

# INTRODUCTION TO STATISTICAL MECHANICS

## GENERAL FRAMEWORK (STATIC SETTING)

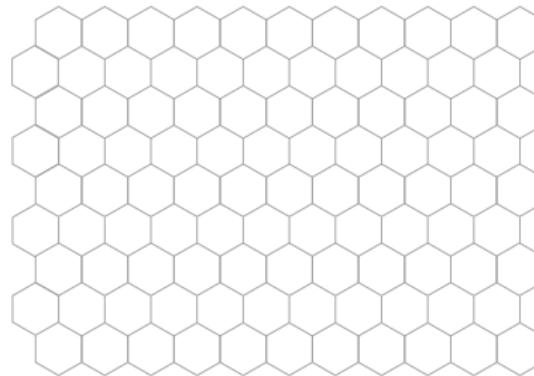
January 23, 2026

# STATISTICAL MECHANICS

*Understand the large scale behavior of a physics system whose interactions are described on the microscopic level*

Start from a **model**, whose general framework is the following.

- ▶ Structure of the physics system is represented by a **graph**  $G = (V, E)$ , finite.

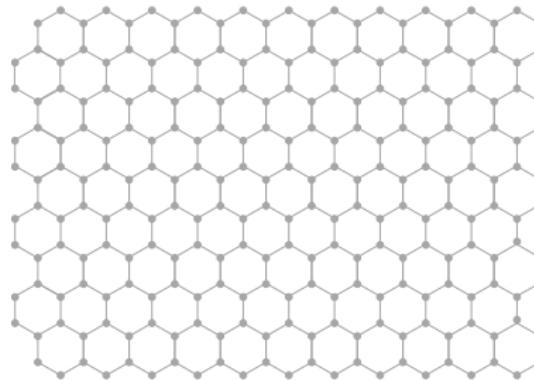


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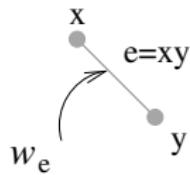
- ▶ Structure of the physics system is represented by a **graph**  $G = (V, E)$ , finite.



# STATISTICAL MECHANICS

- ▶ Set of configurations on the graph  $G$ :  $\mathcal{C}(G)$ ,
  - ▶ vertex configurations,
  - ▶ edge configurations.
- ▶ Parameters:
  - ▶ intensity of interactions between microscopic components
  - ▶ external temperature.

⇒ Positive weight function  $w = (w_e)_{e \in E}$  on the edges.



# STATISTICAL MECHANICS

- ▶ To a configuration  $C$ , one assigns an **energy**  $\mathcal{E}_w(C)$ .
- ▶ **Boltzmann measure** on configurations:

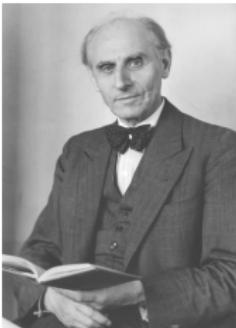
$$\forall C \in \mathcal{C}(G), \quad \mathbb{P}(C) = \frac{e^{-\mathcal{E}_w(C)}}{Z(G, w)},$$

where  $Z(G, w) = \sum_{C \in \mathcal{C}(G)} e^{-\mathcal{E}_w(C)}$  is the **partition function**.

*Understand the behavior of configurations  
when the graph is large (infinite).*

# THE ISING MODEL

*Model of ferromagnetism - mixture of two materials*



Wilhelm Lenz (1888-1957)

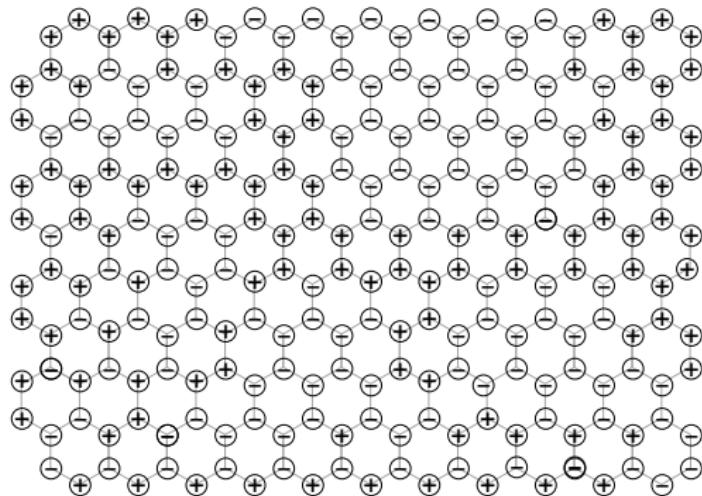


Ernst Ising (1900-1998)

- ▶ Graph  $G = (V, E)$ .
- ▶ A **spin configuration**  $\sigma$  assigns a spin  $\sigma_x \in \{-1, 1\}$  to each vertex  $x$  of the graph  $G$ .  
 $\Rightarrow \mathcal{C}(G) = \{-1, 1\}^V = \text{set of spin configurations.}$

