

# Singular limit problems related to the Onofri inequality

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*Singular limit problems in nonlinear PDEs*  
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# Weighted Onofri inequality and Caffarelli-Kohn-Nirenberg inequalities in two space dimensions

# Onofri's inequality

[E. Onofri. On the positivity of the effective action in a theory of random surfaces. *Comm. Math. Phys.*, 86 (3): 321-326, 1982]

$$\log \left( \int_{\mathbb{S}^2} e^{2u} d\sigma \right) - 2 \int_{\mathbb{S}^2} u d\sigma \leq \|\nabla u\|_{L^2(\mathbb{S}^2, d\sigma)}^2$$

for all  $u \in \mathcal{E} = \{u \in L^1(\mathbb{S}^2, d\sigma) : |\nabla u| \in L^2(\mathbb{S}^2, d\sigma)\}$

By the stereographic projection from  $\mathbb{S}^2$  onto  $\mathbb{R}^2$ , we get an Onofri type inequality in  $\mathbb{R}^2$

$$\log \left( \int_{\mathbb{R}^2} e^v d\mu \right) - \int_{\mathbb{R}^2} v d\mu \leq \frac{1}{16\pi} \|\nabla v\|_{L^2(\mathbb{R}^2, dx)}^2$$

for all  $v \in \mathcal{D} = \{v \in L^1(\mathbb{R}^2, d\mu) : |\nabla v| \in L^2(\mathbb{R}^2, dx)\}$  and

$$d\mu = \frac{dx}{\pi(1+|x|^2)^2}$$

# A first result: generalized Onofri inequalities

On  $\mathbb{R}^2$  for  $\alpha > -1$ , consider the family of probability measures

$$d\mu_\alpha = \frac{\alpha + 1}{\pi} \frac{|x|^{2\alpha} dx}{(1 + |x|^2)^{(\alpha+1)^2}}$$

Theorem (JD, M. Esteban, G. Tarantello)

$$\log \left( \int_{\mathbb{R}^2} e^v d\mu_\alpha \right) - \int_{\mathbb{R}^2} v d\mu_\alpha \leq \frac{1}{16\pi(\alpha+1)} \|\nabla v\|_{L^2(\mathbb{R}^2, dx)}^2$$

holds in the space  $\mathcal{E}_\alpha = \{v \in L^1(\mathbb{R}^2, d\mu_\alpha) : |\nabla v| \in L^2(\mathbb{R}^2, dx)\}$   
 restricted to radially symmetric functions  $\forall \alpha > -1$ , and without  
 restriction if  $\alpha \in (-1, 0]$

# Proof

## Lemma

Let  $\alpha > -1$ . For all  $v \in \mathcal{E}_\alpha$ , there holds

$$\int_{\mathbb{R}^2} e^{v - \int_{\mathbb{R}^2} v d\mu_\alpha} d\mu_\alpha \leq e^{\frac{1}{16\pi(\alpha+1)} (\|\nabla v\|_2^2 + \alpha(\alpha+2) \|\frac{1}{r} \partial_\theta v\|_2^2)}$$

$\mathbb{C} \approx \mathbb{R}^2 \ni x = r e^{i\theta}$ ,  $r \geq 0$ ,  $\theta \in [0, 2\pi)$ . Stereographic projection:  $\Sigma_0$

Let  $\alpha > -1$  and define the inverse of a dilated stereographic projection

$$\Sigma_\alpha^{-1}(r e^{i\theta}) := \left( \frac{2 r^{\alpha+1} e^{i\theta}}{1 + r^{2(\alpha+1)}}, \frac{r^{2(\alpha+1)} - 1}{1 + r^{2(\alpha+1)}} \right) = \Sigma_0^{-1}(r^{1+\alpha} e^{i\theta})$$

If  $f \in C(\mathbb{R})$ ,  $f(u)$ ,  $|\nabla u|^2 \in L^1(\mathbb{S}^2)$  and  $v = u \circ \Sigma_\alpha^{-1}$ , then

- $\int_{\mathbb{S}^2} f(u) d\sigma = \int_{\mathbb{R}^2} f(v) d\mu_\alpha$
- $4\pi \int_{\mathbb{S}^2} |\nabla u|^2 d\sigma = \frac{1}{\alpha+1} \int_{\mathbb{R}^2} \left( |\nabla v|^2 + \alpha(\alpha+2) \left| \frac{1}{r} \partial_\theta v \right|^2 \right) dx$

