$  factor), the time is accelerated ( $n^2 t$ ). It is an example of hydrodynamic limit (at diffusive scale).

Idea : we describe a behaviour at the microscopic scale and we deduce the behaviour of the system at the macroscopic scale.

# A weaker result

## Theorem

For  $f \in C_b^2(\mathbb{R})$ ,  $n \geq 1$ ,  $t \geq 0$  and  $x \in \mathbb{R}$ , we consider

$$u_n(t, x) = \mathbb{E}^x[f(Y_n(t))].$$

Then,

$$u_n \xrightarrow[n \rightarrow +\infty]{} u$$

where  $u$  is the unique solution of the heat equation

$$\begin{cases} \partial_t u &= \frac{1}{2} \partial_x^2 u \\ u(0, \cdot) &= f \end{cases}$$

It is a consequence of Donsker's theorem since the density  $p(t, \cdot)$  of the brownian motion  $B_t$  satisfies the heat equation.

# Sketch of the proof

If we show that

- The sequence  $(u_n)_n$  is relatively compact ;
- It has only one subsequential limit.

Then it proves that the sequence converges to this limit.

# Identification of the limit

We assume that  $u_n \xrightarrow[n \rightarrow +\infty]{} u$ . By definition of the system, we have

$$\begin{aligned} u_n\left(t + \frac{1}{n^2}, x\right) - u_n(t, x) &= \mathbb{E}^x \left[ f\left(Y_n\left(t + \frac{1}{n}\right)\right) - f(Y_n(t)) \right] \\ &= \mathbb{E}^x \left[ f\left(Y_n(t) + \frac{Z_{[n^2 t]}}{n}\right) - f(Y_n(t)) \right] \\ &= \mathbb{E}^x \left[ \frac{1}{2} \left( f\left(Y_n(t) + \frac{1}{n}\right) + f\left(Y_n(t) - \frac{1}{n}\right) \right) \right. \\ &\quad \left. - f(Y_n(t)) \right] \\ &= \frac{1}{n^2} \mathbb{E}^x [\mathcal{A}_n f(Y_n(t))] \end{aligned}$$

with  $\mathcal{A}_n f(x) = \frac{n^2}{2} \left( f\left(x + \frac{1}{n}\right) + f\left(x - \frac{1}{n}\right) - 2f(x) \right)$ .

## Definition

The operator  $\mathcal{A}_n$  is called the generator of the dynamic.

# Identification of the limit

Then, we have

$$u_n(t, x) = u_n(0, x) + \sum_{k=0}^{[n^2 t]-1} \frac{1}{n^2} \mathbb{E}^x \left[ \mathcal{A}_n f \left( Y_n \left( \frac{k}{n^2} \right) \right) \right].$$

and

- $\mathcal{A}_n \xrightarrow{n \rightarrow +\infty} \frac{1}{2} \partial_x^2$
- $\mathbb{E}^x [Y_n(t)] \xrightarrow{n \rightarrow +\infty} u(t, x).$

Heuristically, we recognize a Riemann sum and we obtain, taking the limit when  $n \rightarrow +\infty$  :

$$u(t, x) - u(0, x) = \frac{1}{2} \int_0^t \partial_x^2 u(t, x) dx.$$

We "differentiate" this expression : we obtain

$$\partial_t u(t, x) = \frac{1}{2} \partial_x^2 u(t, x)$$

and we obviously have  $u(0, x) = f(x)$ . Since this equation admits a unique solution, the sequence  $(u_n)_n$  has at most one subsequential limit.

# Compactness

Since  $f$  is bounded, the sets

$$A(t, x) = \{u_n(t, x), \quad n \in \mathbb{N}\}$$

are relatively compact for every  $t, x$ .

Moreover,  $(u_n)_n$  is uniformly Lipschitz. Indeed,

$$\begin{aligned} |u_n(t, x) - u_n(t, y)| &\leq \mathbb{E} \left[ \left| f \left( x + \frac{X_{[n^2 t]}}{n} \right) - f \left( y + \frac{X_{[n^2 t]}}{n} \right) \right| \right] \\ &\leq \|f'\|_\infty |x - y| \end{aligned}$$

and

$$|u_n(t, x) - u_n(s, x)| \leq \|\mathcal{A}_n f\|_\infty |t - s| \leq 2\|f\|_\infty |t - s|$$

In particular, it is equicontinuous. Then, due to Ascoli theorem, we know that  $(u_n)_n$  is relatively compact.

