

# Diffusions non-linéaires: entropies relatives, “best matching” et délais

Jean Dolbeault

<http://www.ceremade.dauphine.fr/~dolbeaul>

Ceremade, Université Paris-Dauphine

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# A. Fast diffusion equations: entropy, linearization, inequalities, improvements

- 1 entropy methods
- 2 linearization of the entropy
- 3 improved Gagliardo-Nirenberg inequalities

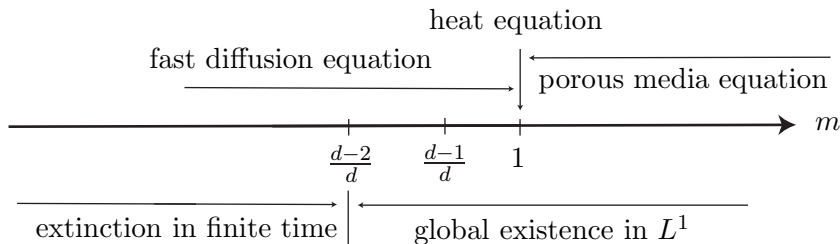
# Fast diffusion equations: entropy methods

# Existence, classical results

$$u_t = \Delta u^m \quad x \in \mathbb{R}^d, \quad t > 0$$

Self-similar (Barenblatt) function:  $\mathcal{U}(t) = O(t^{-d/(2-d(1-m))})$  as  $t \rightarrow +\infty$

[Friedmann, Kamin, 1980]  $\|u(t, \cdot) - \mathcal{U}(t, \cdot)\|_{L^\infty} = o(t^{-d/(2-d(1-m))})$



Existence theory, critical values of the parameter  $m$

# Time-dependent rescaling, Free energy

● *Time-dependent rescaling*: Take  $u(\tau, y) = R^{-d}(\tau) v(t, y/R(\tau))$  where

$$\frac{dR}{d\tau} = R^{d(1-m)-1}, \quad R(0) = 1, \quad t = \log R$$

● The function  $v$  solves a Fokker-Planck type equation

$$\frac{\partial v}{\partial t} = \Delta v^m + \nabla \cdot (x v), \quad v|_{\tau=0} = u_0$$

● [Ralston, Newman, 1984] Lyapunov functional:

*Generalized entropy* or *Free energy*

$$\mathcal{F}[v] := \int_{\mathbb{R}^d} \left( \frac{v^m}{m-1} + \frac{1}{2} |x|^2 v \right) dx - \mathcal{F}_0$$

Entropy production is measured by the *Generalized Fisher information*

$$\frac{d}{dt} \mathcal{F}[v] = -\mathcal{I}[v], \quad \mathcal{I}[v] := \int_{\mathbb{R}^d} v \left| \frac{\nabla v^m}{v} + x \right|^2 dx$$

# Relative entropy and entropy production

🔴 *Stationary solution:* choose  $C$  such that  $\|v_\infty\|_{L^1} = \|u\|_{L^1} = M > 0$

$$v_\infty(x) := \left(C + \frac{1-m}{2m} |x|^2\right)_+^{-1/(1-m)}$$

*Relative entropy:* Fix  $\mathcal{F}_0$  so that  $\mathcal{F}[v_\infty] = 0$

🔴 *Entropy – entropy production inequality*

## Theorem

$d \geq 3$ ,  $m \in [\frac{d-1}{d}, +\infty)$ ,  $m > \frac{1}{2}$ ,  $m \neq 1$

$$\mathcal{I}[v] \geq 2 \mathcal{F}[v]$$

## Corollary

A solution  $v$  with initial data  $u_0 \in L^1_+(\mathbb{R}^d)$  such that  $|x|^2 u_0 \in L^1(\mathbb{R}^d)$ ,  $u_0^m \in L^1(\mathbb{R}^d)$  satisfies  $\mathcal{F}[v(t, \cdot)] \leq \mathcal{F}[u_0] e^{-2t}$

# An equivalent formulation: Gagliardo-Nirenberg inequalities

$$\mathcal{F}[v] = \int_{\mathbb{R}^d} \left( \frac{v^m}{m-1} + \frac{1}{2} |x|^2 v \right) dx - \mathcal{F}_0 \leq \frac{1}{2} \int_{\mathbb{R}^d} v \left| \frac{\nabla v^m}{v} + x \right|^2 dx = \frac{1}{2} \mathcal{I}[v]$$

Rewrite it with  $p = \frac{1}{2m-1}$ ,  $v = w^{2p}$ ,  $v^m = w^{p+1}$  as

$$\frac{1}{2} \left( \frac{2m}{2m-1} \right)^2 \int_{\mathbb{R}^d} |\nabla w|^2 dx + \left( \frac{1}{1-m} - d \right) \int_{\mathbb{R}^d} |w|^{1+p} dx - K \geq 0$$

- for some  $\gamma$ ,  $K = K_0 \left( \int_{\mathbb{R}^d} v dx = \int_{\mathbb{R}^d} w^{2p} dx \right)^\gamma$
- $w = w_\infty = v_\infty^{1/2p}$  is optimal

## Theorem

[Del Pino, J.D.] With  $1 < p \leq \frac{d}{d-2}$  (fast diffusion case) and  $d \geq 3$

$$\|w\|_{L^{2p}(\mathbb{R}^d)} \leq C_{p,d}^{\text{GN}} \|\nabla w\|_{L^2(\mathbb{R}^d)}^\theta \|w\|_{L^{p+1}(\mathbb{R}^d)}^{1-\theta}$$

$$C_{p,d}^{\text{GN}} = \left( \frac{y(p-1)^2}{2\pi d} \right)^{\frac{\theta}{2}} \left( \frac{2y-d}{2y} \right)^{\frac{1}{2p}} \left( \frac{\Gamma(y)}{\Gamma(y-\frac{d}{2})} \right)^{\frac{\theta}{d}}, \quad \theta = \frac{d(p-1)}{p(d+2-(d-2)p)}, \quad y = \frac{p+1}{p-1}$$

## ... a proof by the Bakry-Emery method

Consider the generalized Fisher information

$$\mathcal{I}[v] := \int_{\mathbb{R}^d} v |z|^2 dx \quad \text{with} \quad z := \frac{\nabla v^m}{v} + x$$

and compute

$$\frac{d}{dt} \mathcal{I}[v(t, \cdot)] + 2 \mathcal{I}[v(t, \cdot)] = -2(m-1) \int_{\mathbb{R}^d} u^m (\operatorname{div} z)^2 dx - 2 \sum_{i,j=1}^d \int_{\mathbb{R}^d} u^m (\partial_i z^j)^2 dx$$

- the Fisher information decays exponentially:

$$\mathcal{I}[v(t, \cdot)] \leq \mathcal{I}[u_0] e^{-2t}$$

- $\lim_{t \rightarrow \infty} \mathcal{I}[v(t, \cdot)] = 0$  and  $\lim_{t \rightarrow \infty} \mathcal{F}[v(t, \cdot)] = 0$
- $\frac{d}{dt} \left( \mathcal{I}[v(t, \cdot)] - 2 \mathcal{F}[v(t, \cdot)] \right) \leq 0$  means  $\mathcal{I}[v] \geq 2 \mathcal{F}[v]$

[Carrillo, Toscani], [Juengel, Markowich, Toscani], [Carrillo, Juengel, Markowich, Toscani, Unterreiter], [Carrillo, Vázquez]



## The Bakry-Emery method: details (1/2)

With  $z(x, t) := \eta \nabla u^{m-1} - 2x$ , the equation can be rewritten as

$$\frac{\partial u}{\partial t} + \nabla \cdot (u z) = 0$$

(up to a time rescaling, which introduces a factor 2) and we have

$$\frac{\partial z}{\partial t} = \eta(1-m)\nabla(u^{m-2}\nabla \cdot (u z)) \quad \text{and} \quad \nabla \otimes z = \eta \nabla \otimes \nabla u^{m-1} - 2\text{Id}$$

$$\frac{d}{dt} \int_{\mathbb{R}^d} u |z|^2 dx = \underbrace{\int_{\mathbb{R}^d} \frac{\partial u}{\partial t} |z|^2 dx}_{\text{(I)}} + 2 \underbrace{\int_{\mathbb{R}^d} u z \cdot \frac{\partial z}{\partial t} dx}_{\text{(II)}}$$

$$\begin{aligned} \text{(I)} &= \int_{\mathbb{R}^d} \frac{\partial u}{\partial t} |z|^2 dx = \int_{\mathbb{R}^d} \nabla \cdot (u z) |z|^2 dx \\ &= 2\eta(1-m) \int_{\mathbb{R}^d} u^{m-2} (\nabla u \cdot z)^2 dx + 2\eta(1-m) \int_{\mathbb{R}^d} u^{m-1} (\nabla u \cdot z) (\nabla \cdot z) dx \\ &\quad + 2\eta(1-m) \int_{\mathbb{R}^d} u^{m-1} (z \otimes \nabla u) : (\nabla \otimes z) dx - 4 \int_{\mathbb{R}^d} u |z|^2 dx \end{aligned}$$

## The Bakry-Emery method: details (2/2)

$$\begin{aligned}
 \text{(II)} &= 2 \int_{\mathbb{R}^d} u z \cdot \frac{\partial z}{\partial t} dx \\
 &= -2 \eta (1 - m) \int_{\mathbb{R}^d} \left[ u^m (\nabla \cdot z)^2 + 2 u^{m-1} (\nabla u \cdot z) (\nabla \cdot z) + u^{m-2} (\nabla u \cdot z)^2 \right] dx \\
 &\quad + \int_{\mathbb{R}^d} \frac{\partial u}{\partial t} |z|^2 dx + 4 \int_{\mathbb{R}^d} u |z|^2 dx \\
 &= -2 \eta (1 - m) \int_{\mathbb{R}^d} u^{m-2} \left[ u^2 (\nabla \cdot z)^2 + u (\nabla u \cdot z) (\nabla \cdot z) \right. \\
 &\quad \left. + \frac{1-m}{m} \int_{\mathbb{R}^d} u^m \left( |\nabla z|^2 - (1-m) (\nabla \cdot z)^2 \right) dx \right] dx
 \end{aligned}$$

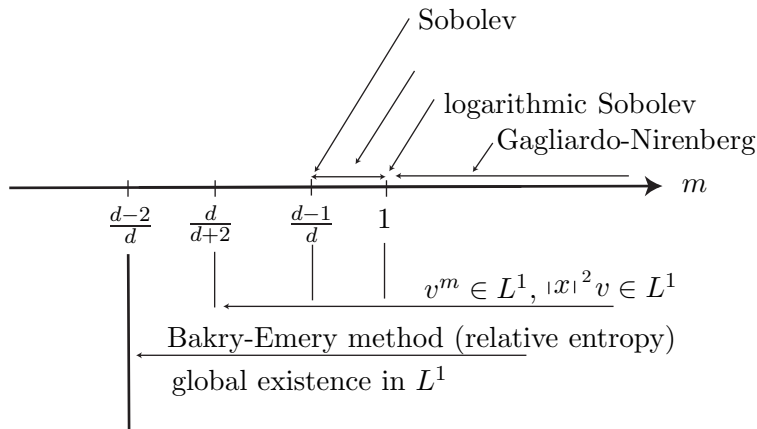
By the arithmetic geometric inequality, we know that

$$|\nabla z|^2 - (1 - m) (\nabla \cdot z)^2 \geq 0$$

if  $1 - m \leq 1/d$ , that is, if  $m \geq m_1 = 1 - 1/d$

# Fast diffusion: finite mass regime

Inequalities...



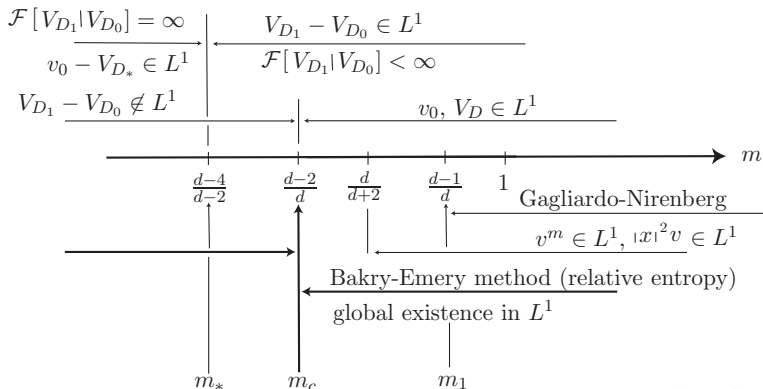
... existence of solutions of  $u_t = \Delta u^m$

# Fast diffusion equations: the infinite mass regime by linearization of the entropy

# Extension to the infinite mass regime, finite time vanishing

- If  $m > m_c := \frac{d-2}{d} \leq m < m_1$ , solutions globally exist in  $L^1(\mathbb{R}^d)$  and the Barenblatt self-similar solution has finite mass
- For  $m \leq m_c$ , the Barenblatt self-similar solution has infinite mass

*Extension to  $m \leq m_c$  ? Work in relative variables !*



# Entropy methods and linearization: intermediate asymptotics, vanishing

[A. Blanchet, M. Bonforte, J.D., G. Grillo, J.L. Vázquez]

$$\frac{\partial u}{\partial \tau} = -\nabla \cdot (u \nabla u^{m-1}) = \frac{1-m}{m} \Delta u^m \quad (1)$$

- $m_c < m < 1$ ,  $T = +\infty$ : intermediate asymptotics,  $\tau \rightarrow +\infty$

$$R(\tau) := (T + \tau)^{\frac{1}{d(m-m_c)}}$$

- $0 < m < m_c$ ,  $T < +\infty$ : vanishing in finite time  $\lim_{\tau \nearrow T} u(\tau, y) = 0$

$$R(\tau) := (T - \tau)^{-\frac{1}{d(m_c-m)}}$$

Self-similar *Barenblatt type solutions* exists for any  $m$

$$t := \frac{1-m}{2} \log \left( \frac{R(\tau)}{R(0)} \right) \quad \text{and} \quad x := \sqrt{\frac{1}{2d|m-m_c|} \frac{y}{R(\tau)}}$$