The result follows from Onofri's inequality

## Application to Caffarelli-Kohn-Nirenberg inequalities

$$\left( \int_{\mathbb{R}^d} \frac{|u|^p}{|x|^{bp}} dx \right)^{2/p} \leq C_{a,b} \int_{\mathbb{R}^d} \frac{|\nabla u|^2}{|x|^{2a}} dx \quad \forall u \in \mathcal{D}_{a,b}$$

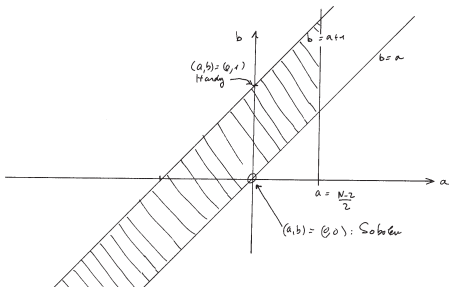
with  $a \leq b \leq a + 1$  if  $d \geq 3$ ,  $a < b \leq a + 1$  if  $d = 2$ , and  $a \neq \frac{d-2}{2} =: a_c$

$$p = \frac{2d}{d - 2 + 2(b - a)}$$

$$\mathcal{D}_{a,b} := \left\{ |x|^{-b} u \in L^p(\mathbb{R}^d, dx) : |x|^{-a} |\nabla u| \in L^2(\mathbb{R}^d, dx) \right\}$$

[Catrina, Wang]  
 [Felli, Schneider]  
 [Smets, Willem]  
 [Lin, Wang]

...



# Approaching Onofri's inequality ( $d = 2$ )

## Theorem (JD, Esteban, Tarantello)

For all  $\varepsilon > 0 \exists \eta > 0$  s.t. for  $a < 0$ ,  $|a| < \eta$

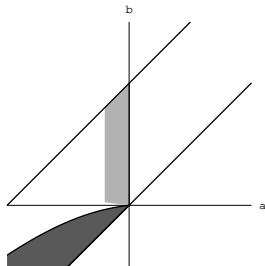
(i) if  $|a| > \frac{2}{p-\varepsilon} (1 + |a|^2)$ , then  $C_{a,b} > C_{a,b}^*$  (symmetry breaking)

(ii) if  $|a| < \frac{2}{p+\varepsilon} (1 + |a|^2)$ , then  $C_{a,b} = C_{a,b}^*$  and  $u_{a,b} = u_{a,b}^*$

For  $d = 2$ , radial symmetry holds if  $-\eta < a < 0$  and  $-\varepsilon(\eta)a \leq b < a + 1$

$$\varepsilon = \frac{2}{p} \alpha = -1 + (1 - \varepsilon) \frac{a}{\varepsilon}$$

Blow-up + Liouville equation





# The Onofri inequality as an endpoint of Gagliardo-Nirenberg inequalities

## Gagliardo-Nirenberg inequalities

Consider the following sub-family of Gagliardo-Nirenberg inequalities

$$\|f\|_{L^{2p}(\mathbb{R}^d)} \leq C_{p,d} \|\nabla f\|_{L^2(\mathbb{R}^d)}^\theta \|f\|_{L^{p+1}(\mathbb{R}^d)}^{1-\theta}$$

with  $\theta = \theta(p) := \frac{p-1}{p} \frac{d}{d+2-p(d-2)}$ ,  $p = \frac{1}{2m-1}$

- $1 < p \leq \frac{d}{d-2}$  if  $d \geq 3$ ,  $\frac{d-1}{d} \leq m < 1$
- $1 < p < \infty$  if  $d = 2$ ,  $\frac{1}{2} < m < 1$

[M. del Pino, JD] equality holds if  $f = F_p$  with

$$F_p(x) = (1 + |x|^2)^{-\frac{1}{p-1}} \quad \forall x \in \mathbb{R}^d$$

All extremal functions are equal to  $F_p$  up to a multiplication by a constant, a translation and a scaling

- When  $p \rightarrow 1$ , we recover the euclidean logarithmic Sobolev inequality in optimal scale invariant form [F. Weissler]
- If  $d \geq 3$ , the limit case  $p = d/(d-2)$  corresponds to Sobolev's inequality [T. Aubin, G. Talenti]
- If  $d = 2$  and  $p \rightarrow \infty$ ...