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# Model

$n$  particles ;  $q_j$  is the position of the  $j^{\text{th}}$  particle,  $p_j$  its velocity.

$$dq_j(t) = p_j(t) dt$$

$$dp_1(t) = -2\delta_0(q_1(t))p_1(t) dt \\ + (p_2(t^-) - p_1(t^-))\delta_0(q_2(t) - q_1(t)) dt \\ - \gamma p_1(t) dt + \sqrt{2\gamma T} dW_t^1$$

$$dp_j(t) = (p_{j-1}(t^-) - p_j(t^-))\delta_0(q_j(t) - q_{j-1}(t)) dt \\ + (p_{j+1}(t^-) - p_j(t^-))\delta_0(q_{j+1}(t) - q_j(t)) dt \\ - \gamma p_j(t) dt + \sqrt{2\gamma T} dW_t^j$$

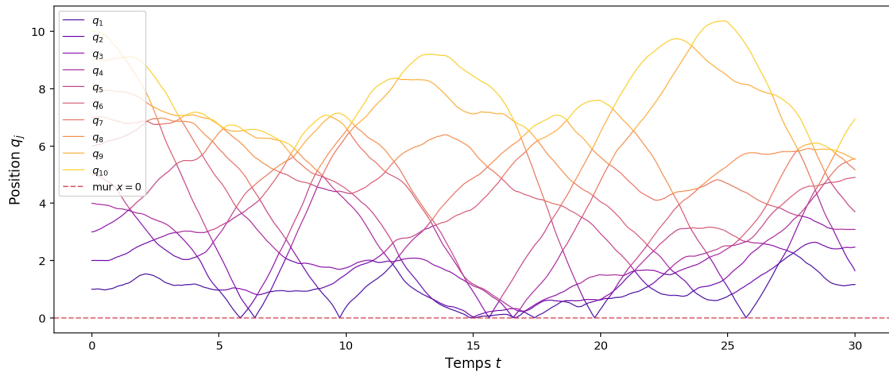
$$dp_n(t) = (p_{n-1}(t^-) - p_n(t^-))\delta_0(q_n(t) - q_{n-1}(t)) dt \\ - \gamma p_n(t) dt + \sqrt{2\gamma T} dW_t^n \quad -P dt.$$

$T$  is the temperature,  $P$  the pressure,  $\gamma$  a parameter.

$$\Gamma_n = \left\{ (q, p) \in \mathbb{R}_+^n \times \mathbb{R}^n, \quad q_1 \leq \dots \leq q_n \right\}$$

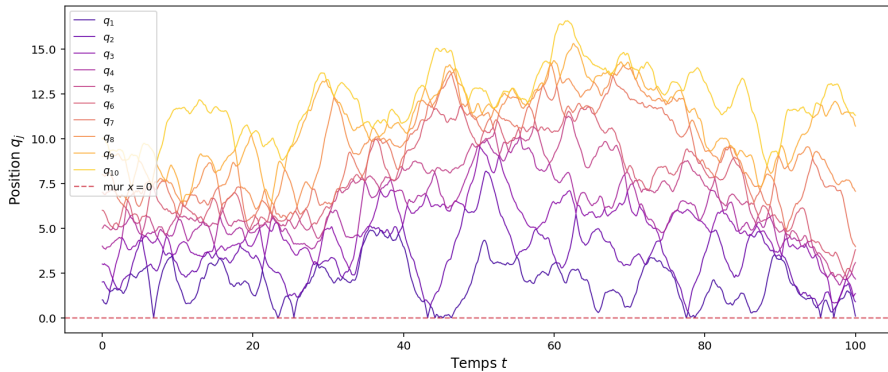
# Simulation

$n = 10$ ,  $T = 1$ ,  $\gamma = 0.1$ ,  $P = 1$ ,  $t_{\text{end}} = 30$



# Simulation

$n = 10$ ,  $T = 1$ ,  $\gamma = 1$ ,  $P = 1$ ,  $t_{\text{end}} = 100$



# Generator

The generator of this system is  $\mathcal{L} = \mathcal{A} + \gamma\mathcal{S}$ , with

$$\mathcal{A} = \sum_{j=1}^n p_j \partial_{q_j} - P \partial_{p_n};$$

$$\mathcal{S} = \sum_{j=1}^n (T \partial_{p_j}^2 - p_j \partial_{p_j}).$$

A function  $\varphi$  is in the domain  $\mathcal{D}_n(\mathcal{L})$  of the generator  $\mathcal{L}$  if

