## INTRODUCTION TO STATISTICAL MECHANICS

## **General framework** (Static setting)

Dauphine, April 5, 2025

Understand the large scale behavior of a physics system whose interactions are described on the microscopic level

Start from a model, whose general framework is the following.

Structure of the physics system is represented by a graph G = (V, E), finite.



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• Set of configurations on the graph G: C(G),

- vertex configurations,
- edge configurations.

Parameters:

- intensity of interactions between microscopic components
- external temperature.
- $\Rightarrow$  Positive weight function  $w = (w_e)_{e \in E}$  on the edges.



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- To a configuration C, one assigns an energy  $\mathcal{E}_w(C)$ .
- Boltzmann measure on configurations:

$$\forall \mathbf{C} \in \mathcal{C}(\mathbf{G}), \quad \mathbb{P}(\mathbf{C}) = \frac{e^{-\mathcal{E}_w(\mathbf{C})}}{Z(\mathbf{G}, w)},$$

where  $Z(\mathbf{G}, w) = \sum_{\mathbf{C} \in \mathcal{C}(\mathbf{G})} e^{-\mathcal{E}_w(\mathbf{C})}$  is the partition function.

Understand the behavior of configurations when the graph is large (infinite).

#### Model of ferromagnetism - mixture of two materials





Wilhelm Lenz (1888-1957)

Ernst Ising (1900-1998)

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- Graph G = (V, E).
- A spin configuration σ assigns a spin σ<sub>x</sub> ∈ {−1,1} to a each vertex x of the graph G.

 $\Rightarrow C(G) = \{-1, 1\}^{V} = \text{set of spin configurations.}$ 

► A spin configuration

#### ► A spin configuration / two interpretations

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#### Magnetic moments:

+1/ $\rightarrow$ , -1/ $\leftarrow$ 



► A spin configuration / two interpretations

# Magnetic moments: $+1/\rightarrow, -1/\leftarrow$



Mixture of two materials: +1/0, -1/0.



## THE ISING MODEL

- ▶ Positive weight function: coupling constants  $J = (J_e)_{e \in E}$ .
- Energy of a spin configuration:  $\mathcal{E}_J(\sigma) = -\sum_{e=xy\in E} J_{xy}\sigma_x\sigma_y$ .
- Ising Boltzmann measure:

$$\forall \sigma \in \{-1,1\}^{\vee}, \quad \mathbb{P}_{\text{Ising}}(\sigma) = \frac{e^{-\mathcal{E}_J(\sigma)}}{Z_{\text{Ising}}(\mathsf{G},J)}.$$

- Two neighboring spins  $\sigma_x, \sigma_y$  tend to align.
- The higher the coupling  $J_{xy}$ , the higher this tendency.

#### Adsorption of di-atomic molecules on the surface of a cristal





Sir Ralph H. Fowler (1889-1944) Congrès Solvay 1927.

George S. Rushbrooke (1915-1995)

- Graph G = (V, E).
- A dimer configuration or perfect matching: subset of edges such that each vertex touches exactly one edge of this subset.

 $\Rightarrow C(G) = M(G) = \text{set of dimer configurations.}$ 

► A dimer configuration.



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► A dimer configuration.





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► A dimer configuration.



- Positive weight function:  $v = (v_e)_{e \in E}$ .
- Energy of a configuration M:  $\mathcal{E}_{\nu}(M) = -\sum_{e \in M} \log \nu_e$ .
- Dimer Botzmann measure:

$$\forall M \in \mathcal{M}(G), \quad \mathbb{P}_{dimer}(M) = \frac{\prod_{e \in M} \nu_e}{Z_{dimer}(G, \nu)}.$$

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Edges with higher weights are more likely to occur.

#### Related to electrical networks



Gustav Kirchhoff (1824-1887)

- Graph G = (V, E).
- A spanning tree: subset of edges covering all vertices of the graph, connected, with no cycle.