## Onofri's inequality as a limit case

When  $d = 2$ , Onofri's inequality can be seen as an endpoint case of the family of the Gagliardo-Nirenberg inequalities

### Proposition (JD)

Assume that  $g \in \mathcal{D}(\mathbb{R}^d)$  is such that  $\int_{\mathbb{R}^2} g \, d\mu = 0$  and let

$f_p := F_p \left(1 + \frac{g}{2p}\right)$ . With  $\mu(x) := \frac{1}{\pi} (1 + |x|^2)^{-2}$ , and  $d\mu(x) := \mu(x) \, dx$ , we have

$$1 \leq \lim_{p \rightarrow \infty} C_{p,2} \frac{\|\nabla f_p\|_{L^2(\mathbb{R}^2)}^{\theta(p)} \|f_p\|_{L^{p+1}(\mathbb{R}^2)}^{1-\theta(p)}}{\|f_p\|_{L^{2p}(\mathbb{R}^2)}} = \frac{e^{\frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla g|^2 \, dx}}{\int_{\mathbb{R}^2} e^g \, d\mu}$$

The standard form of the euclidean version of Onofri's inequality is

$$\log \left( \int_{\mathbb{R}^2} e^g \, d\mu \right) - \int_{\mathbb{R}^2} g \, d\mu \leq \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla g|^2 \, dx$$

## The Onofri inequality in higher dimensions ?

On  $\mathbb{R}^d$  with  $d \geq 3$ , let us define

$$Q_d[u] := \frac{\int_{\mathbb{R}^d} H_d(x, \nabla u) dx}{\log\left(\int_{\mathbb{R}^d} e^u d\mu_d\right) - \int_{\mathbb{R}^2} u d\mu}, \quad d\mu_d(x) := \frac{dx}{|\mathbb{S}^{d-1}| \left(1 + |x|^{\frac{d}{d-1}}\right)^d}$$

$$R_d(X, Y) := |X + Y|^d - |X|^d - d|X|^{d-2} X \cdot Y, \quad (X, Y) \in \mathbb{R}^d \times \mathbb{R}^d$$

$$H_d(x, p) := R_d\left(-\frac{d|x|^{-\frac{d-2}{d-1}}}{1+|x|^{\frac{d}{d-1}}} x, \frac{d-1}{d} p\right), \quad (x, p) \in \mathbb{R}^d \times \mathbb{R}^d$$

Theorem (JD, del Pino)

$$\log\left(\int_{\mathbb{R}^d} e^u d\mu_d\right) - \int_{\mathbb{R}^2} u d\mu \leq \alpha_d \int_{\mathbb{R}^d} H_d(x, \nabla u) dx \quad (1)$$

The optimal constant  $\alpha_d$  is explicit and given by

$$\alpha_d = \frac{d^{1-d} \Gamma(d/2)}{2(d-1) \pi^{d/2}}$$

## Comments

Let  $e \in \mathbb{S}^{d-1}$

$$\lim_{\varepsilon \rightarrow 0} Q_d[\varepsilon v] = \frac{1}{\alpha_d} \quad \text{if} \quad v(x) = -d \frac{x \cdot e}{|x|^{\frac{d-2}{d-1}} \left(1 + |x|^{\frac{d}{d-1}}\right)}$$

### Example

- If  $d = 2$ ,  $\int_{\mathbb{R}^d} H_2(x, \nabla u) dx = \frac{1}{4} \int_{\mathbb{R}^2} |\nabla u|^2 dx$ : Onofri's inequality with optimal constant  $1/\alpha_2 = 4\pi$ .
- If  $d = 4$ , we find that  $H_4(x, \nabla u)$  is a fourth order polynomial and  $R_4(X, Y) = 4(X \cdot Y)^2 + |Y|^2(|Y|^2 + 4X \cdot Y + 2|X|^2)$ .