*Generalized Barenblatt profiles:*  $V_D(x) := (D + |x|^2)^{\frac{1}{m-1}}$

# Sharp rates of convergence

Assumptions on the initial datum  $v_0$

**(H1)**  $V_{D_0} \leq v_0 \leq V_{D_1}$  for some  $D_0 > D_1 > 0$

**(H2)** if  $d \geq 3$  and  $m \leq m_*$ ,  $(v_0 - V_D)$  is integrable for a suitable  $D \in [D_1, D_0]$

## Theorem

[Blanchet, Bonforte, J.D., Grillo, Vázquez] Under Assumptions (H1)-(H2), if  $m < 1$  and  $m \neq m_* := \frac{d-4}{d-2}$ , the entropy decays according to

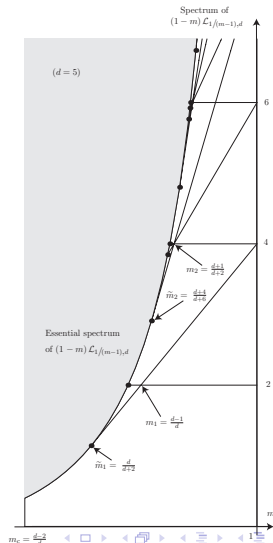
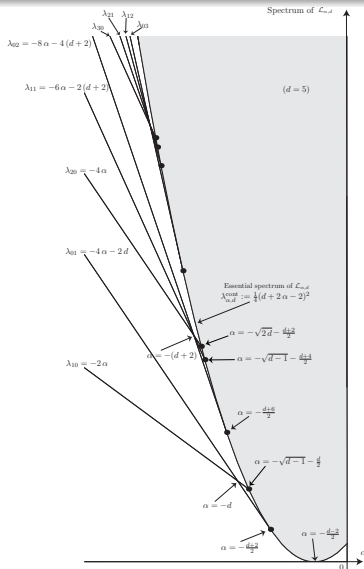
$$\mathcal{F}[v(t, \cdot)] \leq C e^{-2(1-m)\Lambda_{\alpha,d} t} \quad \forall t \geq 0$$

where  $\Lambda_{\alpha,d} > 0$  is the best constant in the Hardy-Poincaré inequality

$$\Lambda_{\alpha,d} \int_{\mathbb{R}^d} |f|^2 d\mu_{\alpha-1} \leq \int_{\mathbb{R}^d} |\nabla f|^2 d\mu_{\alpha} \quad \forall f \in H^1(d\mu_{\alpha})$$

with  $\alpha := 1/(m-1) < 0$ ,  $d\mu_{\alpha} := h_{\alpha} dx$ ,  $h_{\alpha}(x) := (1 + |x|^2)^{\alpha}$

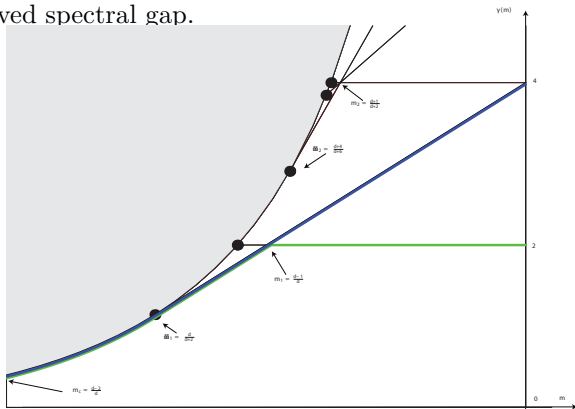
# Plots ( $d = 5$ )





# Improved asymptotic rates

[Bonforte, J.D., Grillo, Vázquez] Assume that  $m \in (m_1, 1)$ ,  $d \geq 3$ . Under Assumption (H1), if  $v$  is a solution of the fast diffusion equation with initial datum  $v_0$  such that  $\int_{\mathbb{R}^d} x v_0 dx = 0$ , then the asymptotic convergence holds with an improved rate corresponding to the improved spectral gap.



# Higher order matching asymptotics

[J.D., G. Toscani] For some  $m \in (m_c, 1)$  with  $m_c := (d-2)/d$ , we consider on  $\mathbb{R}^d$  the fast diffusion equation

$$\frac{\partial u}{\partial \tau} + \nabla \cdot (u \nabla u^{m-1}) = 0$$

*Without choosing  $R$* , we may define the function  $v$  such that

$$u(\tau, y + x_0) = R^{-d} v(t, x), \quad R = R(\tau), \quad t = \frac{1}{2} \log R, \quad x = \frac{y}{R}$$

Then  $v$  has to be a solution of

$$\frac{\partial v}{\partial t} + \nabla \cdot \left[ v \left( \sigma^{\frac{d}{2}(m-m_c)} \nabla v^{m-1} - 2x \right) \right] = 0 \quad t > 0, \quad x \in \mathbb{R}^d$$

with (as long as we make no assumption on  $R$ )

$$2 \sigma^{-\frac{d}{2}(m-m_c)} = R^{1-d(1-m)} \frac{dR}{d\tau}$$

## Refined relative entropy

Consider the family of the Barenblatt profiles

$$B_\sigma(x) := \sigma^{-\frac{d}{2}} \left( C_M + \frac{1}{\sigma} |x|^2 \right)^{\frac{1}{m-1}} \quad \forall x \in \mathbb{R}^d \quad (2)$$

Note that  $\sigma$  is a function of  $t$ : as long as  $\frac{d\sigma}{dt} \neq 0$ , the Barenblatt profile  $B_\sigma$  is *not* a solution (it plays the role of a *local Gibbs state*) but we may still consider the relative entropy

$$\mathcal{F}_\sigma[v] := \frac{1}{m-1} \int_{\mathbb{R}^d} [v^m - B_\sigma^m - m B_\sigma^{m-1} (v - B_\sigma)] dx$$

The time derivative of this relative entropy is

$$\frac{d}{dt} \mathcal{F}_{\sigma(t)}[v(t, \cdot)] = \underbrace{\frac{d\sigma}{dt} \left( \frac{d}{d\sigma} \mathcal{F}_\sigma[v] \right)_{|\sigma=\sigma(t)}}_{\text{choose it} = 0} + \frac{m}{m-1} \int_{\mathbb{R}^d} \left( v^{m-1} - B_{\sigma(t)}^{m-1} \right) \frac{\partial v}{\partial t} dx$$

$$\iff \text{Minimize } \mathcal{F}_\sigma[v] \text{ w.r.t. } \sigma \iff \int_{\mathbb{R}^d} |x|^2 B_\sigma dx = \int_{\mathbb{R}^d} |x|^2 v dx$$

# The entropy / entropy production estimate

Using the new change of variables, we know that

$$\frac{d}{dt} \mathcal{F}_{\sigma(t)}[v(t, \cdot)] = - \frac{m \sigma(t)^{\frac{d}{2}(m-m_c)}}{1-m} \int_{\mathbb{R}^d} v \left| \nabla \left[ v^{m-1} - B_{\sigma(t)}^{m-1} \right] \right|^2 dx$$

Let  $w := v/B_\sigma$  and observe that the relative entropy can be written as

$$\mathcal{F}_\sigma[v] = \frac{m}{1-m} \int_{\mathbb{R}^d} \left[ w - 1 - \frac{1}{m} (w^m - 1) \right] B_\sigma^m dx$$

(Repeating) define the *relative Fisher information* by

$$\mathcal{I}_\sigma[v] := \int_{\mathbb{R}^d} \left| \frac{1}{m-1} \nabla \left[ (w^{m-1} - 1) B_\sigma^{m-1} \right] \right|^2 B_\sigma w dx$$

so that 
$$\frac{d}{dt} \mathcal{F}_{\sigma(t)}[v(t, \cdot)] = -m(1-m) \sigma(t) \mathcal{I}_{\sigma(t)}[v(t, \cdot)] \quad \forall t > 0$$

*When linearizing, one more mode is killed and  $\sigma(t)$  scales out*

# Improved rates of convergence



## Theorem (J.D., G. Toscani)

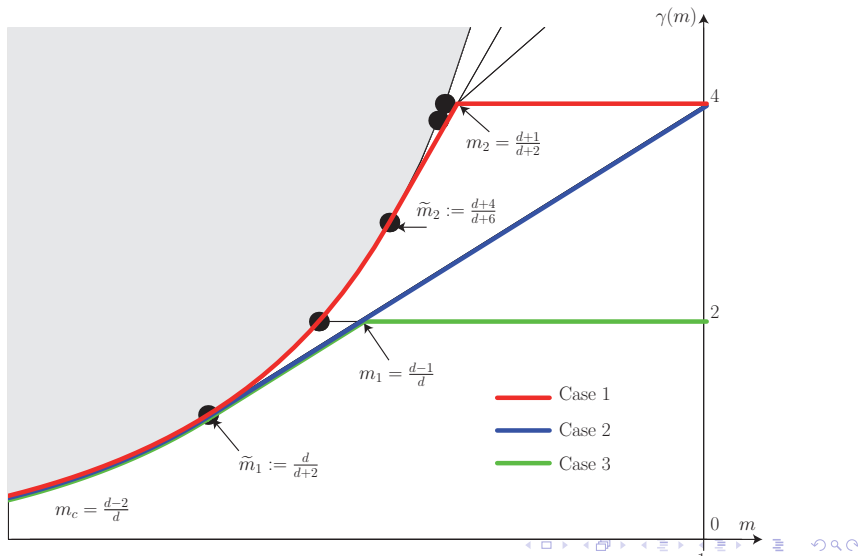
Let  $m \in (\tilde{m}_1, 1)$ ,  $d \geq 2$ ,  $v_0 \in L^1_+(\mathbb{R}^d)$  such that  $v_0^m, |y|^2 v_0 \in L^1(\mathbb{R}^d)$

$$\mathcal{F}[v(t, \cdot)] \leq C e^{-2\gamma(m)t} \quad \forall t \geq 0$$

where

$$\gamma(m) = \begin{cases} \frac{((d-2)m - (d-4))^2}{4(1-m)} & \text{if } m \in (\tilde{m}_1, \tilde{m}_2] \\ 4(d+2)m - 4d & \text{if } m \in [\tilde{m}_2, m_2] \\ 4 & \text{if } m \in [m_2, 1) \end{cases}$$

# Spectral gaps and best constants



# Comments

- 1 A result by [Denzler, Koch, McCann] *Higher order time asymptotics of fast diffusion in Euclidean space: a dynamical systems approach*
- 2 The constant  $C$  in

$$\mathcal{F}[v(t, \cdot)] \leq C e^{-2\gamma(m)t} \quad \forall t \geq 0$$

can be made explicit, under additional restrictions on the initial data [Bonforte, J.D., Grillo, Vázquez]

# An explicit constant $C$ ?

$$\frac{d}{dt} \mathcal{F}[w(t, \cdot)] = -\mathcal{I}[w(t, \cdot)] \quad \forall t > 0$$

$$h^{m-2} \int_{\mathbb{R}^d} |f|^2 V_D^{2-m} dx \leq 2 \mathcal{F}[w] \leq h^{2-m} \int_{\mathbb{R}^d} |f|^2 V_D^{2-m} dx$$

where  $f := (w - 1) V_D^{m-1}$ ,  $h := \max\{\sup_{\mathbb{R}^d} w(t, \cdot), 1/\inf_{\mathbb{R}^d} w(t, \cdot)\}$

$$\int_{\mathbb{R}^d} |\nabla f|^2 V_D dx \leq h^{5-2m} \mathcal{I}[w] + d(1-m) [h^{4(2-m)} - 1] \int_{\mathbb{R}^d} |f|^2 V_D^{2-m} dx$$

$$0 \leq h - 1 \leq C \mathcal{F}^{\frac{1-m}{d+2-(d+1)m}}$$

## Corollary

$\mathcal{F}[w(t, \cdot)] \leq G(t, h(0), \mathcal{F}[w(0, \cdot)])$  for any  $t \geq 0$ , where

$$\frac{dG}{dt} = -2 \frac{\Lambda_{\alpha,d} - Y(h)}{[1 + X(h)] h^{2-m}} G, \quad h = 1 + C G^{\frac{1-m}{d+2-(d+1)m}}, \quad G(0) = \mathcal{F}[w(0, \cdot)]$$



# Gagliardo-Nirenberg and Sobolev inequalities : improvements

[J.D., G. Toscani]

# Best matching Barenblatt profiles

(Repeating) Consider the *fast diffusion equation*

$$\frac{\partial u}{\partial t} + \nabla \cdot \left[ u \left( \sigma^{\frac{d}{2}(m-m_c)} \nabla u^{m-1} - 2x \right) \right] = 0 \quad t > 0, \quad x \in \mathbb{R}^d$$

with a nonlocal, time-dependent diffusion coefficient

$$\sigma(t) = \frac{1}{K_M} \int_{\mathbb{R}^d} |x|^2 u(x, t) dx, \quad K_M := \int_{\mathbb{R}^d} |x|^2 B_1(x) dx$$

where

$$B_\lambda(x) := \lambda^{-\frac{d}{2}} \left( C_M + \frac{1}{\lambda} |x|^2 \right)^{\frac{1}{m-1}} \quad \forall x \in \mathbb{R}^d$$

and define the relative entropy

$$\mathcal{F}_\lambda[u] := \frac{1}{m-1} \int_{\mathbb{R}^d} \left[ u^m - B_\lambda^m - m B_\lambda^{m-1} (u - B_\lambda) \right] dx$$

## Three ingredients for *global improvements*

- 1  $\inf_{\lambda>0} \mathcal{F}_\lambda[u(x, t)] = \mathcal{F}_{\sigma(t)}[u(x, t)]$  so that

$$\frac{d}{dt} \mathcal{F}_{\sigma(t)}[u(x, t)] = -\mathcal{J}_{\sigma(t)}[u(\cdot, t)]$$

where the relative Fisher information is

$$\mathcal{J}_\lambda[u] := \lambda^{\frac{d}{2}(m-m_c)} \frac{m}{1-m} \int_{\mathbb{R}^d} u \left| \nabla u^{m-1} - \nabla B_\lambda^{m-1} \right|^2 dx$$