# THE ISING MODEL

- ▶ A spin configuration

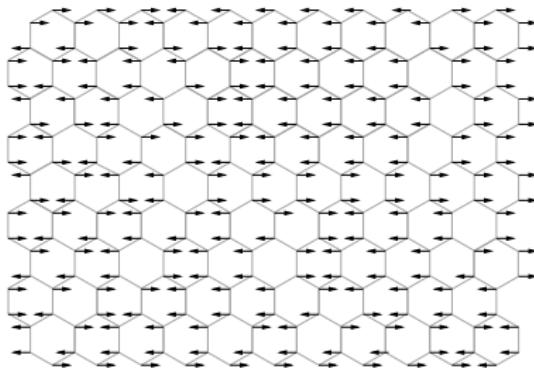


# THE ISING MODEL

- ▶ A spin configuration / two interpretations

Magnetic moments:

$+1/\rightarrow, -1/\leftarrow$

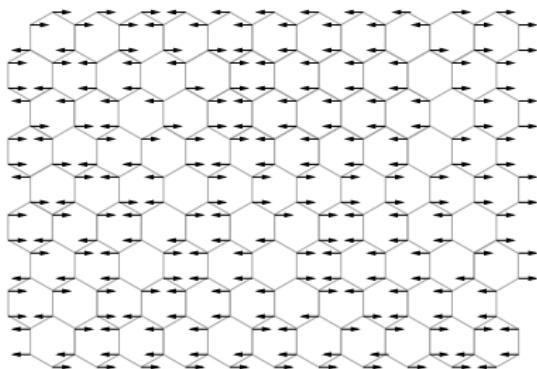


# THE ISING MODEL

- ▶ A spin configuration / two interpretations

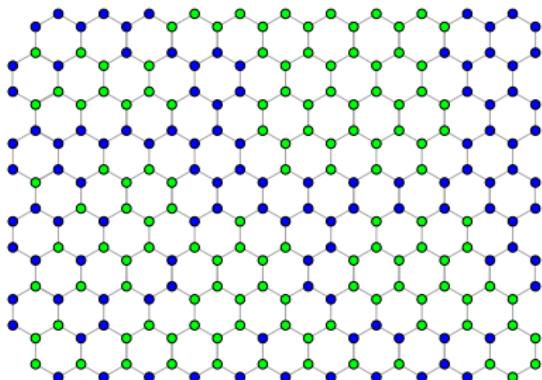
Magnetic moments:

$+1 \rightarrow, -1 \leftarrow$



Mixture of two materials:

$+1/\bullet, -1/\circ$ .



# THE ISING MODEL

- ▶ Positive weight function: coupling constants  $J = (J_e)_{e \in E}$ .
- ▶ Energy of a spin configuration:  $\mathcal{E}_J(\sigma) = - \sum_{e=xy \in E} J_{xy} \sigma_x \sigma_y$ .
- ▶ Ising Boltzmann measure:

$$\forall \sigma \in \{-1, 1\}^V, \quad \mathbb{P}_{\text{Ising}}(\sigma) = \frac{e^{-\mathcal{E}_J(\sigma)}}{Z_{\text{Ising}}(G, J)}.$$

- ▶ Two neighboring spins  $\sigma_x, \sigma_y$  tend to align.
- ▶ The higher the coupling  $J_{xy}$ , the higher this tendency.

# THE DIMER MODEL

*Adsorption of di-atomic molecules on the surface of a cristal*



Sir Ralph H. Fowler (1889-1944)  
Congrès Solvay 1927.

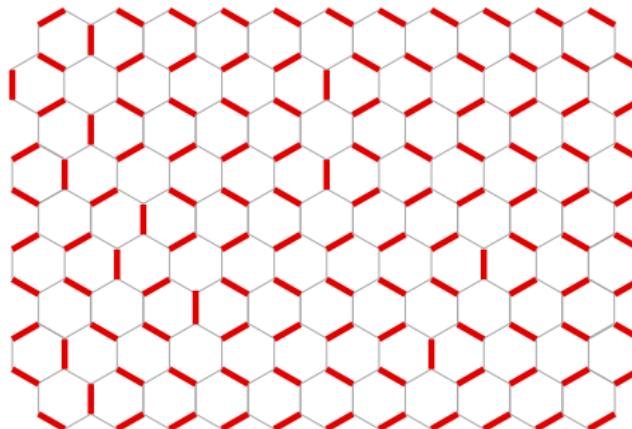


George S. Rushbrooke (1915-1995)

- ▶ Graph  $G = (V, E)$ .
- ▶ A **dimer configuration** or **perfect matching**: subset of edges such that each vertex touches exactly one edge of this subset.  
 $\Rightarrow \mathcal{C}(G) = \mathcal{M}(G) = \text{set of dimer configurations.}$

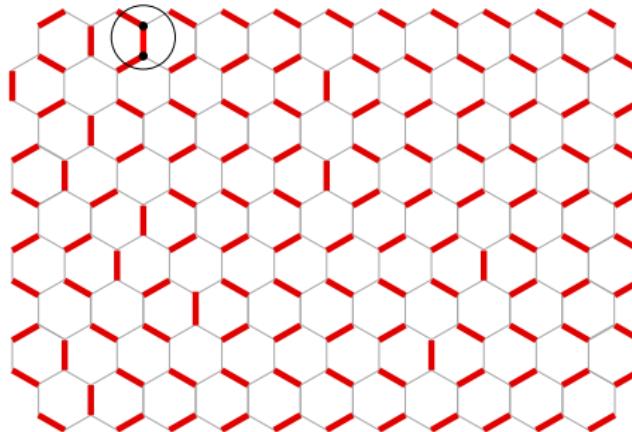
## THE DIMER MODEL

- ▶ A dimer configuration.



