

# A kinetic Fisher-KPP equation : traveling waves and front acceleration.

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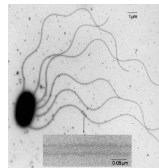
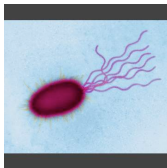
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- 1 Introduction.
- 2 Traveling waves when  $V$  is bounded
- 3 Finally, what happens when  $V$  is unbounded ?
- 4 Conclusion.

# How do bacteria move ...

The bacteria *E. Coli*  
moves thanks to  
flagella :

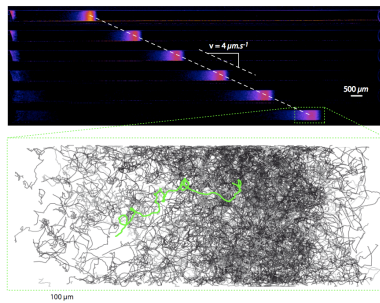


and with a so-called *run and tumble*  
process :  
straight swimming for 1s  
and  
change of direction for 0.1s.

From Howard Berg's lab

# Collective migration

Bacterial traveling pulses :



**The kinetic point of view is the most relevant** for this situation.



Saragosti J, Calvez V, Bournaveas N, Buguin A, Silberzan P, et al. (2010) Mathematical Description of Bacterial Traveling Pulses. PLoS Comput Biol 6(8)

# Kinetic equations

- Probability density of bacteria  $f(t, x, v)$  at time  $t$ , position  $x$  and speed  $v$ .  
Total density  $\rho := \int_V f(v) dv$ .
- The velocity set  $V$  : symmetric, **bounded** or **unbounded**.

The model ( e.g. Schmeiser and al, 2010 ) :

$$\underbrace{\varepsilon \partial_t f + v \partial_x f}_{\text{Free run}} = \underbrace{\frac{1}{\varepsilon} (M(v)\rho - f)}_{\text{Tumbling}} + \underbrace{\varepsilon r \rho (M(v) - f)}_{\text{Growth with saturation}} \quad (1)$$

where the Maxwellian  $M$  on the space  $V$  satisfies

$$\int_V M(v) dv = 1, \quad \int_V v M(v) dv = 0, \quad \int_V v^2 M(v) dv = D. \quad (2)$$

# Aim of the talk : traveling waves

**Formally**, in the limit  $\varepsilon \rightarrow 0$ ,  $\rho$  satisfies the classical Fisher-KPP equation (1937) :

$$\partial_t \rho - D \partial_{xx} \rho = r \rho (1 - \rho). \quad (3)$$

It is now well known that (3) admits traveling waves solutions with minimal speed  $c_{KPP} = 2\sqrt{rD}$ .

**Aim of the talk** : Discuss **existence** and **stability** of traveling fronts for (2), for all values of  $\varepsilon$  ( i.e. in the **full kinetic** regime ).

# Plan.

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Without loss of generality, we assume that  $V = [-1; 1]$ .

### Definition

We say that a function  $f(t, x, v)$  is a **traveling front solution** of **speed**  $c \in \mathbb{R}^+$  of equation (1) if it can be written  $f(t, x, v) = \mu(\xi = x - ct, v)$ , where **the profile**  $\mu \in \mathcal{C}^2(\mathbb{R} \times V)$  is nonnegative, satisfies  $\mu(-\infty, \cdot) = M$ ,  $\mu(+\infty, \cdot) = 0$ , and  $\mu$  solves

$$\varepsilon(v - c\varepsilon)\partial_\xi \mu = (M(v)\nu - \mu) + r\varepsilon^2\nu(M(v) - \mu), \quad \xi \in \mathbb{R}, v \in V. \quad (4)$$

where  $\nu$  is the macroscopic density associated to  $\mu$ , that is  $\nu(\xi) = \int_V \mu(\xi, v) dv$ .

# Existence result

## Theorem

Assume that  $\varepsilon > 0$  and that  $\text{Supp}(M) = [-1, 1]$ . There exists a **minimal** speed  $c^*(\varepsilon) \in (0, \frac{1}{\varepsilon})$  such that **there exists a traveling wave solution of (1) of speed  $c$  for  $c \in [c^*(\varepsilon), \frac{1}{\varepsilon}]$** . Moreover, this traveling wave is **nonincreasing** with respect to  $x$ .

This is **not** a perturbative result from the Fisher-KPP equation.

# Finding the speed : Dispersion relation

- As for the Fisher-KPP equation, the rate of **exponential decay** in space and the speed are given by the linearized problem at the edge of the front. In this case, we have  $D\lambda^2 - c\lambda + r = 0$ .
- Similarly, looking for a solution of the linearized problem of type  $e^{-\lambda\xi}Q(v)$  yields the following **spectral problem** : For all  $\lambda$ , find  $c(\lambda)$  such that there exists a Maxwellian  $Q_\lambda$  such that

$$\forall v \in V, \quad (1 + \varepsilon\lambda(c(\lambda)\varepsilon - v)) Q_\lambda(v) = (1 + r\varepsilon^2) \int_V M(v) Q_\lambda(v) dv. \quad (5)$$

## Proposition

For all  $\varepsilon > 0$ , the minimal speed  $c^*(\varepsilon)$  is given by  $c^*(\varepsilon) = \min_{\lambda > 0} c(\lambda)$  where  $c(\lambda)$  is for all  $\lambda$  a solution of the following **dispersion relation** :

$$(1 + r\varepsilon^2) \int_V \frac{M(v)}{1 + \varepsilon\lambda(c(\lambda)\varepsilon - v)} dv = 1. \quad (6)$$

This relation is **not** compatible with and unbounded velocity set.

# Remarks on the results

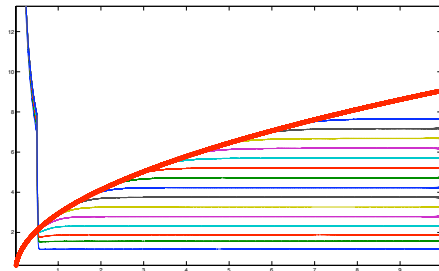
- ① The existence result is proved using a sub- and super-solutions technique (see e.g. Berestycki and Hamel).
- ② We can provide estimates for the critical speed  $c^*(\varepsilon)$ . In particular,

## Proposition

Assume that  $\text{Supp}(M) = [-1, 1]$ , then  $c^*(\varepsilon) \xrightarrow{\varepsilon \rightarrow 0} 2\sqrt{rD} := c_{\text{KPP}}$ .

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# Numerical simulations



**Figure:** Evolution of the speed of the front for different values of the maximal speed. The Maxwellian here is a Gaussian :  $M(v) = C(V_{max}) \exp\left(-\frac{v^2}{2}\right)$ . "Bell" initial condition.

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## Concluding remarks.

- When  $V$  is bounded we construct traveling wave profiles. The speed is obtained through a dispersion relation.
- On the contrary, when  $V$  is unbounded, acceleration pattern is observed.

Thank you for your attention !



# References

-  E. Bouin, V. Calvez, G. Nadin, *A kinetic KPP equation : traveling waves and front acceleration*, in preparation,
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-  D.G. Aronson and H.F. Weinberger, *Multidimensional nonlinear diffusion arising in population genetics*, Adv. Math. 30, pp. 33-76, (1978).