- (i) Regularity :  $\varphi$  is of class  $C^2$  on  $\Gamma_n$  with compact support ;
- (ii) Collision rule : for any  $k \in \llbracket 1, n-1 \rrbracket$ ,

$$\varphi(\dots, q_k, q_k, \dots, p_k, p_{k+1}, \dots) = \varphi(\dots, q_k, q_k, \dots, p_{k+1}, p_k, \dots);$$

- (iii) Reflexion at 0 :

$$\varphi(0, q_2, \dots, q_n, p_1, p_2, \dots, p_n) = \varphi(0, q_2, \dots, q_n, -p_1, p_2, \dots, p_n).$$

An invariant measure of this dynamic is

$$d\mu_{P,T}^n = \mathbb{1}_{\{q_1 \leq \dots \leq q_n\}} e^{-\frac{p}{T} q_n} \prod_{j=1}^n \frac{e^{-p_j^2/2T}}{\sqrt{2\pi T}} dq dp.$$

It means that if the initial configuration  $(q_1(0), \dots, q_n(0), p_1(0), \dots, p_n(0))$  is distributed by  $\mu_{P,T}^n$ , then at each time  $t \geq 0$ , we have

$$(q_1(t), \dots, q_n(t), p_1(t), \dots, p_n(t)) \sim \mu_{P,T}^n$$

## Definition

We set

$$\pi_t^n = \frac{1}{n} \sum_{j=1}^n \delta_{\frac{q_j(n^2 t)}{n}}(dx)$$

the empirical measure in the diffusive scale.

We define also

$$\langle \pi_t^n, \varphi \rangle = \int \varphi d\pi_t^n = \frac{1}{n} \sum_{j=1}^n \varphi \left( \frac{q_j(n^2 t)}{n} \right).$$

$$\begin{aligned} -\langle \pi_0^n, \psi \rangle &= \int_0^t \langle \pi_s^n, \partial_t \psi \rangle ds \\ &+ \frac{1}{n^2 \gamma} \sum_{j=1}^n \rho_j(0) \partial_x \psi \left( 0, \frac{q_j(0)}{n} \right) \\ &+ \frac{1}{n \gamma} \sum_{j=1}^n \int_0^t \rho_j^2(n^2 s) \partial_x^2 \psi \left( s, \frac{q_j(n^2 s)}{n} \right) ds \\ &- \frac{P}{\gamma} \int_0^t \partial_x \psi \left( s, \frac{q_n(n^2 s)}{n} \right) ds \\ &+ \frac{1}{n^2 \gamma} \sum_{j=1}^n \int_0^t \rho_j(n^2 s) \partial_{xt} \psi \left( s, \frac{q_j(n^2 s)}{n} \right) ds \\ &+ \frac{1}{n^2 \gamma} \mathcal{M}_{n,t}. \end{aligned}$$

# Identification of the limit

We can show that

$$\frac{1}{n\gamma} \sum_{j=1}^n p_j^2(n^2s) \partial_x^2 \psi \left( s, \frac{q_j(n^2s)}{n} \right) ds \underset{n \rightarrow +\infty}{\approx} \frac{T}{\gamma} \int_0^t \langle \pi_s^n, \partial_x^2 \psi \rangle ds$$

Then if we know that  $\pi^n \xrightarrow[n \rightarrow +\infty]{d} \pi$  and if  $\frac{q_j(n^2s)}{n} \xrightarrow[n \rightarrow +\infty]{} L(s)$ , then, considering  $R$  the cdf of  $\pi$ , we have the expression

$$\begin{aligned} - \int_0^t \int_0^{+\infty} R(s, x) \partial_t \varphi(s, x) dx ds - \frac{T}{\gamma} \int_0^t \int_0^{+\infty} R(s, x) \partial_x^2 \varphi(s, x) dx ds \\ = \frac{P}{\gamma} \int_0^t \varphi(s, L(s)) ds + \int_0^{+\infty} R_0(x) \varphi(0, x) dx. \end{aligned}$$

where  $\varphi = -\partial_x \psi$ .