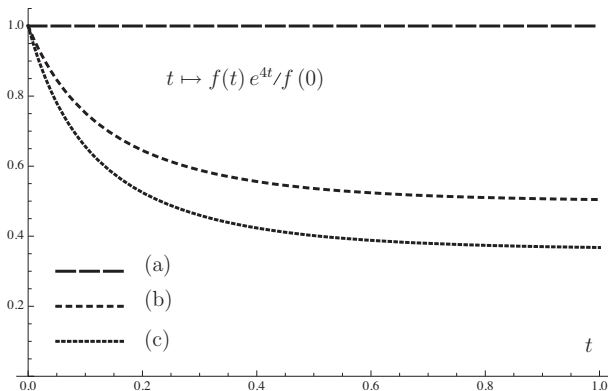
- 2 In the *Bakry-Emery method*, there is *an additional (good) term*

$$4 \left[ 1 + 2 C_{m,d} \frac{\mathcal{F}_{\sigma(t)}[u(\cdot, t)]}{M^\gamma \sigma_0^{\frac{d}{2}(1-m)}} \right] \frac{d}{dt} (\mathcal{F}_{\sigma(t)}[u(\cdot, t)]) \geq \frac{d}{dt} (\mathcal{J}_{\sigma(t)}[u(\cdot, t)])$$

- 3 The *Csiszár-Kullback inequality* is also improved

$$\mathcal{F}_\sigma[u] \geq \frac{m}{8 \int_{\mathbb{R}^d} B_1^m dx} C_M^2 \|u - B_\sigma\|_{L^1(\mathbb{S}^d)}^2$$

# improved decay for the relative entropy



**Figure:** Upper bounds on the decay of the relative entropy:  $t \mapsto f(t) e^{4t} / f(0)$

(a): estimate given by the entropy-entropy production method

(b): exact solution of a simplified equation

(c): numerical solution (found by a shooting method)

# A Csiszár-Kullback(-Pinsker) inequality

Let  $m \in (\tilde{m}_1, 1)$  with  $\tilde{m}_1 = \frac{d}{d+2}$  and consider the relative entropy

$$\mathcal{F}_\sigma[u] := \frac{1}{m-1} \int_{\mathbb{R}^d} [u^m - B_\sigma^m - m B_\sigma^{m-1} (u - B_\sigma)] dx$$

## Theorem

Let  $d \geq 1$ ,  $m \in (\tilde{m}_1, 1)$  and assume that  $u$  is a nonnegative function in  $\mathcal{L}^1(\mathbb{R}^d)$  such that  $u^m$  and  $x \mapsto |x|^2 u$  are both integrable on  $\mathbb{R}^d$ . If  $\|u\|_{L^1(\mathbb{R}^d)} = M$  and  $\int_{\mathbb{R}^d} |x|^2 u dx = \int_{\mathbb{R}^d} |x|^2 B_\sigma dx$ , then

$$\frac{\mathcal{F}_\sigma[u]}{\sigma^{\frac{d}{2}(1-m)}} \geq \frac{m}{8 \int_{\mathbb{R}^d} B_1^m dx} \left( C_M \|u - B_\sigma\|_{L^1(\mathbb{R}^d)} + \frac{1}{\sigma} \int_{\mathbb{R}^d} |x|^2 |u - B_\sigma| dx \right)^2$$

# Csiszár-Kullback(-Pinsker): proof (1/2)

Let  $v := u/B_\sigma$  and  $d\mu_\sigma := B_\sigma^m dx$

$$\begin{aligned} \int_{\mathbb{R}^d} (v - 1) d\mu_\sigma &= \int_{\mathbb{R}^d} B_\sigma^{m-1} (u - B_\sigma) dx \\ &= \sigma^{\frac{d}{2}(1-m)} C_M \int_{\mathbb{R}^d} (u - B_\sigma) dx + \sigma^{\frac{d}{2}(m_c-m)} \int_{\mathbb{R}^d} |x|^2 (u - B_\sigma) dx = 0 \end{aligned}$$

$$\int_{\mathbb{R}^d} (v - 1) d\mu_\sigma = \int_{v>1} (v - 1) d\mu_\sigma - \int_{v<1} (1 - v) d\mu_\sigma = 0$$

$$\int_{\mathbb{R}^d} |v - 1| d\mu_\sigma = \int_{v>1} (v - 1) d\mu_\sigma + \int_{v<1} (1 - v) d\mu_\sigma$$

$$\int_{\mathbb{R}^d} |u - B_\sigma| B_\sigma^{m-1} dx = \int_{\mathbb{R}^d} |v - 1| d\mu_\sigma = 2 \int_{v<1} |v - 1| d\mu_\sigma$$

## Csiszár-Kullback(-Pinsker): proof (2/2)

A Taylor expansion shows that

$$\begin{aligned}\mathcal{F}_\sigma[u] &= \frac{1}{m-1} \int_{\mathbb{R}^d} [v^m - 1 - m(v-1)] d\mu_\sigma = \frac{m}{2} \int_{\mathbb{R}^d} \xi^{m-2} |v-1|^2 d\mu_\sigma \\ &\geq \frac{m}{2} \int_{v < 1} |v-1|^2 d\mu_\sigma\end{aligned}$$

Using the Cauchy-Schwarz inequality, we get

$$\left( \int_{v < 1} |v-1| d\mu_\sigma \right)^2 = \left( \int_{v < 1} |v-1| B_\sigma^{\frac{m}{2}} B_\sigma^{\frac{m}{2}} dx \right)^2 \leq \int_{v < 1} |v-1|^2 d\mu_\sigma \int_{\mathbb{R}^d} B_\sigma^m dx$$

and finally obtain that

$$\mathcal{F}_\sigma[u] \geq \frac{m}{2} \frac{\left( \int_{v < 1} |v-1| d\mu_\sigma \right)^2}{\int_{\mathbb{R}^d} B_\sigma^m dx} = \frac{m}{8} \frac{\left( \int_{\mathbb{R}^d} |u - B_\sigma| B_\sigma^{m-1} dx \right)^2}{\int_{\mathbb{R}^d} B_\sigma^m dx}$$

# An improved Gagliardo-Nirenberg inequality: the setting

The inequality

$$\|f\|_{L^{2p}(\mathbb{S}^d)} \leq C_{p,d}^{\text{GN}} \|\nabla f\|_{L^2(\mathbb{S}^d)}^\theta \|f\|_{L^{p+1}(\mathbb{S}^d)}^{1-\theta}$$

with  $\theta = \theta(p) := \frac{p-1}{p} \frac{d}{d+2-p(d-2)}$ ,  $1 < p \leq \frac{d}{d-2}$  if  $d \geq 3$  and  $1 < p < \infty$  if  $d = 2$ , can be rewritten, in a non-scale invariant form, as

$$\int_{\mathbb{R}^d} |\nabla f|^2 dx + \int_{\mathbb{R}^d} |f|^{p+1} dx \geq K_{p,d} \left( \int_{\mathbb{R}^d} |f|^{2p} dx \right)^\gamma$$

with  $\gamma = \gamma(p, d) := \frac{d+2-p(d-2)}{d-p(d-4)}$ . Optimal function are given by

$$f_{M,y,\sigma}(x) = \frac{1}{\sigma^{\frac{d}{2}}} \left( C_M + \frac{|x-y|^2}{\sigma} \right)^{-\frac{1}{p-1}} \quad \forall x \in \mathbb{R}^d$$

where  $C_M$  is determined by  $\int_{\mathbb{R}^d} f_{M,y,\sigma}^{2p} dx = M$

$$\mathfrak{M}_d := \{f_{M,y,\sigma} : (M, y, \sigma) \in \mathcal{M}_d := (0, \infty) \times \mathbb{R}^d \times (0, \infty)\}$$



# An improved Gagliardo-Nirenberg inequality (1/2)

*Relative entropy functional*

$$\mathcal{R}^{(p)}[f] := \inf_{g \in \mathfrak{M}_d^{(p)}} \int_{\mathbb{R}^d} \left[ g^{1-p} (|f|^{2p} - g^{2p}) - \frac{2p}{p+1} (|f|^{p+1} - g^{p+1}) \right] dx$$

## Theorem

Let  $d \geq 2$ ,  $p > 1$  and assume that  $p < d/(d-2)$  if  $d \geq 3$ . If

$$\frac{\int_{\mathbb{R}^d} |x|^2 |f|^{2p} dx}{\left( \int_{\mathbb{R}^d} |f|^{2p} dx \right)^\gamma} = \frac{d(p-1) \sigma_* M_*^{\gamma-1}}{d+2-p(d-2)}, \quad \sigma_*(p) := \left( 4 \frac{d+2-p(d-2)}{(p-1)^2 (p+1)} \right)^{\frac{4p}{d-p(d-4)}}$$

for any  $f \in \mathcal{L}^{p+1} \cap \mathcal{D}^{1,2}(\mathbb{R}^d)$ , then we have

$$\int_{\mathbb{R}^d} |\nabla f|^2 dx + \int_{\mathbb{R}^d} |f|^{p+1} dx - K_{p,d} \left( \int_{\mathbb{R}^d} |f|^{2p} dx \right)^\gamma \geq C_{p,d} \frac{(\mathcal{R}^{(p)}[f])^2}{\left( \int_{\mathbb{R}^d} |f|^{2p} dx \right)^\gamma}$$

# An improved Gagliardo-Nirenberg inequality (2/2)



A Csiszár-Kullback inequality

$$\mathcal{R}^{(p)}[f] \geq C_{CK} \|f\|_{L^{2p}(\mathbb{S}^d)}^{2p(\gamma-2)} \inf_{g \in \mathfrak{M}_d^{(p)}} \| |f|^{2p} - g^{2p} \|_{L^1(\mathbb{S}^d)}^2$$

with  $C_{CK} = \frac{p-1}{p+1} \frac{d+2-p(d-2)}{32p} \sigma_*^d \frac{p-1}{4p} M_*^{1-\gamma}$ . Let

$$\mathfrak{C}_{p,d} := C_{d,p} C_{CK}^2$$

## Corollary

*Under previous assumptions, we have*

$$\begin{aligned} \int_{\mathbb{R}^d} |\nabla f|^2 dx + \int_{\mathbb{R}^d} |f|^{p+1} dx - K_{p,d} \left( \int_{\mathbb{R}^d} |f|^{2p} dx \right)^\gamma \\ \geq \mathfrak{C}_{p,d} \|f\|_{L^{2p}(\mathbb{S}^d)}^{2p(\gamma-4)} \inf_{g \in \mathfrak{M}_d^{(p)}} \| |f|^{2p} - g^{2p} \|_{L^1(\mathbb{S}^d)}^4 \end{aligned}$$

## Conclusion 1: improved inequalities

- 1 We have found an improvement of an optimal Gagliardo-Nirenberg inequality, which provides an explicit measure of the distance to the manifold of optimal functions.
- 2 The method is based on the nonlinear flow
- 3 The explicit improvement gives (is equivalent to) an improved entropy – entropy production inequality

## Conclusion 2: improved rates

If  $m \in (m_1, 1)$ , with

$$f(t) := \mathcal{F}_{\sigma(t)}[u(\cdot, t)]$$

$$\sigma(t) = \frac{1}{K_M} \int_{\mathbb{R}^d} |x|^2 u(x, t) dx$$

$$j(t) := \mathcal{J}_{\sigma(t)}[u(\cdot, t)]$$

$$\mathcal{J}_{\sigma}[u] := \frac{m \sigma^{\frac{d}{2}(m-m_c)}}{1-m} \int_{\mathbb{R}^d} u |\nabla u^{m-1} - \nabla \mathfrak{B}_{\sigma}^{m-1}|^2 dx$$

we can write a system of coupled ODEs

$$\begin{cases} f' = -j \leq 0 \\ \sigma' = -2d \frac{(1-m)^2}{m K_M} \sigma^{\frac{d}{2}(m-m_c)} f \leq 0 \\ j' + 4j = \frac{d}{2}(m-m_c) \left[ \frac{j}{\sigma} + 4d(1-m) \frac{f}{\sigma} \right] \sigma' - r \end{cases} \quad (3)$$

In the rescaled variables, we have found an *improved decay* (algebraic rate) of the relative entropy. This is a new nonlinear effect, which matters for the initial time layer

## Conclusion 3: Best matching Barenblatt profiles are delayed

Let  $u$  be such that

$$v(\tau, x) = \frac{\mu^d}{R(D\tau)^d} u\left(\frac{1}{2} \log R(D\tau), \frac{\mu x}{R(D\tau)}\right)$$

with  $\tau \mapsto R(\tau)$  given as the solution to

$$\frac{1}{R} \frac{dR}{d\tau} = \left( \frac{\mu^2}{K_M} \int_{\mathbb{R}^d} |x|^2 v(\tau, x) dx \right)^{-\frac{d}{2}(m-m_c)}, \quad R(0) = 1$$

Then

$$\frac{1}{R} \frac{dR}{d\tau} = \left[ R^2(\tau) \sigma\left(\frac{1}{2} \log R(D\tau)\right) \right]^{-\frac{d}{2}(m-m_c)}$$

that is  $R(\tau) = R_0(\tau) \leq R_0(\tau)$  where  $\frac{1}{R} \frac{dR}{d\tau} = (R_0^2(\tau) \sigma(0))^{-\frac{d}{2}(m-m_c)}$   
 and asymptotically as  $\tau \rightarrow \infty$ ,  $R(\tau) = R_0(\tau - \delta)$  for some **delay**  $\delta > 0$

## B. Fast diffusion equations: New points of view

- 1 improved inequalities and scalings
- 2 scalings and a concavity property
- 3 improved rates and best matching

# Improved inequalities and scalings

# The logarithmic Sobolev inequality

$d\mu = \mu dx$ ,  $\mu(x) = (2\pi)^{-d/2} e^{-|x|^2/2}$ , on  $\mathbb{R}^d$  with  $d \geq 1$

*Gaussian logarithmic Sobolev inequality*

$$\int_{\mathbb{R}^2} |\nabla u|^2 d\mu \geq \frac{1}{2} \int_{\mathbb{R}^2} |u|^2 \log |u|^2 d\mu$$

for any function  $u \in H^1(\mathbb{R}^d, d\mu)$  such that  $\int_{\mathbb{R}^2} |u|^2 d\mu = 1$

$$\varphi(t) := \frac{d}{4} \left[ \exp\left(\frac{2t}{d}\right) - 1 - \frac{2t}{d} \right] \quad \forall t \in \mathbb{R}$$