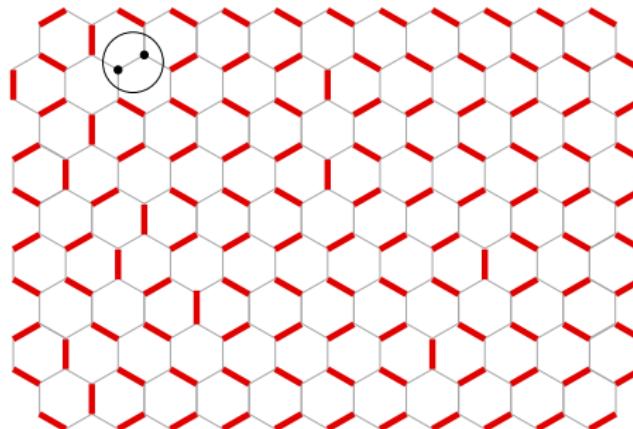
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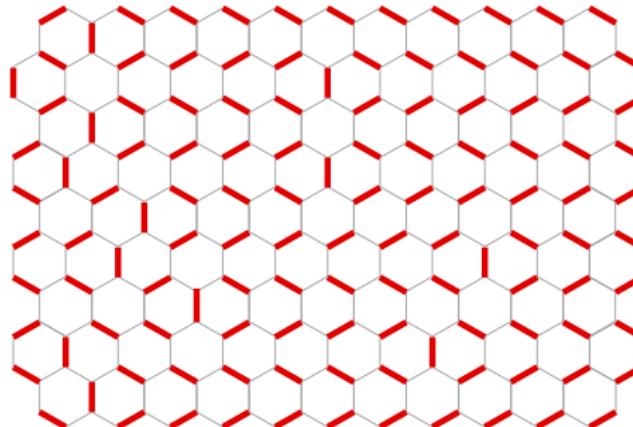
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- ▶ Positive weight function:  $\nu = (\nu_e)_{e \in E}$ .
- ▶ Energy of a configuration  $M$ :  $\mathcal{E}_\nu(M) = -\sum_{e \in M} \log \nu_e$ .
- ▶ Dimer Botzmann measure:

$$\forall M \in \mathcal{M}(G), \quad \mathbb{P}_{\text{dimer}}(M) = \frac{\prod_{e \in M} \nu_e}{Z_{\text{dimer}}(G, \nu)}.$$

- ▶ Edges with higher weights are more likely to occur.

# SPANNING TREES

*Related to electrical networks*

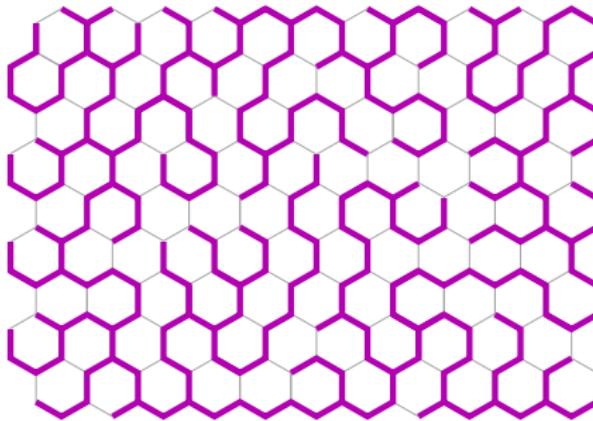


Gustav Kirchhoff (1824-1887)

- ▶ Graph  $G = (V, E)$ .
- ▶ A **spanning tree**: subset of edges covering all vertices of the graph, connected, with no cycle.  
 $\Rightarrow \mathcal{C}(G) = \mathcal{T}(G) = \text{set of spanning trees.}$