Extensions of the inequality to higher dimensions have been obtained by [Carlen & Loss, 1992] and [Beckner, 1993] in the case of  $\mathbb{S}^d$  but there natural conformally invariant, non-local generalizations of the Laplacian were used

## A general family of Gagliardo-Nirenberg inequalities

### Theorem (JD, del Pino)

Let  $p \in (1, d]$ ,  $a > 1$  such that  $a \leq \frac{p(d-1)}{d-p}$  if  $p < d$ , and  $b = p \frac{a-1}{p-1}$ .  
There exists a positive constant  $C_{p,a}$  such that, for any function  $f \in L^a(\mathbb{R}^d, dx)$  with  $\nabla f \in L^p(\mathbb{R}^d, dx)$ , we have

$$\|f\|_{L^b(\mathbb{R}^d)} \leq C_{p,a} \|\nabla f\|_{L^p(\mathbb{R}^d)}^\theta \|f\|_{L^a(\mathbb{R}^d)}^{1-\theta} \quad \text{with } \theta = \frac{(a-p)d}{(a-1)(dp - (d-p)a)} \quad (2)$$

if  $a > p$ . A similar inequality also holds if  $a < p$ , namely

$$\|f\|_{L^a(\mathbb{R}^d)} \leq C_{p,a} \|\nabla f\|_{L^p(\mathbb{R}^d)}^\theta \|f\|_{L^b(\mathbb{R}^d)}^{1-\theta} \quad \text{with } \theta = \frac{(p-a)d}{a(d(p-a) + p(a-1))}$$

In both cases, equality holds for any function taking the form

$$f(x) = A \left( 1 + B |x - x_0|^{\frac{p}{p-1}} \right)_+^{-\frac{p-1}{a-p}} \quad \forall x \in \mathbb{R}^d$$

for some  $(A, B, x_0) \in \mathbb{R} \times \mathbb{R} \times \mathbb{R}^d$ , where  $B$  has the sign of  $a - p$

## Remark: the other endpoint of the family

For  $a = p$ , inequality (2) degenerates into an equality. By subtracting it to the inequality, dividing by  $a - p$  and taking the limit as  $a \rightarrow p_+$ , we obtain an *optimal Euclidean  $L^p$ -Sobolev logarithmic inequality* which goes as follows. Assume that  $1 < p \leq d$ . Then for any  $u \in W^{1,p}(\mathbb{R}^d)$  with  $\int_{\mathbb{R}^d} |u|^p dx = 1$  we have

$$\int_{\mathbb{R}^d} |u|^p \log |u|^p dx \leq \frac{d}{p} \log \left[ \beta_{p,d} \int_{\mathbb{R}^d} |\nabla u|^p dx \right]$$

where the optimal constant  $\beta_{p,d}$  is explicit. Equality holds if and only if for some  $\sigma > 0$  and  $x_0 \in \mathbb{R}^d$

$$u(x) = \left[ \frac{1}{2\pi^{\frac{d}{2}}} \frac{p}{p-1} \frac{\Gamma\left(\frac{d}{2}\right)}{\Gamma\left(d\frac{p-1}{p}\right)} \left(\frac{p}{\sigma}\right)^d \frac{p-1}{p} \right]^{\frac{1}{p}} e^{-\frac{1}{\sigma}|x-x_0|^{\frac{p}{p-1}}} \quad \forall x \in \mathbb{R}^d$$

[del Pino, JD], [del Pino, JD, Gentil]  
[Cordero-Erausquin, Gangbo, Houdré]

# Legendre duality through flows

From Onofri to logarithmic Hardy-Littlewood-Sobolev (log HLS) inequalities: the critical case



## Legendre duality

To a convex functional  $F$ , we may associate the functional  $F^*$  defined by Legendre's duality as

$$F^*[v] := \sup \left( \int_{\mathbb{R}^d} u v \, dx - F[u] \right)$$

- To  $F_1[u] = \frac{1}{2} \|u\|_{L^p(\mathbb{R}^d)}^2$ , we associate  $F_1^*[v] = \frac{1}{2} \|v\|_{L^q(\mathbb{R}^d)}^2$  where  $p$  and  $q$  are Hölder conjugate exponents:  $1/p + 1/q = 1$
- To  $F_2[u] = \frac{1}{2} S_d \|\nabla u\|_{L^2(\mathbb{R}^d)}^2$ , we associate

$$F_2^*[v] = \frac{1}{2} S_d^{-1} \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx$$

where  $(-\Delta)^{-1} v = G_d * v$ ,  $G_d(x) = \frac{1}{d-2} |\mathbb{S}^{d-1}|^{-1} |x|^{2-d}$  if  $d \geq 3$