[Bakry, Ledoux (2006)], [Fathi et al. (2014)], [Dolbeault, Toscani (2014)]

## Proposition

$$\int_{\mathbb{R}^2} |\nabla u|^2 d\mu - \frac{1}{2} \int_{\mathbb{R}^2} |u|^2 \log |u|^2 d\mu \geq \varphi \left( \int_{\mathbb{R}^2} |u|^2 \log |u|^2 d\mu \right)$$

$$\forall u \in H^1(\mathbb{R}^d, d\mu) \quad \text{s.t.} \quad \int |u|^2 d\mu = 1 \quad \text{and} \quad \int |x|^2 |u|^2 d\mu = d$$



# Consequences for the heat equation

Ornstein-Uhlenbeck equation (or backward Kolmogorov equation)

$$\frac{\partial f}{\partial t} = \Delta f - x \cdot \nabla f$$

with initial datum  $f_0 \in L^1_+(\mathbb{R}^d, (1 + |x|^2) d\mu)$  and define the *entropy* as

$$\mathcal{E}[f] := \int_{\mathbb{R}^2} f \log f d\mu, \quad \frac{d}{dt} \mathcal{E}[f] = -4 \int_{\mathbb{R}^2} |\nabla \sqrt{f}|^2 d\mu \leq -2 \mathcal{E}[f]$$

thus proving that  $\mathcal{E}[f(t, \cdot)] \leq \mathcal{E}[f_0] e^{-2t}$ . Moreover,

$$\frac{d}{dt} \int_{\mathbb{R}^2} f |x|^2 d\mu = 2 \int_{\mathbb{R}^2} f (d - |x|^2) d\mu$$

## Theorem

Assume that  $\mathcal{E}[f_0]$  is finite and  $\int_{\mathbb{R}^2} f_0 |x|^2 d\mu = d \int_{\mathbb{R}^2} f_0 d\mu$ . Then

$$\mathcal{E}[f(t, \cdot)] \leq -\frac{d}{2} \log \left[ 1 - \left( 1 - e^{-\frac{2}{d} \mathcal{E}[f_0]} \right) e^{-2t} \right] \quad \forall t \geq 0$$

# Gagliardo-Nirenberg inequalities and the FDE

$$\|\nabla w\|_{L^2(\mathbb{S}^d)}^\vartheta \|w\|_{L^{q+1}(\mathbb{S}^d)}^{1-\vartheta} \geq C_{\text{GN}} \|w\|_{L^{2q}(\mathbb{S}^d)}$$

With the right choice of the constants, the functional

$$J[w] := \frac{1}{4} (q^2 - 1) \int_{\mathbb{R}^d} |\nabla w|^2 dx + \beta \int_{\mathbb{R}^d} |w|^{q+1} dx - \mathcal{K} C_{\text{GN}}^\alpha \left( \int_{\mathbb{R}^d} |w|^{2q} dx \right)^{\frac{\alpha}{2q}}$$

is nonnegative and  $J[w] \geq J[w_*] = 0$

## Theorem

[Dolbeault-Toscani] *For some nonnegative, convex, increasing  $\varphi$*

$$J[w] \geq \varphi \left[ \beta \left( \int_{\mathbb{R}^d} |w_*|^{q+1} dx - \int_{\mathbb{R}^d} |w|^{q+1} dx \right) \right]$$

*for any  $w \in L^{q+1}(\mathbb{R}^d)$  such that  $\int_{\mathbb{R}^d} |\nabla w|^2 dx < \infty$  and  $\int_{\mathbb{R}^d} |w|^{2q} |x|^2 dx = \int_{\mathbb{R}^d} w_*^{2q} |x|^2 dx$*

Consequence for decay rates of relative Rényi entropies:  
 see [Carrillo-Toscani]

# Scalings and a concavity property

# The fast diffusion equation in original variables

Consider the nonlinear diffusion equation in  $\mathbb{R}^d$ ,  $d \geq 1$

$$\frac{\partial u}{\partial t} = \Delta u^p$$

with initial datum  $u(x, t = 0) = u_0(x) \geq 0$  such that  $\int_{\mathbb{R}^d} u_0 dx = 1$  and  $\int_{\mathbb{R}^d} |x|^2 u_0 dx < +\infty$ . The large time behavior of the solutions is governed by the source-type Barenblatt solutions

$$u_\star(t, x) := \frac{1}{(\kappa t^{1/\mu})^d} \mathcal{B}_\star\left(\frac{x}{\kappa t^{1/\mu}}\right)$$

where

$$\mu := 2 + d(p - 1), \quad \kappa := \left| \frac{2\mu p}{p - 1} \right|^{1/\mu}$$

and  $\mathcal{B}_\star$  is the Barenblatt profile

$$\mathcal{B}_\star(x) := \begin{cases} (C_\star - |x|^2)_+^{1/(p-1)} & \text{if } p > 1 \\ (C_\star + |x|^2)^{1/(p-1)} & \text{if } p < 1 \end{cases}$$

# The entropy

The *entropy* is defined by

$$E := \int_{\mathbb{R}^d} u^p dx$$

and the *Fisher information* by

$$I := \int_{\mathbb{R}^d} u |\nabla v|^2 dx \quad \text{with} \quad v = \frac{p}{p-1} u^{p-1}$$

If  $u$  solves the fast diffusion equation, then

$$E' = (1-p)I$$

To compute  $I'$ , we will use the fact that

$$\frac{\partial v}{\partial t} = (p-1) v \Delta v + |\nabla v|^2$$

$$F := E^\sigma \quad \text{with} \quad \sigma = \frac{\mu}{d(1-p)} = 1 + \frac{2}{1-p} \left( \frac{1}{d} + p - 1 \right) = \frac{2}{d} \frac{1}{1-p} - 1$$

has a linear growth asymptotically as  $t \rightarrow +\infty$

# The concavity property

## Theorem

[Toscani-Savaré] Assume that  $p \geq 1 - \frac{1}{d}$  if  $d > 1$  and  $p > 0$  if  $d = 1$ . Then  $F(t)$  is increasing,  $(1 - p) F''(t) \leq 0$  and

$$\lim_{t \rightarrow +\infty} \frac{1}{t} F(t) = (1 - p) \sigma \lim_{t \rightarrow +\infty} E^{\sigma-1} I = (1 - p) \sigma E_{\star}^{\sigma-1} I_{\star}$$

[Dolbeault-Toscani] The inequality

$$E^{\sigma-1} I \geq E_{\star}^{\sigma-1} I_{\star}$$

is equivalent to the Gagliardo-Nirenberg inequality

$$\|\nabla w\|_{L^2(\mathbb{S}^d)}^{\theta} \|w\|_{L^{q+1}(\mathbb{S}^d)}^{1-\theta} \geq C_{\text{GN}} \|w\|_{L^{2q}(\mathbb{S}^d)}$$

if  $1 - \frac{1}{d} \leq p < 1$ . Hint:  $u^{p-1/2} = \frac{w}{\|w\|_{L^{2q}(\mathbb{S}^d)}}, q = \frac{1}{2p-1}$

# The proof

## Lemma

$$I' = \frac{d}{dt} \int_{\mathbb{R}^d} u |\nabla v|^2 dx = -2 \int_{\mathbb{R}^d} u^p \left( \|D^2 v\|^2 + (p-1) (\Delta v)^2 \right) dx$$

$$\|D^2 v\|^2 = \frac{1}{d} (\Delta v)^2 + \left\| D^2 v - \frac{1}{d} \Delta v \text{Id} \right\|^2$$

$$\begin{aligned} \frac{1}{\sigma(1-p)} E^{2-\sigma} (E^\sigma)'' &= (1-p)(\sigma-1) \left( \int_{\mathbb{R}^d} u |\nabla v|^2 dx \right)^2 \\ &\quad - 2 \left( \frac{1}{d} + p - 1 \right) \int_{\mathbb{R}^d} u^p dx \int_{\mathbb{R}^d} u^p (\Delta v)^2 dx \\ &\quad - 2 \int_{\mathbb{R}^d} u^p dx \int_{\mathbb{R}^d} u^p \left\| D^2 v - \frac{1}{d} \Delta v \text{Id} \right\|^2 dx \end{aligned}$$

# Improved rates and best matching



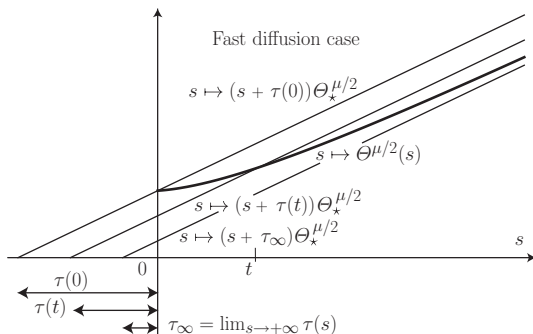
# Temperature (fast diffusion case)

The *second moment functional* (temperature) is defined by

$$\Theta(t) := \frac{1}{d} \int_{\mathbb{R}^d} |x|^2 u(t, x) dx$$

and such that

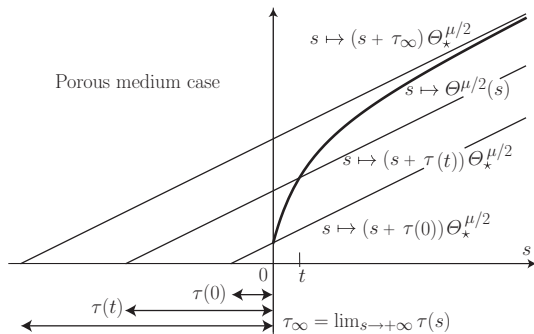
$$\Theta' = 2E$$



# Temperature (porous medium case) and delay

Let  $\mathcal{U}_\star^s$  be the *best matching Barenblatt* function, in the sense of relative entropy  $\mathcal{F}[u | \mathcal{U}_\star^s]$ , among all Barenblatt functions  $(\mathcal{U}_\star^s)_{s>0}$ . We define  $s$  as a function of  $t$  and consider the *delay* given by

$$\tau(t) := \left( \frac{\Theta(t)}{\Theta_\star} \right)^{\frac{\mu}{2}} - t$$



# A result on delays

## Theorem

Assume that  $p \geq 1 - \frac{1}{d}$  and  $p \neq 1$ . The best matching Barenblatt function of a solution  $u$  is  $(t, x) \mapsto \mathcal{U}_*(t + \tau(t), x)$  and the function  $t \mapsto \tau(t)$  is nondecreasing if  $p > 1$  and nonincreasing if  $1 - \frac{1}{d} \leq p < 1$

With  $G := \Theta^{1-\frac{\eta}{2}}$ ,  $\eta = d(1-p) = 2 - \mu$ , the Rényi entropy power functional  $H := \Theta^{-\frac{\eta}{2}} E$  is such that

$$G' = \mu H \quad \text{with} \quad H := \Theta^{-\frac{\eta}{2}} E$$

$$\frac{H'}{1-p} = \Theta^{-1-\frac{\eta}{2}} (\Theta I - d E^2) = \frac{d E^2}{\Theta^{\frac{\eta}{2}+1}} (q - 1) \quad \text{with} \quad q := \frac{\Theta I}{d E^2} \geq 1$$

$$\begin{aligned} d E^2 &= \frac{1}{d} \left( - \int_{\mathbb{R}^d} x \cdot \nabla(u^p) dx \right)^2 = \frac{1}{d} \left( \int_{\mathbb{R}^d} x \cdot u \nabla v dx \right)^2 \\ &\leq \frac{1}{d} \int_{\mathbb{R}^d} u |x|^2 dx \int_{\mathbb{R}^d} u |\nabla v|^2 dx = \Theta I \end{aligned}$$

# An estimate of the delay

## Theorem

*If  $p > 1 - \frac{1}{d}$  and  $p \neq 1$ , then the delay satisfies*

$$\lim_{t \rightarrow +\infty} |\tau(t) - \tau(0)| \geq |1 - p| \frac{\Theta(0)^{1 - \frac{d}{2}(1-p)}}{2 H_\star} \frac{(H_\star - H(0))^2}{\Theta(0) I(0) - d E(0)^2}$$

# C. Fast diffusion equations on manifolds and sharp functional inequalities

- 1 The sphere
- 2 The line
- 3 Compact Riemannian manifolds
- 4 The cylinder: Caffarelli-Kohn-Nirenberg inequalities

# Interpolation inequalities on the sphere

Joint work with M.J. Esteban, M. Kowalczyk and M. Loss

# A family of interpolation inequalities on the sphere

The following interpolation inequality holds on the sphere

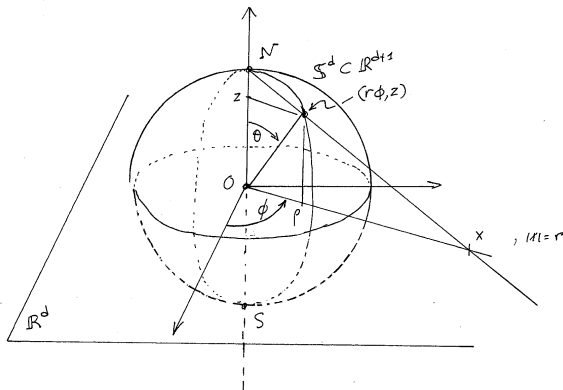
$$\frac{p-2}{d} \int_{\mathbb{S}^d} |\nabla u|^2 d\nu_g + \int_{\mathbb{S}^d} |u|^2 d\nu_g \geq \left( \int_{\mathbb{S}^d} |u|^p d\nu_g \right)^{2/p} \quad \forall u \in H^1(\mathbb{S}^d, d\nu_g)$$

- for any  $p \in (2, 2^*]$  with  $2^* = \frac{2d}{d-2}$  if  $d \geq 3$
- for any  $p \in (2, \infty)$  if  $d = 2$

Here  $d\nu_g$  is the uniform probability measure:  $\nu_g(\mathbb{S}^d) = 1$

- 1 is the optimal constant, equality achieved by constants
- $p = 2^*$  corresponds to Sobolev's inequality...