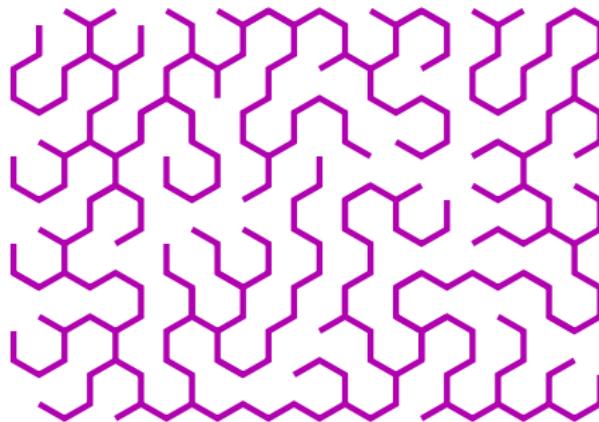
## SPANNING TREES

- ▶ A spanning tree



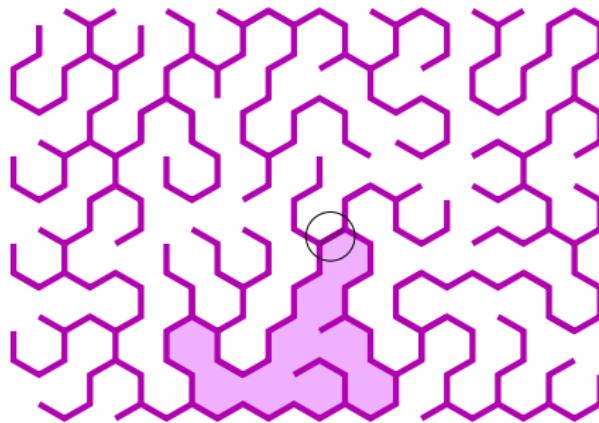
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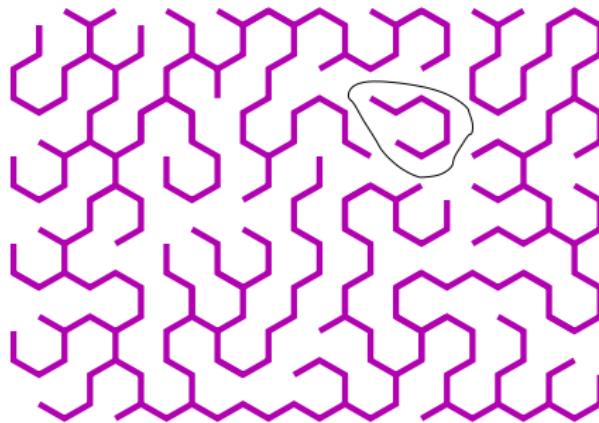
## SPANNING TREES

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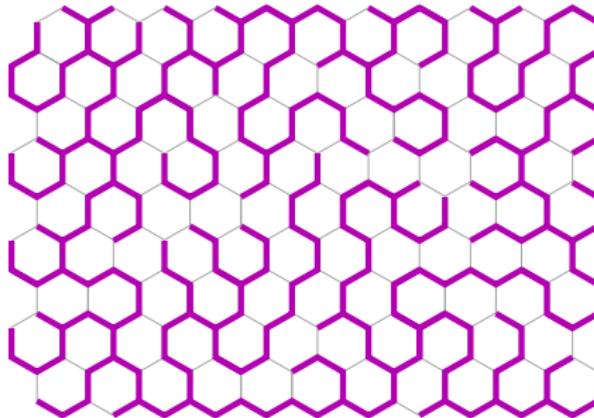
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- ▶ Positive weight function:  $\rho = (\rho_e)_{e \in E}$ .
- ▶ Energy of a tree  $T$ :  $\mathcal{E}_\rho(T) = - \sum_{e \in T} \log \rho_e$ .
- ▶ Tree Boltzmann measure:

$$\forall T \in \mathcal{T}(G), \quad \mathbb{P}_{\text{tree}}(T) = \frac{\prod_{e \in T} \rho_e}{Z_{\text{tree}}(G, \rho)}.$$

- ▶ Edges with higher weights are more likely to occur.

# PERCOLATION

*Flow of a liquid through a porous material*



Simon Broadbent (1928-2002)



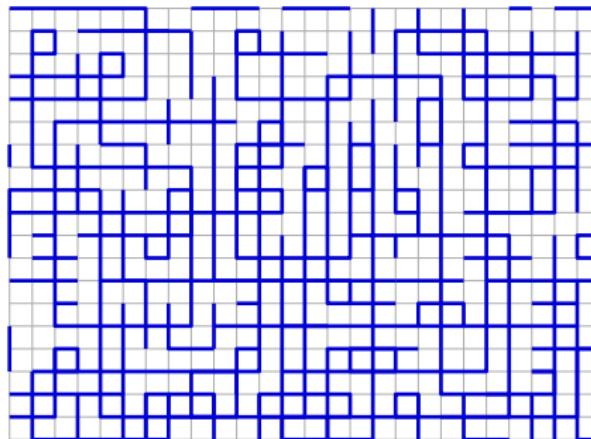
John Hammersley (1920-2004)

- ▶ Graph  $G = (V, E)$ .
- ▶ Configuration of opened and closed edges:  $\forall e \in E, \omega_e \in \{0, 1\}$ .