 $\Rightarrow C(G) = T(G) = set of spanning trees.$ 

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► A spanning tree



► A spanning tree



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► A spanning tree



- Positive weight function:  $\rho = (\rho_e)_{e \in E}$ .
- Energy of a tree T:  $\mathcal{E}_{\rho}(T) = -\sum_{e \in T} \log \rho_e$ .
- ► Tree Boltzmann measure:

$$\forall \mathsf{T} \in \mathfrak{T}(\mathsf{G}), \quad \mathbb{P}_{\text{tree}}(\mathsf{T}) = \frac{\prod_{\mathsf{e} \in \mathsf{T}} \rho_{\mathsf{e}}}{Z_{\text{tree}}(\mathsf{G}, \rho)}.$$

Edges with higher weights are more likely to occur.

## PERCOLATION

Flow of a liquid through a porous material





Simon Broadbent (1928-2002)

John Hammersley (1920-2004)

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- Graph G = (V, E).
- ► Configuration of opened and closed edges:  $\forall e \in E, \omega_e \in \{0, 1\}$ .

$$\Rightarrow \quad \mathcal{C}(\mathsf{G}) = \{0,1\}^{\mathsf{E}}.$$

# Percolation

► A percolation configuration



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## Percolation

► A percolation configuration



- Let  $p \in [0,1]$ . Each edge is opened/closed with probability p / 1-p, independently.
- Percolation measure:

$$\forall \, \omega \in \{0,1\}^{\mathsf{E}}, \quad \mathbb{P}_{\mathsf{perco}}(\omega) = p^{\sum_{\mathsf{e} \in \mathsf{E}} \omega_{\mathsf{e}}} (1-p)^{|\mathsf{E}| - \sum_{\mathsf{e} \in \mathsf{E}} \omega_{\mathsf{e}}}.$$

- ▶ The higher *p* is, the more open edges there are
- ▶ For which values of *p* do we percolate ?

Let the edge-length tend to 0 Look at a "typical" configuration.

► Ising model (Illustrations of R. Cerf)



J small

J critical



- On  $\mathbb{Z}^2$ :  $J_c \equiv \frac{1}{2} \log(1 + \sqrt{2})$  [Kramers et Wannier]
- ▶ Phase transition: studied through magnetization.

Dimer model (Illustration of R. Kenyon)



- One sees two phases on the same figure.
- Phase transition studied through decay of correlations.

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Percolation (Illustration of Erzbischof)



p = 0.51

- On  $\mathbb{Z}^2$ :  $p_c = 0.5$  [Kesten]
- Phase transition: (non) existence of an infinite connected component.

- Identification of the phase transition.
- Understand the sub/super critical models.
- Understand the critical model (at the phase transition):
  - Universality and conformal invariance.
  - Conjectures: Nienhuis, Cardy, Duplantier ...
    Proofs: Lawler, Schramm, Werner (Fields 2006), D. Chelkak, S. Smirnov (Fields 2010), H. Duminil-Copin (Fields 2022), ...

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▶ Ising model on G.



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▶ Low temperature expansion [Kramers-Wannier].



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Polygon contour configurations on G<sup>\*</sup>.



► Fisher's correspondence: exactly keep the polygon contour edges.



Fill the decorations:  $2^{|V^*|}$  possibilities.



#### EXACTLY SOLVABLE MODELS

One of the tools to study the macroscopic behavior is the partition function:

$$Z(\mathsf{G},w) = \sum_{\mathsf{C}\in\mathcal{C}(\mathsf{G})} e^{-\mathcal{E}_w(\mathsf{C})},$$

the normalizing constant in the Boltzmann measure.

$$\forall C \in \mathcal{C}(G), \quad \mathbb{P}(C) = \frac{e^{-\mathcal{E}_w(C)}}{Z(G, w)}.$$

- The model is exactly solvable if there exists an exact, explicit formula for the partition function
- Three exactly solvable models:
  - ▶ Ising-2d: Onsager (1944) Fisher (1966).
  - Dimers-2d: Kasteleyn, Temperley-Fisher (1961).
  - Spanning trees: Kirchhoff (1848).

• Let  $M_1$ ,  $M_2$  be two dimer configurations of G, and  $M_1 \cup M_2$  be their superposition.



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► M<sub>1</sub> ∪ M<sub>2</sub> is a disjoint union of alternating cycles, where an alternating cycles has edges alternating between M<sub>1</sub> and M<sub>2</sub>. Alternating cycles of length 2 are called doubled edges.