# Stereographic projection





# Sobolev inequality

The stereographic projection of  $\mathbb{S}^d \subset \mathbb{R}^d \times \mathbb{R} \ni (\rho \phi, z)$  onto  $\mathbb{R}^d$ :  
 to  $\rho^2 + z^2 = 1$ ,  $z \in [-1, 1]$ ,  $\rho \geq 0$ ,  $\phi \in \mathbb{S}^{d-1}$  we associate  $x \in \mathbb{R}^d$  such  
 that  $r = |x|$ ,  $\phi = \frac{x}{|x|}$

$$z = \frac{r^2 - 1}{r^2 + 1} = 1 - \frac{2}{r^2 + 1}, \quad \rho = \frac{2r}{r^2 + 1}$$

and transform any function  $u$  on  $\mathbb{S}^d$  into a function  $v$  on  $\mathbb{R}^d$  using

$$u(y) = \left(\frac{r}{\rho}\right)^{\frac{d-2}{2}} v(x) = \left(\frac{r^2+1}{2}\right)^{\frac{d-2}{2}} v(x) = (1-z)^{-\frac{d-2}{2}} v(x)$$

•  $p = 2^*$ ,  $S_d = \frac{1}{4} d(d-2) |\mathbb{S}^d|^{2/d}$ : Euclidean Sobolev inequality

$$\int_{\mathbb{R}^d} |\nabla v|^2 dx \geq S_d \left[ \int_{\mathbb{R}^d} |v|^{\frac{2d}{d-2}} dx \right]^{\frac{d-2}{d}} \quad \forall v \in \mathcal{D}^{1,2}(\mathbb{R}^d)$$

# Extended inequality

$$\int_{\mathbb{S}^d} |\nabla u|^2 d\nu_g \geq \frac{d}{p-2} \left[ \left( \int_{\mathbb{S}^d} |u|^p d\nu_g \right)^{2/p} - \int_{\mathbb{S}^d} |u|^2 d\nu_g \right] \quad \forall u \in H^1(\mathbb{S}^d, d\mu)$$

is valid

- for any  $p \in (1, 2) \cup (2, \infty)$  if  $d = 1, 2$
- for any  $p \in (1, 2) \cup (2, 2^*]$  if  $d \geq 3$

● Case  $p = 2$ : Logarithmic Sobolev inequality

$$\int_{\mathbb{S}^d} |\nabla u|^2 d\nu_g \geq \frac{d}{2} \int_{\mathbb{S}^d} |u|^2 \log \left( \frac{|u|^2}{\int_{\mathbb{S}^d} |u|^2 d\nu_g} \right) d\nu_g \quad \forall u \in H^1(\mathbb{S}^d, d\mu)$$

● Case  $p = 1$ : Poincaré inequality

$$\int_{\mathbb{S}^d} |\nabla u|^2 d\nu_g \geq d \int_{\mathbb{S}^d} |u - \bar{u}|^2 d\nu_g \quad \text{with} \quad \bar{u} := \int_{\mathbb{S}^d} u d\nu_g \quad \forall u \in H^1(\mathbb{S}^d, d\mu)$$

## Optimality: a perturbation argument

For any  $p \in (1, 2^*]$  if  $d \geq 3$ , any  $p > 1$  if  $d = 1$  or  $2$ , it is remarkable that

$$Q[u] := \frac{(p-2) \|\nabla u\|_{L^2(\mathbb{S}^d)}^2}{\|u\|_{L^p(\mathbb{S}^d)}^2 - \|u\|_{L^2(\mathbb{S}^d)}^2} \geq \inf_{u \in H^1(\mathbb{S}^d, d\mu)} Q[u] = \frac{1}{d}$$

is achieved in the limiting case

$$Q[1 + \varepsilon v] \sim \frac{\|\nabla v\|_{L^2(\mathbb{S}^d)}^2}{\|v\|_{L^2(\mathbb{S}^d)}^2} \quad \text{as } \varepsilon \rightarrow 0$$

when  $v$  is an eigenfunction associated with the first nonzero eigenvalue of  $\Delta_g$ , thus proving the optimality

$p < 2$ : a proof by semi-groups using Nelson's hypercontractivity lemma.  $p > 2$ : no simple proof based on spectral analysis is available: [Beckner], an approach based on Lieb's duality, the Funk-Hecke formula and some (non-trivial) computations

elliptic methods /  $\Gamma_2$  formalism of Bakry-Emery / nonlinear flows

# Schwarz symmetrization and the ultraspherical setting

$$(\xi_0, \xi_1, \dots, \xi_d) \in \mathbb{S}^d, \xi_d = z, \sum_{i=0}^d |\xi_i|^2 = 1 \text{ [Smets-Willem]}$$

## Lemma

*Up to a rotation, any minimizer of  $Q$  depends only on  $\xi_d = z$*

- Let  $d\sigma(\theta) := \frac{(\sin \theta)^{d-1}}{Z_d} d\theta$ ,  $Z_d := \sqrt{\pi} \frac{\Gamma(\frac{d}{2})}{\Gamma(\frac{d+1}{2})}$ :  $\forall v \in H^1([0, \pi], d\sigma)$

$$\frac{p-2}{d} \int_0^\pi |v'(\theta)|^2 d\sigma + \int_0^\pi |v(\theta)|^2 d\sigma \geq \left( \int_0^\pi |v(\theta)|^p d\sigma \right)^{\frac{2}{p}}$$

- Change of variables  $z = \cos \theta$ ,  $v(\theta) = f(z)$

$$\frac{p-2}{d} \int_{-1}^1 |f'|^2 \nu d\nu_d + \int_{-1}^1 |f|^2 d\nu_d \geq \left( \int_{-1}^1 |f|^p d\nu_d \right)^{\frac{2}{p}}$$

where  $\nu_d(z) dz = d\nu_d(z) := Z_d^{-1} \nu^{\frac{d}{2}-1} dz$ ,  $\nu(z) := 1 - z^2$

# The ultraspherical operator

With  $d\nu_d = Z_d^{-1} \nu^{\frac{d}{2}-1} dz$ ,  $\nu(z) := 1 - z^2$ , consider the space  $L^2((-1, 1), d\nu_d)$  with scalar product

$$\langle f_1, f_2 \rangle = \int_{-1}^1 f_1 f_2 d\nu_d, \quad \|f\|_p = \left( \int_{-1}^1 f^p d\nu_d \right)^{\frac{1}{p}}$$

The self-adjoint *ultraspherical* operator is

$$\mathcal{L}f := (1 - z^2) f'' - dz f' = \nu f'' + \frac{d}{2} \nu' f'$$

which satisfies  $\langle f_1, \mathcal{L} f_2 \rangle = - \int_{-1}^1 f_1' f_2' \nu d\nu_d$

## Proposition

Let  $p \in [1, 2) \cup (2, 2^*]$ ,  $d \geq 1$

$$-\langle f, \mathcal{L} f \rangle = \int_{-1}^1 |f'|^2 \nu d\nu_d \geq d \frac{\|f\|_p^2 - \|f\|_2^2}{p - 2} \quad \forall f \in H^1([-1, 1], d\nu_d)$$

# Flows on the sphere

- Heat flow and the Bakry-Emery method
- Fast diffusion (porous media) flow and the choice of the exponents

Joint work with M.J. Esteban, M. Kowalczyk and M. Loss

# Heat flow and the Bakry-Emery method

With  $g = f^p$ , i.e.  $f = g^\alpha$  with  $\alpha = 1/p$

$$(\text{Ineq.}) \quad -\langle f, \mathcal{L} f \rangle = -\langle g^\alpha, \mathcal{L} g^\alpha \rangle =: \mathcal{I}[g] \geq d \frac{\|g\|_1^{2\alpha} - \|g^{2\alpha}\|_1}{p-2} =: \mathcal{F}[g]$$

Heat flow

$$\frac{\partial g}{\partial t} = \mathcal{L} g$$

$$\frac{d}{dt} \|g\|_1 = 0, \quad \frac{d}{dt} \|g^{2\alpha}\|_1 = -2(p-2) \langle f, \mathcal{L} f \rangle = 2(p-2) \int_{-1}^1 |f'|^2 \nu \, d\nu_d$$

which finally gives

$$\frac{d}{dt} \mathcal{F}[g(t, \cdot)] = -\frac{d}{p-2} \frac{d}{dt} \|g^{2\alpha}\|_1 = -2d \mathcal{I}[g(t, \cdot)]$$

$$\text{Ineq.} \iff \frac{d}{dt} \mathcal{F}[g(t, \cdot)] \leq -2d \mathcal{F}[g(t, \cdot)] \iff \frac{d}{dt} \mathcal{I}[g(t, \cdot)] \leq -2d \mathcal{I}[g(t, \cdot)]$$

The equation for  $g = f^p$  can be rewritten in terms of  $f$  as

$$\frac{\partial f}{\partial t} = \mathcal{L} f + (p-1) \frac{|f'|^2}{f} \nu$$

$$-\frac{1}{2} \frac{d}{dt} \int_{-1}^1 |f'|^2 \nu \, d\nu_d = \frac{1}{2} \frac{d}{dt} \langle f, \mathcal{L} f \rangle = \langle \mathcal{L} f, \mathcal{L} f \rangle + (p-1) \left\langle \frac{|f'|^2}{f} \nu, \mathcal{L} f \right\rangle$$

$$\begin{aligned} \frac{d}{dt} \mathcal{I}[g(t, \cdot)] + 2d \mathcal{I}[g(t, \cdot)] &= \frac{d}{dt} \int_{-1}^1 |f'|^2 \nu \, d\nu_d + 2d \int_{-1}^1 |f'|^2 \nu \, d\nu_d \\ &= -2 \int_{-1}^1 \left( |f''|^2 + (p-1) \frac{d}{d+2} \frac{|f'|^4}{f^2} - 2(p-1) \frac{d-1}{d+2} \frac{|f'|^2 f''}{f} \right) \nu^2 \, d\nu_d \end{aligned}$$

is nonpositive if

$$|f''|^2 + (p-1) \frac{d}{d+2} \frac{|f'|^4}{f^2} - 2(p-1) \frac{d-1}{d+2} \frac{|f'|^2 f''}{f}$$

is pointwise nonnegative, which is granted if

$$\left[ (p-1) \frac{d-1}{d+2} \right]^2 \leq (p-1) \frac{d}{d+2} \iff p \leq \frac{2d^2+1}{(d-1)^2} = 2^\# < \frac{2d}{d-2} = 2^*$$



... up to the critical exponent: a proof in two slides

$$\left[ \frac{d}{dz}, \mathcal{L} \right] u = (\mathcal{L} u)' - \mathcal{L} u' = -2z u'' - d u'$$

$$\begin{aligned} \int_{-1}^1 (\mathcal{L} u)^2 d\nu_d &= \int_{-1}^1 |u''|^2 \nu^2 d\nu_d + d \int_{-1}^1 |u'|^2 \nu d\nu_d \\ \int_{-1}^1 (\mathcal{L} u) \frac{|u'|^2}{u} \nu d\nu_d &= \frac{d}{d+2} \int_{-1}^1 \frac{|u'|^4}{u^2} \nu^2 d\nu_d - 2 \frac{d-1}{d+2} \int_{-1}^1 \frac{|u'|^2 u''}{u} \nu^2 d\nu_d \end{aligned}$$

On  $(-1, 1)$ , let us consider the *porous medium (fast diffusion)* flow

$$u_t = u^{2-2\beta} \left( \mathcal{L} u + \kappa \frac{|u'|^2}{u} \nu \right)$$

If  $\kappa = \beta(p-2) + 1$ , the  $L^p$  norm is conserved

$$\frac{d}{dt} \int_{-1}^1 u^{\beta p} d\nu_d = \beta p (\kappa - \beta(p-2) - 1) \int_{-1}^1 u^{\beta(p-2)} |u'|^2 \nu d\nu_d = 0$$

$$f = u^\beta, \|f'\|_{L^2(\mathbb{S}^d)}^2 + \frac{d}{p-2} \left( \|f\|_{L^2(\mathbb{S}^d)}^2 - \|f\|_{L^p(\mathbb{S}^d)}^2 \right) \geq 0 ?$$

$$\begin{aligned} \mathcal{A} := \int_{-1}^1 |u''|^2 \nu^2 d\nu_d - 2 \frac{d-1}{d+2} (\kappa + \beta - 1) \int_{-1}^1 u'' \frac{|u'|^2}{u} \nu^2 d\nu_d \\ + \left[ \kappa(\beta - 1) + \frac{d}{d+2} (\kappa + \beta - 1) \right] \int_{-1}^1 \frac{|u'|^4}{u^2} \nu^2 d\nu_d \end{aligned}$$

$\mathcal{A}$  is nonnegative for some  $\beta$  if

$$\frac{8d^2}{(d+2)^2} (p-1)(2^* - p) \geq 0$$

$\mathcal{A}$  is a sum of squares if  $p \in (2, 2^*)$  for an arbitrary choice of  $\beta$  in a certain interval (depending on  $p$  and

$$\mathcal{A} = \int_{-1}^1 \left| u'' - \frac{p+2}{6-p} \frac{|u'|^2}{u} \right|^2 \nu^2 d\nu_d \geq 0 \quad \text{if } p = 2^* \text{ and } \beta = \frac{4}{6-p}$$

# The rigidity point of view

Which computation have we done ?  $u_t = u^{2-2\beta} \left( \mathcal{L} u + \kappa \frac{|u'|^2}{u} \nu \right)$

$$- \mathcal{L} u - (\beta - 1) \frac{|u'|^2}{u} \nu + \frac{\lambda}{p-2} u = \frac{\lambda}{p-2} u^\kappa$$

Multiply by  $\mathcal{L} u$  and integrate

$$\dots \int_{-1}^1 \mathcal{L} u u^\kappa d\nu_d = - \kappa \int_{-1}^1 u^\kappa \frac{|u'|^2}{u} d\nu_d$$

Multiply by  $\kappa \frac{|u'|^2}{u}$  and integrate

$$\dots = + \kappa \int_{-1}^1 u^\kappa \frac{|u'|^2}{u} d\nu_d$$

The two terms cancel and we are left only with the two-homogenous terms

# Improvements of the inequalities (subcritical range)

- as long as the exponent is either in the range  $(1, 2)$  or in the range  $(2, 2^*)$ , one can establish *improved inequalities*
- An improvement automatically gives an explicit stability result of the optimal functions in the (non-improved) inequality
- By duality, this provides a stability result for Keller-Lieb-Tirring inequalities