$$\Rightarrow \mathcal{C}(G) = \{0, 1\}^E.$$

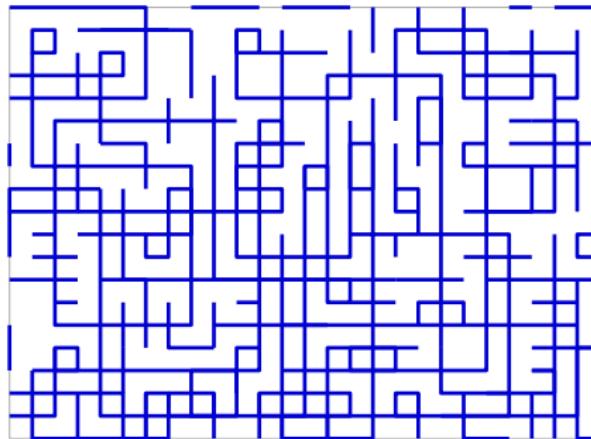
# PERCOLATION

- ▶ A percolation configuration



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- ▶ Let  $p \in [0, 1]$ . Each edge is opened/closed with probability  $p$  /  $1 - p$ , independently.
- ▶ **Percolation measure:**

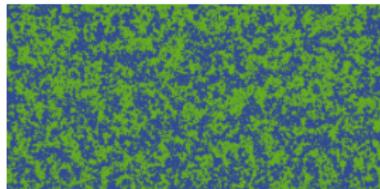
$$\forall \omega \in \{0,1\}^E, \quad \mathbb{P}_{\text{perco}}(\omega) = p^{\sum_{e \in E} \omega_e} (1-p)^{|E| - \sum_{e \in E} \omega_e}.$$

- ▶ The higher  $p$  is, the more open edges there are
- ▶ For which values of  $p$  do we percolate ?

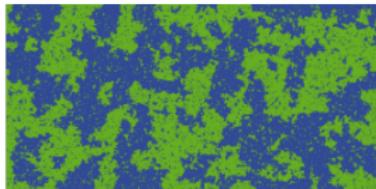
# MACROSCOPIC BEHAVIOR

*Let the edge-length tend to 0  
Look at a “typical” configuration.*

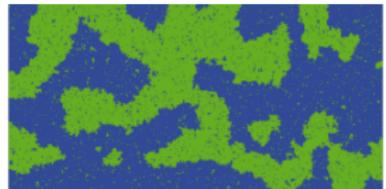
- ▶ Ising model (Illustrations of R. Cerf)



$J$  small



$J$  critical

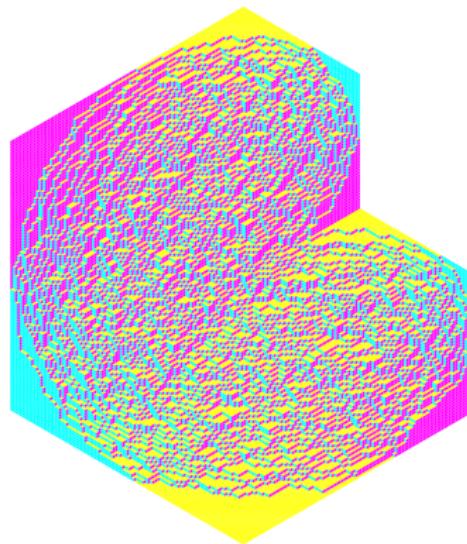


$J$  large

- ▶ On  $\mathbb{Z}^2$ :  $J_c \equiv \frac{1}{2} \log(1 + \sqrt{2})$  [Kramers et Wannier]
- ▶ Phase transition: studied through magnetization.

## MACROSCOPIC BEHAVIOR

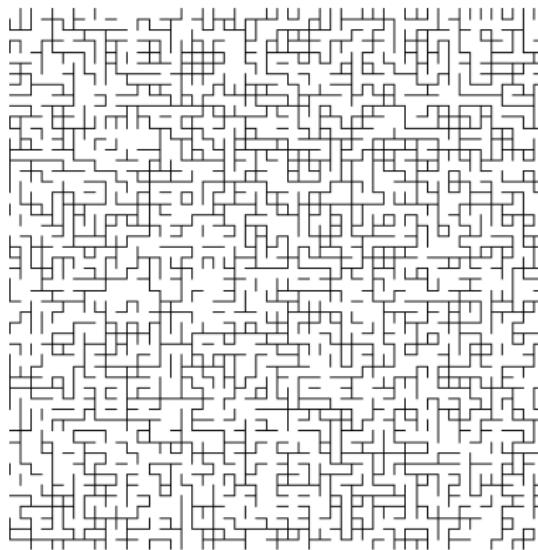
- Dimer model (Illustration of R. Kenyon)



- One sees two phases on the same figure.
- Phase transition studied through decay of correlations.

# MACROSCOPIC BEHAVIOR

- ▶ Percolation (Illustration of Erzbischof)