# Identification of the limit

It is the weak formulation of the equation

$$\begin{aligned}\partial_t R &= \frac{T}{\gamma} \partial_x^2 R & t \in \mathbb{R}_+, x \in [0, L(t)] \\ R(t, 0) &= 0 & t \in \mathbb{R}_+ \\ \partial_x R(t, L(t)) &= \frac{P}{T} & t \in \mathbb{R}_+ \\ R(0, x) &= R_0(x) & x \in [0, L_0] \\ L'(t) &= -\frac{T^2}{\gamma P} \partial_x^2 R(t, L(t)) & t \geq 0\end{aligned}$$

This equation admits a unique weak solution  $(R, L)$ .

# The border

What is new in this system ? The border.

Problem : we don't know how to show that the sequence  $\left(\frac{q_n(n^2 \cdot)}{n}\right)_{n \in \mathbb{N}}$  is relatively compact.

Solution ? Setting  $\beta = \delta_{L(t)}(dx)dt$ , the weak formulation becomes

$$\begin{aligned} - \int_0^t \int_0^{+\infty} R(s, x) \partial_t \varphi(s, x) dx ds - \frac{T}{\gamma} \int_0^t \int_0^{+\infty} R(s, x) \partial_x^2 \varphi(s, x) dx ds \\ = \frac{P}{\gamma} \int_{[0, t] \times \mathbb{R}_+} \varphi d\beta + \int_0^{+\infty} R_0(x) \varphi(0, x) dx. \end{aligned}$$

and it is easier to show that the sequence  $(\beta^n)_n$ , defined by

$$\beta^n(dx, dt) = \delta_{\frac{q_n(n^2 t)}{n}}(dx)dt$$

is relatively compact.

# Relative Entropy

We start at an initial configuration  $\nu_0^n$  and we note  $\nu_t^n$  the distribution of the system at time  $t$ . We assume that  $\nu_t^n$  is absolutely continuous with respect to  $\mu_{P,T}^n$ , with density  $f_n(t, \cdot)$ . We define

$$H_n(t) = H(\nu_t^n | \mu_{\beta,P}^n) = \int_{\Gamma_n} f_n(t, q, p) \log(f_n(t, q, p)) d\mu_{P,T}^n.$$

## Proposition

$H_n$  is a non-increasing function.

## Theorem

For all functions  $\varphi$  such that the quantities are well-defined, if  $X$  is a random variable with law  $\nu$ , we have

$$\forall \alpha > 0, \quad \mathbb{E}[\varphi(X)] \leq \frac{1}{\alpha} \left( H(\nu|\mu) + \log \left( \mathbb{E}_\mu [\exp(\alpha\varphi)] \right) \right)$$

Application :

$$\forall \alpha > 0, \quad \mathbb{E}[\varphi(q(t), p(t))] \leq \frac{1}{\alpha} \left( H_n(t) + \log \mathbb{E}_{\mu_{P,T}^n} [\exp(\alpha\varphi(q, p))] \right)$$

Since the entropy is non-increasing, we can write :

$$\forall \alpha > 0, \quad \mathbb{E}[\varphi(q(t), p(t))] \leq \frac{1}{\alpha} \left( H_n(0) + \log \mathbb{E}_{\mu_{P,T}^n} [\exp(\alpha\varphi(q, p))] \right)$$

## Application : tightness of $(\beta^n)_n$

We apply the inequality with  $\alpha = cn$ .

$$\mathbb{E} \left[ \frac{q_n(n^2 t)}{n} \right] \leq \frac{1}{cn} H_n(0) + \frac{1}{cn} \log \left( \mathbb{E}_{\mu_{P,T}} [e^{cq_n}] \right).$$

and for  $c \leq P/T$ ,

$$\frac{1}{cn} \log \mathbb{E}_{\mu_{P,T}^n} [e^{cq_n}] = \frac{1}{c} \log \left( \frac{P}{P - cT} \right).$$

Then, if  $H_n(0) = O(n)$ , we obtain that this quantity is bounded by a constant  $C > 0$ .

# Application : tightness of $(\beta^n)_n$

It implies that

$$\sup_{n \in \mathbb{N}} \sup_{s \in \mathbb{R}_+} \mathbb{P} \left( \frac{q_n(n^2 s)}{n} \geq r \right) \leq \frac{C}{r}$$

and then by Markov inequality, we can show that for every  $\varepsilon > 0$ , it exists  $r > 0$  such that

$$\sup_{n \in \mathbb{N}} \mathbb{P}(\beta^n([0, t] \times [r, +\infty]) > \varepsilon) \leq \varepsilon.$$

It is sufficient to use Prokhorov theorem and show that the sequence of the laws of  $\beta^n$  is relatively compact.

Thank you for your attention