# What does “improvement” mean ?

An *improved* inequality is

$$d \Phi(e) \leq i \quad \forall u \in H^1(\mathbb{S}^d) \quad \text{s.t.} \quad \|u\|_{L^2(\mathbb{S}^d)}^2 = 1$$

for some function  $\Phi$  such that  $\Phi(0) = 0$ ,  $\Phi'(0) = 1$ ,  $\Phi' > 0$  and  $\Phi(s) > s$  for any  $s$ . With  $\Psi(s) := s - \Phi^{-1}(s)$

$$i - de \geq d (\Psi \circ \Phi)(e) \quad \forall u \in H^1(\mathbb{S}^d) \quad \text{s.t.} \quad \|u\|_{L^2(\mathbb{S}^d)}^2 = 1$$

## Lemma (Generalized Csiszár-Kullback inequalities)

$$\begin{aligned} & \|\nabla u\|_{L^2(\mathbb{S}^d)}^2 - \frac{d}{p-2} \left[ \|u\|_{L^p(\mathbb{S}^d)}^2 - \|u\|_{L^2(\mathbb{S}^d)}^2 \right] \\ & \geq d \|u\|_{L^2(\mathbb{S}^d)}^2 (\Psi \circ \Phi) \left( C \frac{\|u\|_{L^s(\mathbb{S}^d)}^{2(1-r)}}{\|u\|_{L^2(\mathbb{S}^d)}^2} \|u^r - \bar{u}^r\|_{L^q(\mathbb{S}^d)}^2 \right) \quad \forall u \in H^1(\mathbb{S}^d) \end{aligned}$$

$s(p) := \max\{2, p\}$  and  $p \in (1, 2)$ :  $q(p) := 2/p$ ,  $r(p) := p$ ;  $p \in (2, 4)$ :  
 $q = p/2$ ,  $r = 2$ ;  $p \geq 4$ :  $q = p/(p-2)$ ,  $r = p-2$

# Linear flow: improved Bakry-Emery method

Cf. [Arnold, JD]

$$w_t = \mathcal{L} w + \kappa \frac{|w'|^2}{w} \nu$$

With  $2^\# := \frac{2d+1}{(d-1)^2}$

$$\gamma_1 := \left( \frac{d-1}{d+2} \right)^2 (p-1)(2^\# - p) \quad \text{if } d > 1, \quad \gamma_1 := \frac{p-1}{3} \quad \text{if } d = 1$$

If  $p \in [1, 2) \cup (2, 2^\#]$  and  $w$  is a solution, then

$$\frac{d}{dt} (\mathbf{i} - d \mathbf{e}) \leq -\gamma_1 \int_{-1}^1 \frac{|w'|^4}{w^2} d\nu_d \leq -\gamma_1 \frac{|\mathbf{e}'|^2}{1 - (p-2)\mathbf{e}}$$

Recalling that  $\mathbf{e}' = -\mathbf{i}$ , we get a differential inequality

$$\mathbf{e}'' + d \mathbf{e}' \geq \gamma_1 \frac{|\mathbf{e}'|^2}{1 - (p-2)\mathbf{e}}$$

After integration:  $d \Phi(\mathbf{e}(0)) \leq \mathbf{i}(0)$

# Nonlinear flow: the Hölder estimate of J. Demange

$$w_t = w^{2-2\beta} \left( \mathcal{L} w + \kappa \frac{|w'|^2}{w} \right)$$

For all  $p \in [1, 2^*]$ ,  $\kappa = \beta(p-2) + 1$ ,  $\frac{d}{dt} \int_{-1}^1 w^{\beta p} d\nu_d = 0$

$$-\frac{1}{2\beta^2} \frac{d}{dt} \int_{-1}^1 \left( |(w^\beta)'|^2 \nu + \frac{d}{p-2} (w^{2\beta} - \overline{w}^{2\beta}) \right) d\nu_d \geq \gamma \int_{-1}^1 \frac{|w'|^4}{w^2} \nu^2 d\nu_d$$

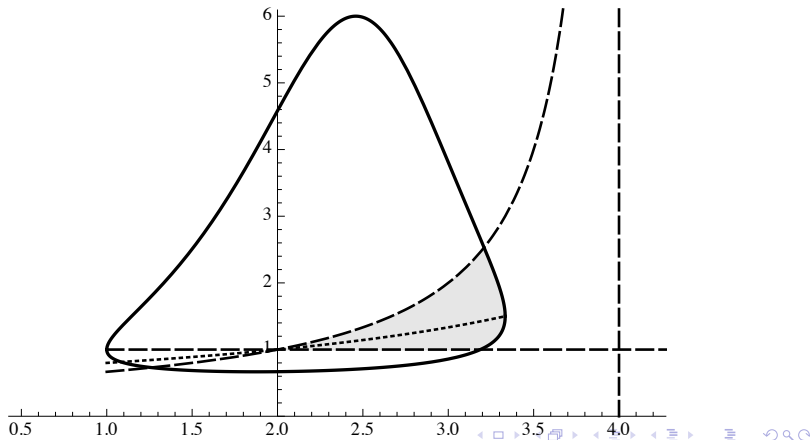
## Lemma

For all  $w \in H^1((-1, 1), d\nu_d)$ , such that  $\int_{-1}^1 w^{\beta p} d\nu_d = 1$

$$\int_{-1}^1 \frac{|w'|^4}{w^2} \nu^2 d\nu_d \geq \frac{1}{\beta^2} \frac{\int_{-1}^1 |(w^\beta)'|^2 \nu d\nu_d \int_{-1}^1 |w'|^2 \nu d\nu_d}{\left( \int_{-1}^1 w^{2\beta} d\nu_d \right)^\delta}$$

.... but there are conditions on  $\beta$

## Admissible $(p, \beta)$ for $d = 5$





# The line

🟢 A first example of a non-compact manifold

Joint work with M.J. Esteban, A. Laptev and M. Loss

# One-dimensional Gagliardo-Nirenberg-Sobolev inequalities

$$\|f\|_{L^p(\mathbb{R})} \leq C_{GN}(p) \|f'\|_{L^2(\mathbb{R})}^\theta \|f\|_{L^2(\mathbb{R})}^{1-\theta} \quad \text{if } p \in (2, \infty)$$

$$\|f\|_{L^2(\mathbb{R})} \leq C_{GN}(p) \|f'\|_{L^2(\mathbb{R})}^\eta \|f\|_{L^p(\mathbb{R})}^{1-\eta} \quad \text{if } p \in (1, 2)$$

with  $\theta = \frac{p-2}{2p}$  and  $\eta = \frac{2-p}{2+p}$

The threshold case corresponding to the limit as  $p \rightarrow 2$  is the logarithmic Sobolev inequality

$$\int_{\mathbb{R}} u^2 \log \left( \frac{u^2}{\|u\|_{L^2(\mathbb{R})}^2} \right) dx \leq \frac{1}{2} \|u\|_{L^2(\mathbb{R})}^2 \log \left( \frac{2}{\pi e} \frac{\|u'\|_{L^2(\mathbb{R})}^2}{\|u\|_{L^2(\mathbb{R})}^2} \right)$$

If  $p > 2$ ,  $u_*(x) = (\cosh x)^{-\frac{2}{p-2}}$  solves

$$-(p-2)^2 u'' + 4u - 2p|u|^{p-2}u = 0$$

If  $p \in (1, 2)$  consider  $u_*(x) = (\cos x)^{\frac{2}{2-p}}$ ,  $x \in (-\pi/2, \pi/2)$

# Flow

Let us define on  $H^1(\mathbb{R})$  the functional

$$\mathcal{F}[v] := \|v'\|_{L^2(\mathbb{R})}^2 + \frac{4}{(p-2)^2} \|v\|_{L^2(\mathbb{R})}^2 - C \|v\|_{L^p(\mathbb{R})}^2 \quad \text{s.t. } \mathcal{F}[u_\star] = 0$$

With  $z(x) := \tanh x$ , consider the *flow*

$$v_t = \frac{v^{1-\frac{p}{2}}}{\sqrt{1-z^2}} \left[ v'' + \frac{2p}{p-2} z v' + \frac{p}{2} \frac{|v'|^2}{v} + \frac{2}{p-2} v \right]$$

## Theorem (Dolbeault-Esteban-Laptev-Loss)

Let  $p \in (2, \infty)$ . Then

$$\frac{d}{dt} \mathcal{F}[v(t)] \leq 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \mathcal{F}[v(t)] = 0$$

$$\frac{d}{dt} \mathcal{F}[v(t)] = 0 \quad \Longleftrightarrow \quad v_0(x) = u_\star(x - x_0)$$

Similar results for  $p \in (1, 2)$

# The inequality ( $p > 2$ ) and the ultraspherical operator

🟢 *The problem on the line is equivalent to the critical problem for the ultraspherical operator*

$$\int_{\mathbb{R}} |v'|^2 dx + \frac{4}{(p-2)^2} \int_{\mathbb{R}} |v|^2 dx \geq C \left( \int_{\mathbb{R}} |v|^p dx \right)^{\frac{2}{p}}$$

With

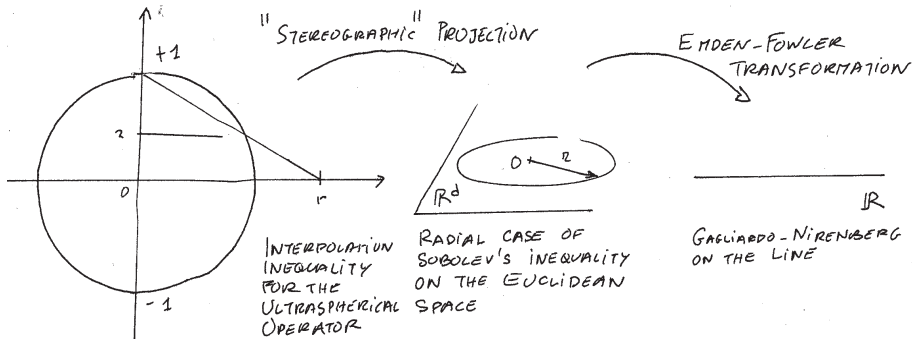
$$z(x) = \tanh x, \quad v_{\star} = (1 - z^2)^{\frac{1}{p-2}} \quad \text{and} \quad v(x) = v_{\star}(x) f(z(x))$$

equality is achieved for  $f = 1$  and, if we let  $\nu(z) := 1 - z^2$ , then

$$\int_{-1}^1 |f'|^2 \nu d\nu_d + \frac{2p}{(p-2)^2} \int_{-1}^1 |f|^2 d\nu_d \geq \frac{2p}{(p-2)^2} \left( \int_{-1}^1 |f|^p d\nu_d \right)^{\frac{2}{p}}$$

where  $d\nu_p$  denotes the probability measure  $d\nu_p(z) := \frac{1}{\zeta_p} \nu^{\frac{2}{p-2}} dz$

$$d = \frac{2p}{p-2} \quad \Longleftrightarrow \quad p = \frac{2d}{d-2}$$



Change of variables = stereographic projection + Emden-Fowler

# Compact Riemannian manifolds

- no sign is required on the Ricci tensor and an improved integral criterion is established
- the flow explores the energy landscape... and shows the non-optimality of the improved criterion

# Riemannian manifolds with positive curvature

$(\mathfrak{M}, g)$  is a smooth closed compact connected Riemannian manifold dimension  $d$ , no boundary,  $\Delta_g$  is the Laplace-Beltrami operator  $\text{vol}(\mathfrak{M}) = 1$ ,  $\mathfrak{R}$  is the Ricci tensor,  $\lambda_1 = \lambda_1(-\Delta_g)$

$$\rho := \inf_{\mathfrak{M}} \inf_{\xi \in \mathbb{S}^{d-1}} \mathfrak{R}(\xi, \xi)$$

## Theorem (Licois-Véron, Bakry-Ledoux)

Assume  $d \geq 2$  and  $\rho > 0$ . If

$$\lambda \leq (1 - \theta) \lambda_1 + \theta \frac{d \rho}{d - 1} \quad \text{where} \quad \theta = \frac{(d - 1)^2 (p - 1)}{d(d + 2) + p - 1} > 0$$

then for any  $p \in (2, 2^*)$ , the equation

$$-\Delta_g v + \frac{\lambda}{p - 2} (v - v^{p-1}) = 0$$

has a unique positive solution  $v \in C^2(\mathfrak{M})$ :  $v \equiv 1$

# Riemannian manifolds: first improvement

## Theorem (Dolbeault-Esteban-Loss)

For any  $p \in (1, 2) \cup (2, 2^*)$

$$0 < \lambda < \lambda_\star = \inf_{u \in H^2(\mathfrak{M})} \frac{\int_{\mathfrak{M}} \left[ (1 - \theta) (\Delta_g u)^2 + \frac{\theta d}{d-1} \Re(\nabla u, \nabla u) \right] d v_g}{\int_{\mathfrak{M}} |\nabla u|^2 d v_g}$$

there is a unique positive solution in  $C^2(\mathfrak{M})$ :  $u \equiv 1$

$\lim_{p \rightarrow 1_+} \theta(p) = 0 \implies \lim_{p \rightarrow 1_+} \lambda_\star(p) = \lambda_1$  if  $\rho$  is bounded  
 $\lambda_\star = \lambda_1 = d \rho / (d-1) = d$  if  $\mathfrak{M} = \mathbb{S}^d$  since  $\rho = d-1$

$$(1 - \theta) \lambda_1 + \theta \frac{d \rho}{d-1} \leq \lambda_\star \leq \lambda_1$$



# Riemannian manifolds: second improvement

$H_g u$  denotes Hessian of  $u$  and  $\theta = \frac{(d-1)^2(p-1)}{d(d+2)+p-1}$

$$Q_g u := H_g u - \frac{g}{d} \Delta_g u - \frac{(d-1)(p-1)}{\theta(d+3-p)} \left[ \frac{\nabla u \otimes \nabla u}{u} - \frac{g}{d} \frac{|\nabla u|^2}{u} \right]$$

$$\Lambda_\star := \inf_{u \in H^2(\mathfrak{M}) \setminus \{0\}} \frac{(1-\theta) \int_{\mathfrak{M}} (\Delta_g u)^2 d\nu_g + \frac{\theta d}{d-1} \int_{\mathfrak{M}} [\|Q_g u\|^2 + \Re(\nabla u, \nabla u)]}{\int_{\mathfrak{M}} |\nabla u|^2 d\nu_g}$$