$$p = 0.51$$

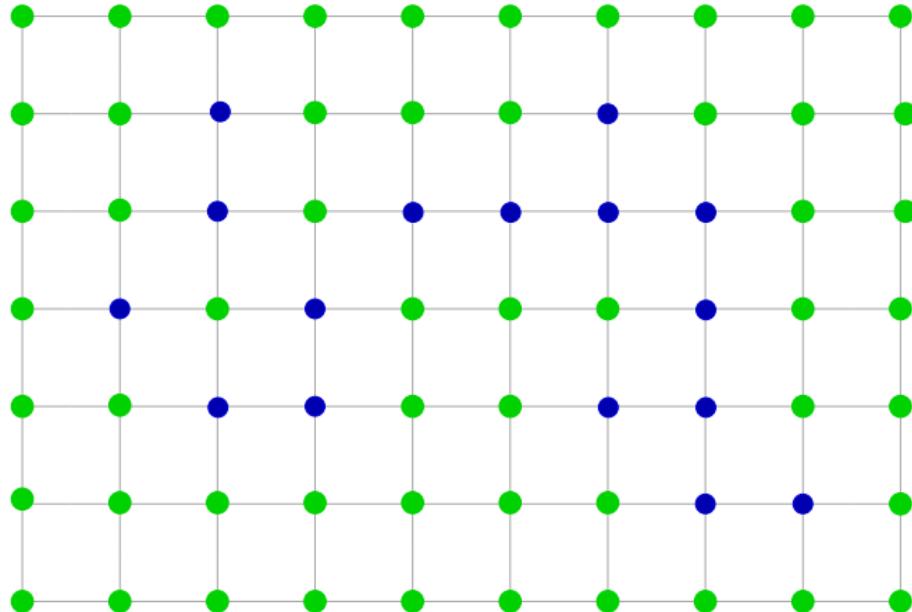
- ▶ On  $\mathbb{Z}^2$ :  $p_c = 0.5$  [Kesten]
- ▶ Phase transition: (non) existence of an infinite connected component.

## MACROSCOPIC BEHAVIOR

- ▶ Identification of the phase transition.
- ▶ Understand the sub/super critical models.
- ▶ Understand the critical model (at the phase transition):
  - ▶ Universality and conformal invariance.
  - ▶ Conjectures: Nienhuis, Cardy, Duplantier ...  
Proofs: Lawler, Schramm, Werner (Fields 2006), D. Chelkak, S. Smirnov (Fields 2010), H. Duminil-Copin (Fields 2022), ...

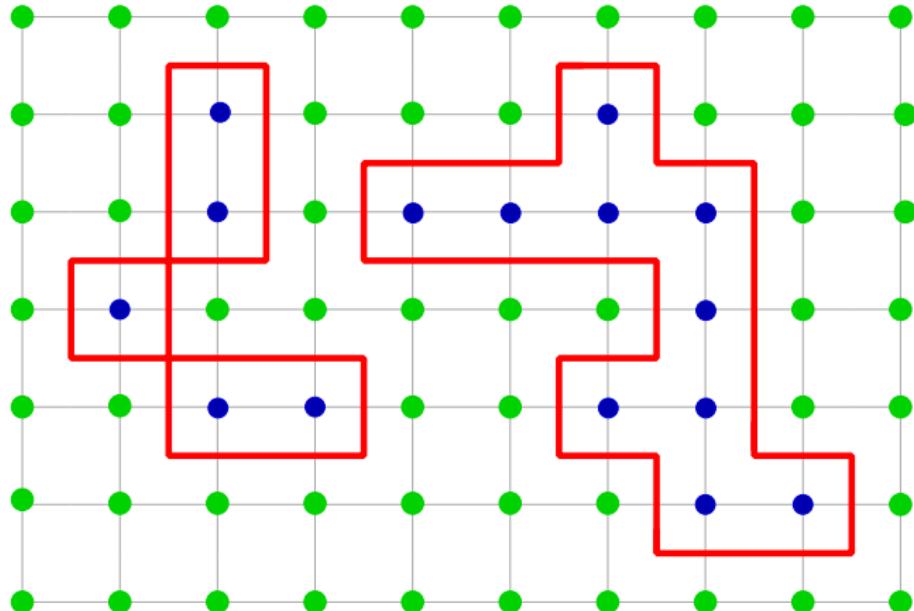
## 2. RELATIONS BETWEEN MODELS: ISING - DIMERS [FISHER'66]

- ▶ Ising model on  $G$ .



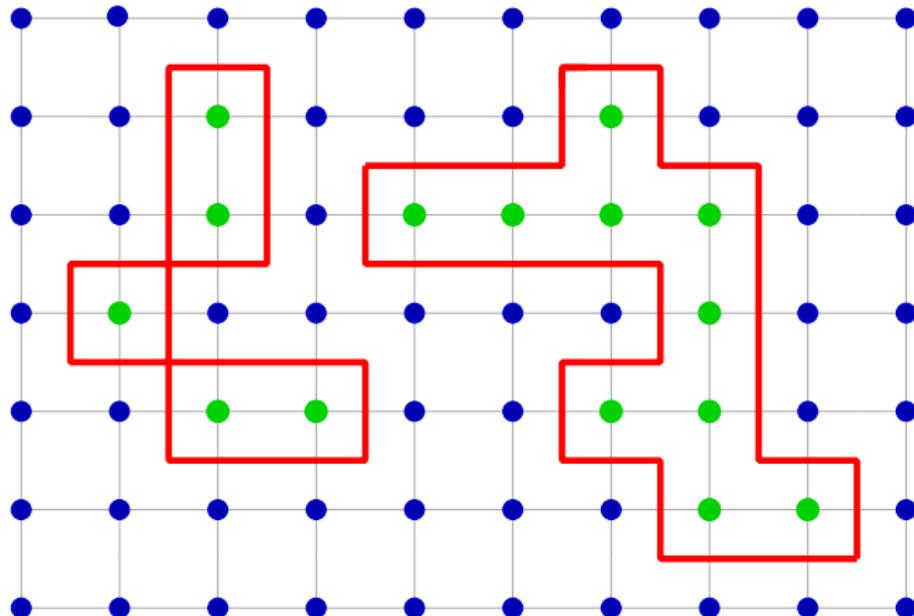
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- ▶ Low temperature expansion [Kramers-Wannier].