## Theorem (Dolbeault-Esteban-Loss)

Assume that  $\Lambda_\star > 0$ . For any  $p \in (1, 2) \cup (2, 2^*)$ , the equation has a unique positive solution in  $C^2(\mathfrak{M})$  if  $\lambda \in (0, \Lambda_\star)$ :  $u \equiv 1$

# Optimal interpolation inequality

For any  $p \in (1, 2) \cup (2, 2^*)$  or  $p = 2^*$  if  $d \geq 3$

$$\|\nabla v\|_{L^2(\mathfrak{M})}^2 \geq \frac{\lambda}{p-2} \left[ \|v\|_{L^p(\mathfrak{M})}^2 - \|v\|_{L^2(\mathfrak{M})}^2 \right] \quad \forall v \in H^1(\mathfrak{M})$$

## Theorem (Dolbeault-Esteban-Loss)

Assume  $\Lambda_\star > 0$ . The above inequality holds for some  $\lambda = \Lambda \in [\Lambda_\star, \lambda_1]$   
 If  $\Lambda_\star < \lambda_1$ , then the optimal constant  $\Lambda$  is such that

$$\Lambda_\star < \Lambda \leq \lambda_1$$

If  $p = 1$ , then  $\Lambda = \lambda_1$

Using  $u = 1 + \varepsilon \varphi$  as a test function where  $\varphi$  we get  $\lambda \leq \lambda_1$

A minimum of

$$v \mapsto \|\nabla v\|_{L^2(\mathfrak{M})}^2 - \frac{\lambda}{p-2} \left[ \|v\|_{L^p(\mathfrak{M})}^2 - \|v\|_{L^2(\mathfrak{M})}^2 \right]$$

under the constraint  $\|v\|_{L^p(\mathfrak{M})} = 1$  is negative if  $\lambda > \lambda_1$

# The flow

The key tools the flow

$$u_t = u^{2-2\beta} \left( \Delta_g u + \kappa \frac{|\nabla u|^2}{u} \right), \quad \kappa = 1 + \beta(p-2)$$

If  $v = u^\beta$ , then  $\frac{d}{dt} \|v\|_{L^p(\mathfrak{M})} = 0$  and the functional

$$\mathcal{F}[u] := \int_{\mathfrak{M}} |\nabla(u^\beta)|^2 dv_g + \frac{\lambda}{p-2} \left[ \int_{\mathfrak{M}} u^{2\beta} dv_g - \left( \int_{\mathfrak{M}} u^{\beta p} dv_g \right)^{2/p} \right]$$

is monotone decaying

🟢 J. Demange, *Improved Gagliardo-Nirenberg-Sobolev inequalities on manifolds with positive curvature*, J. Funct. Anal., 254 (2008), pp. 593–611. Also see C. Villani, *Optimal Transport, Old and New*

## Elementary observations (1/2)

Let  $d \geq 2$ ,  $u \in C^2(\mathfrak{M})$ , and consider the trace free Hessian

$$L_g u := H_g u - \frac{g}{d} \Delta_g u$$

### Lemma

$$\int_{\mathfrak{M}} (\Delta_g u)^2 d\nu_g = \frac{d}{d-1} \int_{\mathfrak{M}} \|L_g u\|^2 d\nu_g + \frac{d}{d-1} \int_{\mathfrak{M}} \Re(\nabla u, \nabla u) d\nu_g$$

Based on the Bochner-Lichnerovitz-Weitzenböck formula

$$\frac{1}{2} \Delta |\nabla u|^2 = \|H_g u\|^2 + \nabla(\Delta_g u) \cdot \nabla u + \Re(\nabla u, \nabla u)$$

## Elementary observations (2/2)

### Lemma

$$\begin{aligned} \int_{\mathfrak{M}} \Delta_g u \frac{|\nabla u|^2}{u} d v_g \\ = \frac{d}{d+2} \int_{\mathfrak{M}} \frac{|\nabla u|^4}{u^2} d v_g - \frac{2d}{d+2} \int_{\mathfrak{M}} [L_g u] : \left[ \frac{\nabla u \otimes \nabla u}{u} \right] d v_g \end{aligned}$$

### Lemma

$$\int_{\mathfrak{M}} (\Delta_g u)^2 d v_g \geq \lambda_1 \int_{\mathfrak{M}} |\nabla u|^2 d v_g \quad \forall u \in H^2(\mathfrak{M})$$

and  $\lambda_1$  is the optimal constant in the above inequality

# The key estimates

$$\mathcal{G}[u] := \int_{\mathfrak{M}} \left[ \theta (\Delta_g u)^2 + (\kappa + \beta - 1) \Delta_g u \frac{|\nabla u|^2}{u} + \kappa (\beta - 1) \frac{|\nabla u|^4}{u^2} \right] d\nu_g$$

## Lemma

$$\frac{1}{2\beta^2} \frac{d}{dt} \mathcal{F}[u] = - (1 - \theta) \int_{\mathfrak{M}} (\Delta_g u)^2 d\nu_g - \mathcal{G}[u] + \lambda \int_{\mathfrak{M}} |\nabla u|^2 d\nu_g$$

$$Q_g^\theta u := L_g u - \frac{1}{\theta} \frac{d-1}{d+2} (\kappa + \beta - 1) \left[ \frac{\nabla u \otimes \nabla u}{u} - \frac{g}{d} \frac{|\nabla u|^2}{u} \right]$$

## Lemma

$$\mathcal{G}[u] = \frac{\theta d}{d-1} \left[ \int_{\mathfrak{M}} \|Q_g^\theta u\|^2 d\nu_g + \int_{\mathfrak{M}} \Re(\nabla u, \nabla u) d\nu_g \right] - \mu \int_{\mathfrak{M}} \frac{|\nabla u|^4}{u^2} d\nu_g$$

$$\text{with } \mu := \frac{1}{\theta} \left( \frac{d-1}{d+2} \right)^2 (\kappa + \beta - 1)^2 - \kappa (\beta - 1) - (\kappa + \beta - 1) \frac{d}{d+2}$$

# The end of the proof

Assume that  $d \geq 2$ . If  $\theta = 1$ , then  $\mu$  is nonpositive if

$$\beta_-(p) \leq \beta \leq \beta_+(p) \quad \forall p \in (1, 2^*)$$

where  $\beta_{\pm} := \frac{b \pm \sqrt{b^2 - a}}{2a}$  with  $a = 2 - p + \left[ \frac{(d-1)(p-1)}{d+2} \right]^2$  and  $b = \frac{d+3-p}{d+2}$

Notice that  $\beta_-(p) < \beta_+(p)$  if  $p \in (1, 2^*)$  and  $\beta_-(2^*) = \beta_+(2^*)$

$$\theta = \frac{(d-1)^2(p-1)}{d(d+2)+p-1} \quad \text{and} \quad \beta = \frac{d+2}{d+3-p}$$

## Proposition

Let  $d \geq 2$ ,  $p \in (1, 2) \cup (2, 2^*)$  ( $p \neq 5$  or  $d \neq 2$ )

$$\frac{1}{2\beta^2} \frac{d}{dt} \mathcal{F}[u] \leq (\lambda - \Lambda_*) \int_{\mathfrak{M}} |\nabla u|^2 d\nu_g$$

# The Moser-Trudinger-Onofri inequality on Riemannian manifolds

Joint work with G. Jankowiak and M.J. Esteban

📍 Extension to compact Riemannian manifolds of dimension 2...



We shall also denote by  $\mathfrak{R}$  the Ricci tensor, by  $H_g u$  the Hessian of  $u$  and by

$$L_g u := H_g u - \frac{g}{d} \Delta_g u$$

the trace free Hessian. Let us denote by  $M_g u$  the trace free tensor

$$M_g u := \nabla u \otimes \nabla u - \frac{g}{d} |\nabla u|^2$$

We define

$$\lambda_\star := \inf_{u \in H^2(\mathfrak{M}) \setminus \{0\}} \frac{\int_{\mathfrak{M}} \left[ \|L_g u - \frac{1}{2} M_g u\|^2 + \mathfrak{R}(\nabla u, \nabla u) \right] e^{-u/2} d\nu_g}{\int_{\mathfrak{M}} |\nabla u|^2 e^{-u/2} d\nu_g}$$

## Theorem

Assume that  $d = 2$  and  $\lambda_\star > 0$ . If  $u$  is a smooth solution to

$$-\frac{1}{2} \Delta_g u + \lambda = e^u$$

then  $u$  is a constant function if  $\lambda \in (0, \lambda_\star)$

The Moser-Trudinger-Onofri inequality on  $\mathfrak{M}$

$$\frac{1}{4} \|\nabla u\|_{L^2(\mathfrak{M})}^2 + \lambda \int_{\mathfrak{M}} u \, d\nu_g \geq \lambda \log \left( \int_{\mathfrak{M}} e^u \, d\nu_g \right) \quad \forall u \in H^1(\mathfrak{M})$$

for some constant  $\lambda > 0$ . Let us denote by  $\lambda_1$  the first positive eigenvalue of  $-\Delta_g$

## Corollary

If  $d = 2$ , then the MTO inequality holds with  $\lambda = \Lambda := \min\{4\pi, \lambda_\star\}$ . Moreover, if  $\Lambda$  is strictly smaller than  $\lambda_1/2$ , then the optimal constant in the MTO inequality is strictly larger than  $\Lambda$

# The flow

$$\frac{\partial f}{\partial t} = \Delta_g(e^{-f/2}) - \frac{1}{2} |\nabla f|^2 e^{-f/2}$$

$$\begin{aligned} \mathcal{G}_\lambda[f] := \int_{\mathfrak{M}} \|L_g f - \frac{1}{2} M_g f\|^2 e^{-f/2} d\nu_g + \int_{\mathfrak{M}} \Re(\nabla f, \nabla f) e^{-f/2} d\nu_g \\ - \lambda \int_{\mathfrak{M}} |\nabla f|^2 e^{-f/2} d\nu_g \end{aligned}$$

Then for any  $\lambda \leq \lambda_\star$  we have

$$\begin{aligned} \frac{d}{dt} \mathcal{F}_\lambda[f(t, \cdot)] &= \int_{\mathfrak{M}} \left(-\frac{1}{2} \Delta_g f + \lambda\right) \left(\Delta_g(e^{-f/2}) - \frac{1}{2} |\nabla f|^2 e^{-f/2}\right) d\nu_g \\ &= -\mathcal{G}_\lambda[f(t, \cdot)] \end{aligned}$$

Since  $\mathcal{F}_\lambda$  is nonnegative and  $\lim_{t \rightarrow \infty} \mathcal{F}_\lambda[f(t, \cdot)] = 0$ , we obtain that

$$\mathcal{F}_\lambda[u] \geq \int_0^\infty \mathcal{G}_\lambda[f(t, \cdot)] dt$$

# Weighted Moser-Trudinger-Onofri inequalities on the two-dimensional Euclidean space

On the Euclidean space  $\mathbb{R}^2$ , given a general probability measure  $\mu$  does the inequality

$$\frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla u|^2 dx \geq \lambda \left[ \log \left( \int_{\mathbb{R}^2} e^u d\mu \right) - \int_{\mathbb{R}^2} u d\mu \right]$$

hold for some  $\lambda > 0$  ? Let

$$\Lambda_\star := \inf_{x \in \mathbb{R}^2} \frac{-\Delta \log \mu}{8\pi \mu}$$

## Theorem

*Assume that  $\mu$  is a radially symmetric function. Then any radially symmetric solution to the EL equation is a constant if  $\lambda < \Lambda_\star$  and the inequality holds with  $\lambda = \Lambda_\star$  if equality is achieved among radial functions*

# Caffarelli-Kohn-Nirenberg inequalities

Work in progress with M.J. Esteban and M. Loss

# Caffarelli-Kohn-Nirenberg inequalities and the symmetry breaking issue

Let  $\mathcal{D}_{a,b} := \left\{ v \in L^p(\mathbb{R}^d, |x|^{-b} dx) : |x|^{-a} |\nabla v| \in L^2(\mathbb{R}^d, dx) \right\}$

$$\left( \int_{\mathbb{R}^d} \frac{|v|^p}{|x|^{bp}} dx \right)^{2/p} \leq C_{a,b} \int_{\mathbb{R}^d} \frac{|\nabla v|^2}{|x|^{2a}} dx \quad \forall v \in \mathcal{D}_{a,b}$$

hold under the conditions that  $a \leq b \leq a+1$  if  $d \geq 3$ ,  $a < b \leq a+1$  if  $d = 2$ ,  $a + 1/2 < b \leq a+1$  if  $d = 1$ , and  $a < a_c := (d-2)/2$

$$p = \frac{2d}{d-2+2(b-a)}$$

▷ *With*

$$v_\star(x) = \left( 1 + |x|^{(p-2)(a_c-a)} \right)^{-\frac{2}{p-2}} \quad \text{and} \quad C_{a,b}^\star = \frac{\| |x|^{-b} v_\star \|_p^2}{\| |x|^{-a} \nabla v_\star \|_2^2}$$

do we have  $C_{a,b} = C_{a,b}^\star$  (symmetry)  
 or  $C_{a,b} > C_{a,b}^\star$  (symmetry breaking) ?

# The Emden-Fowler transformation and the cylinder

$$v(r, \omega) = r^{a-a_c} \varphi(s, \omega) \quad \text{with} \quad r = |x|, \quad s = -\log r \quad \text{and} \quad \omega = \frac{x}{r}$$

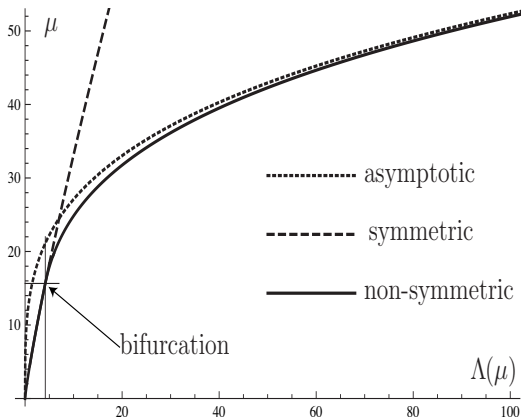
With this transformation, the Caffarelli-Kohn-Nirenberg inequalities can be rewritten as

$$\|\partial_s \varphi\|_{L^2(C_1)}^2 + \|\nabla_\omega \varphi\|_{L^2(C_1)}^2 + \Lambda \|\varphi\|_{L^p(C_1)}^2 \geq \mu(\Lambda) \|\varphi\|_{L^p(C_1)}^2 \quad \forall \varphi \in H^1(\mathcal{C})$$

where  $\Lambda := (a_c - a)^2$ ,  $\mathcal{C} = \mathbb{R} \times \mathbb{S}^{d-1}$  and the optimal constant  $\mu(\Lambda)$  is

$$\mu(\Lambda) = \frac{1}{C_{a,b}} \quad \text{with} \quad a = a_c \pm \sqrt{\Lambda} \quad \text{and} \quad b = \frac{d}{p} \pm \sqrt{\Lambda}$$

# Numerical results



*Parametric plot of the branch of optimal functions for  $p = 2.8$ ,  $d = 5$ ,  $\theta = 1$ . Non-symmetric solutions bifurcate from symmetric ones at a bifurcation point computed by V. Felli and M. Schneider. The branch behaves for large values of  $\Lambda$  as predicted by F. Catrina and Z.-Q. Wang*

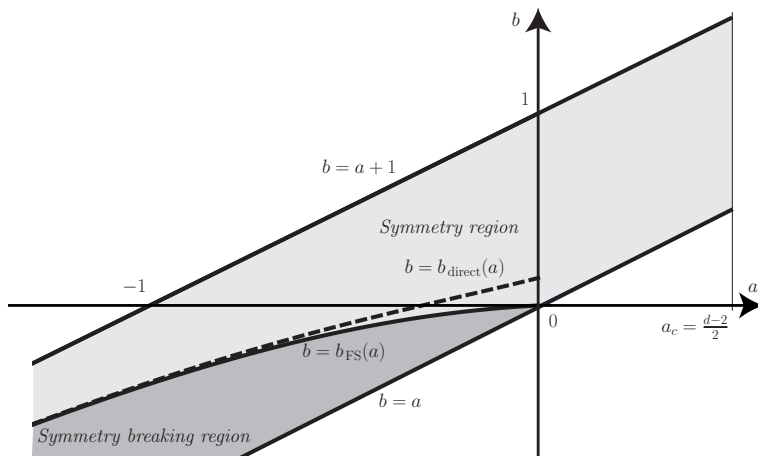


# The symmetry result

$$b_{\text{FS}}(a) := \frac{d(a_c - a)}{2\sqrt{(a_c - a)^2 + d - 1}} + a - a_c$$

## Theorem

*Let  $d \geq 2$  and  $p \leq 4$ . If either  $a \in [0, a_c)$  and  $b > 0$ , or  $a < 0$  and  $b \geq b_{\text{FS}}(a)$ , then the optimal functions for the Caffarelli-Kohn-Nirenberg inequalities are radially symmetric*



*The Felli-Schneider region, or symmetry breaking region, appears in dark grey and is defined by  $a < 0$ ,  $a \leq b < b_{\text{FS}}(a)$ . We prove that symmetry holds in the light grey region defined by  $b \geq b_{\text{FS}}(a)$  when  $a < 0$  and for any  $b \in [a, a + 1]$  if  $a \in [0, a_c]$*

# Sketch of a proof

# A change of variables

With  $(r = |x|, \omega = x/r) \in \mathbb{R}^+ \times \mathbb{S}^{d-1}$ , the Caffarelli-Kohn-Nirenberg inequality is

$$\left( \int_0^\infty \int_{\mathbb{S}^{d-1}} |v|^p r^{d-bp} \frac{dr}{r} d\omega \right)^{\frac{2}{p}} \leq C_{a,b} \int_0^\infty \int_{\mathbb{S}^{d-1}} |\nabla v|^2 r^{d-2a} \frac{dr}{r} d\omega$$

Change of variables  $r \mapsto r^\alpha$ ,  $v(r, \omega) = w(r^\alpha, \omega)$

$$\begin{aligned} \alpha^{1-\frac{2}{p}} \left( \int_0^\infty \int_{\mathbb{S}^{d-1}} |w|^p r^{\frac{d-bp}{\alpha}} \frac{dr}{r} d\omega \right)^{\frac{2}{p}} \\ \leq C_{a,b} \int_0^\infty \int_{\mathbb{S}^{d-1}} \left( \alpha^2 \left| \frac{\partial w}{\partial r} \right|^2 + \frac{1}{r^2} |\nabla_\omega w|^2 \right) r^{\frac{d-2a-2}{\alpha}+2} \frac{dr}{r} d\omega \end{aligned}$$

Choice of  $\alpha$

$$n = \frac{d-bp}{\alpha} = \frac{d-2a-2}{\alpha} + 2$$

Then  $p = \frac{2n}{n-2}$  is the critical Sobolev exponent associated with  $n$

# A Sobolev type inequality

The parameters  $\alpha$  and  $n$  vary in the ranges  $0 < \alpha < \infty$  and  $d < n < \infty$  and the *Felli-Schneider curve* in the  $(\alpha, n)$  variables is given by

$$\alpha = \sqrt{\frac{d-1}{n-1}} =: \alpha_{\text{FS}}$$

With

$$Dw = \left( \alpha \frac{\partial w}{\partial r}, \frac{1}{r} \nabla_{\omega} w \right), \quad d\mu := r^{n-1} dr d\omega$$

the inequality becomes

$$\alpha^{1-\frac{2}{p}} \left( \int_{\mathbb{R}^d} |w|^p d\mu \right)^{\frac{2}{p}} \leq C_{a,b} \int_{\mathbb{R}^d} |Dw|^2 d\mu$$

## Proposition

Let  $d \geq 4$ . Optimality is achieved by radial functions and  $C_{a,b} = C_{a,b}^*$  if  $\alpha \leq \alpha_{\text{FS}}$

Gagliardo-Nirenberg inequalities on general cylinders: similar

# Notations

When there is no ambiguity, we will omit the index  $\omega$  and from now on write that  $\nabla = \nabla_\omega$  denotes the gradient with respect to the angular variable  $\omega \in \mathbb{S}^{d-1}$  and that  $\Delta$  is the Laplace-Beltrami operator on  $\mathbb{S}^{d-1}$ . We define the self-adjoint operator  $\mathcal{L}$  by

$$\mathcal{L} w := -D^* D w = \alpha^2 w'' + \alpha^2 \frac{n-1}{r} w' + \frac{\Delta w}{r^2}$$

The fundamental property of  $\mathcal{L}$  is the fact that

$$\int_{\mathbb{R}^d} w_1 \mathcal{L} w_2 d\mu = - \int_{\mathbb{R}^d} Dw_1 \cdot Dw_2 d\mu \quad \forall w_1, w_2 \in \mathcal{D}(\mathbb{R}^d)$$

▷ Heuristics: we look for a monotonicity formula along a well chosen nonlinear flow, based on the analogy with the decay of the Fisher information along the fast diffusion flow in  $\mathbb{R}^d$

# Fisher information

$$\text{Let } u^{\frac{1}{2}-\frac{1}{n}} = |w| \quad \Longleftrightarrow \quad u = |w|^p, \quad p = \frac{2n}{n-2}$$

$$\mathcal{I}[u] := \int_{\mathbb{R}^d} u |\mathrm{D}p|^2 d\mu, \quad p = \frac{m}{1-m} u^{m-1} \quad \text{and} \quad m = 1 - \frac{1}{n}$$

Here  $\mathcal{I}$  is the *Fisher information* and  $p$  is the *pressure function*

## Proposition

With  $\Lambda = 4\alpha^2/(p-2)^2$  and for some explicit numerical constant  $\kappa$ , we have

$$\kappa \mu(\Lambda) = \inf \left\{ \mathcal{I}[u] : \|u\|_{L^1(\mathbb{S}^d, d\nu_n)} \right\}$$

# The fast diffusion equation

$$\frac{\partial u}{\partial t} = \mathcal{L} u^m, \quad m = 1 - \frac{1}{n}$$

Barenblatt self-similar solutions

$$u_*(t, r, \omega) = t^{-n} \left( c_* + \frac{r^2}{2(n-1)\alpha^2 t^2} \right)^{-n}$$

## Lemma

$$\kappa \mu_*(\Lambda) = \mathcal{I}[u_*(t, \cdot)] \quad \forall t > 0$$

▷ Strategy:

- 1) prove that  $\frac{d}{dt} \mathcal{I}[u(t, \cdot)] \leq 0$ ,
- 2) prove that  $\frac{d}{dt} \mathcal{I}[u(t, \cdot)] = 0$  means that  $u = u_*$  up to a time shift



# Decay of the Fisher information along the flow ?

$$\frac{\partial p}{\partial t} = \frac{1}{n} p \mathcal{L} p - |Dp|^2$$

$$\mathcal{Q}[p] := \frac{1}{2} \mathcal{L} |Dp|^2 - Dp \cdot D\mathcal{L} p$$

$$\mathcal{K}[p] := \int_{\mathbb{R}^d} \left( \mathcal{Q}[p] - \frac{1}{n} (\mathcal{L} p)^2 \right) p^{1-n} d\mu$$

## Lemma

$$\frac{d}{dt} \mathcal{I}[u(t, \cdot)] = -2(n-1)^{n-1} \mathcal{K}[p]$$

If  $u$  is a critical point, then  $\mathcal{K}[p] = 0$

Boundary terms ! Regularity !

## Proving decay (1/2)

$$k[p] := \mathcal{Q}(p) - \frac{1}{n} (\mathcal{L} p)^2 = \frac{1}{2} \mathcal{L} |Dp|^2 - Dp \cdot D \mathcal{L} p - \frac{1}{n} (\mathcal{L} p)^2$$

$$k_{\mathfrak{M}}[p] := \frac{1}{2} \Delta |\nabla p|^2 - \nabla p \cdot \nabla \Delta p - \frac{1}{n-1} (\Delta p)^2 - (n-2) \alpha^2 |\nabla p|^2$$

### Lemma

Let  $n \neq 1$  be any real number,  $d \in \mathbb{N}$ ,  $d \geq 2$ , and consider a function  $p \in C^3((0, \infty) \times \mathfrak{M})$ , where  $(\mathfrak{M}, g)$  is a smooth, compact Riemannian manifold. Then we have

$$k[p] = \alpha^4 \left(1 - \frac{1}{n}\right) \left[ p'' - \frac{p'}{r} - \frac{\Delta p}{\alpha^2 (n-1) r^2} \right]^2 + 2 \alpha^2 \frac{1}{r^2} \left| \nabla p' - \frac{\nabla p}{r} \right|^2 + \frac{1}{r^4} k_{\mathfrak{M}}[p]$$

## Proving decay (2/2)

### Lemma

Assume that  $d \geq 3$ ,  $n > d$  and  $\mathfrak{M} = \mathbb{S}^{d-1}$ . There is a positive constant  $\zeta_\star$  such that

$$\int_{\mathbb{S}^{d-1}} k_{\mathfrak{M}}[p] p^{1-n} d\omega \geq (\lambda_\star - (n-2)\alpha^2) \int_{\mathbb{S}^{d-1}} |\nabla p|^2 p^{1-n} d\omega + \zeta_\star (n-d) \int_{\mathbb{S}^{d-1}} |\nabla p|^4 p^{1-n} d\omega$$

Proof based on the Bochner-Lichnerowicz-Weitzenböck formula

### Corollary

Let  $d \geq 2$  and assume that  $\alpha \leq \alpha_{\text{FS}}$ . Then for any nonnegative function  $u \in L^1(\mathbb{R}^d)$  with  $\mathcal{I}[u] < +\infty$  and  $\int_{\mathbb{R}^d} u d\mu = 1$ , we have

$$\mathcal{I}[u] \geq \mathcal{I}_\star$$

# A perturbation argument

🟢 If  $u$  is a critical point of  $\mathcal{I}$  under the mass constraint  $\int_{\mathbb{R}^d} u \, d\mu = 1$ , then

$$o(\varepsilon) = \mathcal{I}[u + \varepsilon \mathcal{L} u^m] - \mathcal{I}[u] = -2(n-1)^{n-1} \varepsilon \mathcal{K}[p] + o(\varepsilon)$$

because  $\varepsilon \mathcal{L} u^m$  is an admissible perturbation. Indeed, we know that

$$\int_{\mathbb{R}^d} (u + \varepsilon \mathcal{L} u^m) \, d\mu = \int_{\mathbb{R}^d} u \, d\mu = 1$$

and, as we take the limit as  $\varepsilon \rightarrow 0$ ,  $u + \varepsilon \mathcal{L} u^m$  makes sense and, in particular, is positive

🟢 If  $\alpha \leq \alpha_{\text{FS}}$ , then  $\mathcal{K}[p] = 0$  implies that  $u = u_\star$

# A summary

🟢 the sphere: the flow tells us what to do, and provides a simple proof (*choice of the exponents / of the nonlinearity*) once the problem is reduced to the ultraspherical setting + improvements

🟢 [not presented here: Keller-Lieb-Thirring estimates] the spectral point of view on the inequality: how to measure the deviation with respect to the *semi-classical* estimates, a nice example of bifurcation (and *symmetry breaking*)

🟢 *Riemannian manifolds*: no sign is required on the Ricci tensor and an improved integral criterion is established. We extend the theory from pointwise criteria to a non-local Schrödinger type estimate (Rayleigh quotient). The method generically shows the non-optimality of the improved criterion

🟢 the flow is a nice way of exploring an energy space: it explain how to produce a good test function at *any* critical point. A *rigidity* result tells you that a local result is actually global because otherwise the flow would relate (far away) extremal points while keeping the energy minimal

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<http://www.ceremade.dauphine.fr/~dolbeaul>

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▷ Lectures

Thank you for your attention !