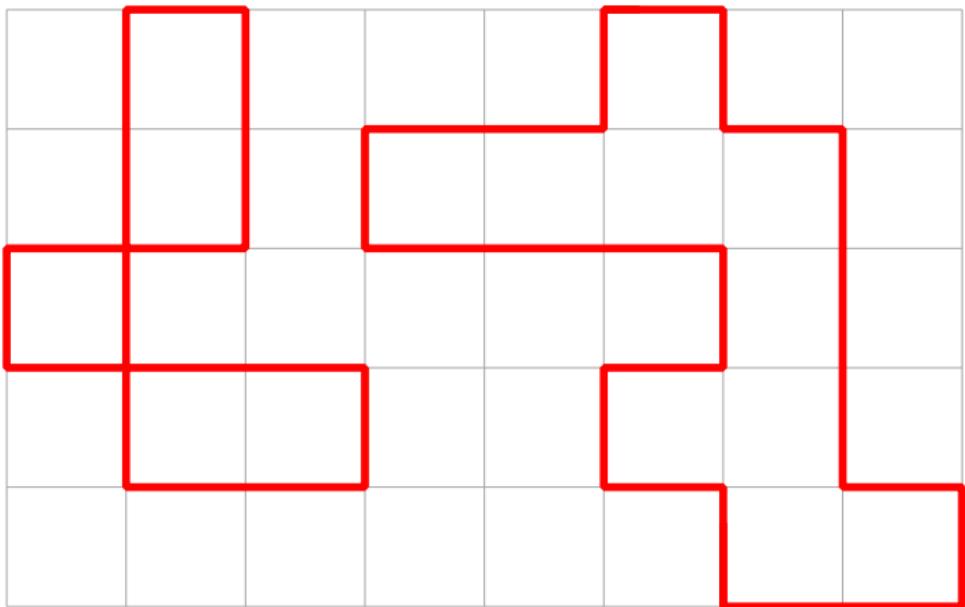
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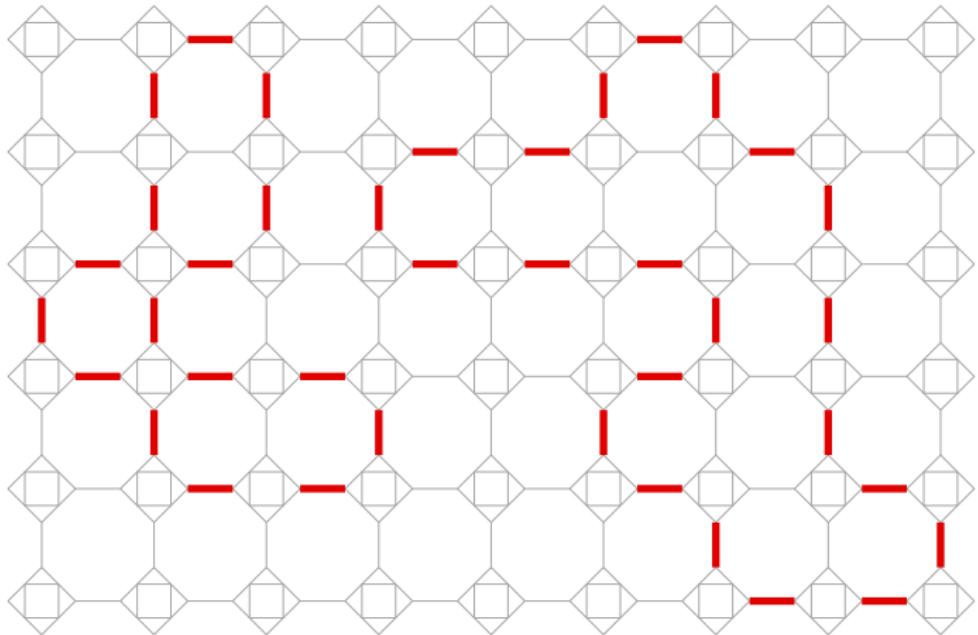
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- ▶ Polygon contour configurations on  $G^*$ .



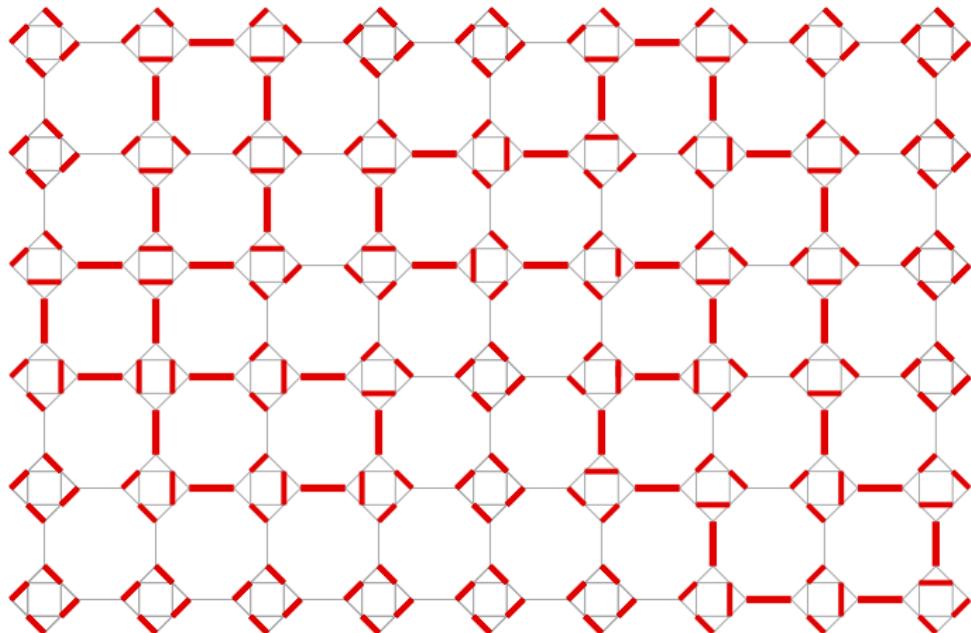
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- ▶ Fisher's correspondence: exactly keep the polygon contour edges.



## 2. RELATIONS BETWEEN MODELS: ISING - DIMERS [FISHER'66]

- ▶ Fill the decorations:  $2^{|V^*|}$  possibilities.



## EXACTLY SOLVABLE MODELS

- One of the tools to study the macroscopic behavior is the **partition function**:

$$Z(G, w) = \sum_{C \in \mathcal{C}(G)} e^{-\mathcal{E}_w(C)},$$

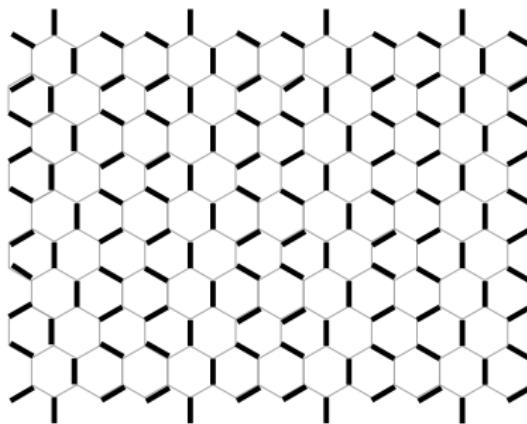
the normalizing constant in the Boltzmann measure.

$$\forall C \in \mathcal{C}(G), \quad \mathbb{P}(C) = \frac{e^{-\mathcal{E}_w(C)}}{Z(G, w)}.$$

- The model is **exactly solvable** if there exists an exact, explicit formula for the partition function
- Three exactly solvable models:
  - Ising-2d: Onsager (1944) - Fisher (1966).
  - Dimers-2d: Kasteleyn, Temperley-Fisher (1961).
  - Spanning trees: Kirchhoff (1848).

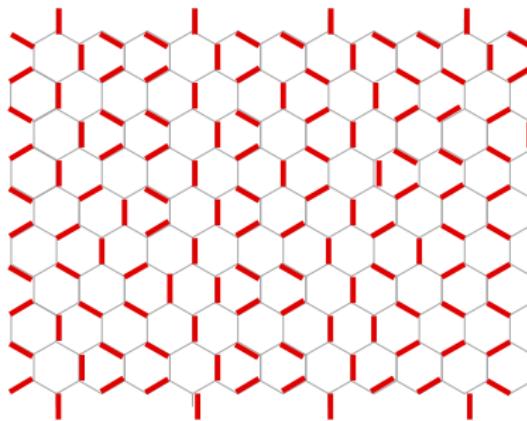
## PRELIMINARIES

- ▶ Let  $M_1, M_2$  be two dimer configurations of  $G$ , and  $M_1 \cup M_2$  be their superposition.



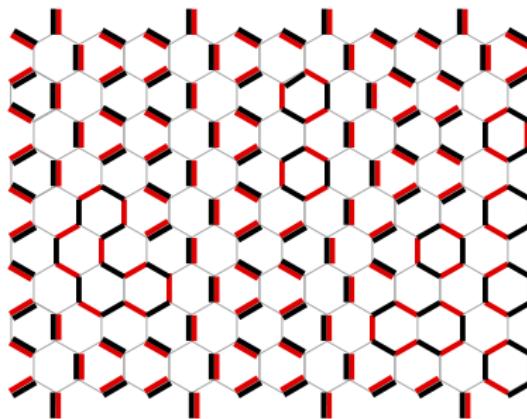
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- $M_1 \cup M_2$  is a disjoint union of alternating cycles, where an alternating cycles has edges alternating between  $M_1$  and  $M_2$ . Alternating cycles of length 2 are called doubled edges.