As a straightforward consequence of Legendre's duality, if we have a functional inequality of the form  $F_1[u] \leq F_2[u]$ , then we have the dual inequality  $F_1^*[v] \geq F_2^*[v]$

## Sobolev and HLS

As it has been noticed by E. Lieb, Sobolev's inequality in  $\mathbb{R}^d$ ,  $d \geq 3$ ,

$$\|u\|_{L^{2^*}(\mathbb{R}^d)}^2 \leq S_d \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 \quad \forall u \in \mathcal{D}^{1,2}(\mathbb{R}^d)$$

and the Hardy-Littlewood-Sobolev inequality

$$S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \geq \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx \quad \forall v \in L^{\frac{2d}{d+2}}(\mathbb{R}^d)$$

are dual of each other. Here  $S_d$  is the Aubin-Talenti constant and  $2^* = \frac{2d}{d-2}$

## Using a nonlinear flow to relate Sobolev and HLS

Consider the *fast diffusion* equation

$$\frac{\partial v}{\partial t} = \Delta v^m \quad t > 0, \quad x \in \mathbb{R}^d \quad (3)$$

If we define  $H(t) := H_d[v(t, \cdot)]$ , with

$$H_d[v] := \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2$$

then we observe that

$$\frac{1}{2} H' = - \int_{\mathbb{R}^d} v^{m+1} \, dx + S_d \left( \int_{\mathbb{R}^d} v^{\frac{2d}{d+2}} \, dx \right)^{\frac{2}{d}} \int_{\mathbb{R}^d} \nabla v^m \cdot \nabla v^{\frac{d-2}{d+2}} \, dx$$

where  $v = v(t, \cdot)$  is a solution of (3). With the choice  $m = \frac{d-2}{d+2}$ , we find that  $m + 1 = \frac{2d}{d+2}$

## Relating Sobolev and HLS inequalities through a flow

### Theorem (JD)

Assume that  $d \geq 3$  and  $m = \frac{d-2}{d+2}$ . If  $v$  is a solution of (3) with nonnegative initial datum in  $L^{2d/(d+2)}(\mathbb{R}^d)$ , then

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left[ \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \right] \\ = \left( \int_{\mathbb{R}^d} v^{m+1} \, dx \right)^{2/d} \left[ S_d \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 - \|u\|_{L^{2^*}(\mathbb{R}^d)}^2 \right] \geq 0 \end{aligned}$$

Assume that  $d \geq 5$  and let  $q = \frac{d+2}{d-2}$ . There exists a positive constant  $C \leq (1 + \frac{2}{d}) (1 - e^{-d/2}) S_d$  such that

$$\begin{aligned} S_d \|w^q\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 - \int_{\mathbb{R}^d} w^q (-\Delta)^{-1} w^q \, dx \\ \leq C \|w\|_{L^{2^*}(\mathbb{R}^d)}^{\frac{8}{d-2}} \left[ \|\nabla w\|_{L^2(\mathbb{R}^d)}^2 - S_d \|w\|_{L^{2^*}(\mathbb{R}^d)}^2 \right] \end{aligned}$$

## The two-dimensional case: Legendre duality

Onofri's inequality

$$F_1[u] := \log \left( \int_{\mathbb{R}^2} e^u d\mu \right) \leq \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} u \mu dx =: F_2[u]$$

By duality:  $F_i^*[v] = \sup \left( \int_{\mathbb{R}^2} v u d\mu - F_i[u] \right)$  we can relate Onofri's inequality with (logHLS)

Proposition (E. Carlen, M. Loss & V. Calvez, L. Corrias)

For any  $v \in L^1_+(\mathbb{R}^2)$  with  $\int_{\mathbb{R}^2} v dx = 1$ , such that  $v \log v$  and  $(1 + \log|x|^2)v \in L^1(\mathbb{R}^2)$ , we have

$$F_1^*[v] - F_2^*[v] = \int_{\mathbb{R}^2} v \log \left( \frac{v}{\mu} \right) dx - 4\pi \int_{\mathbb{R}^2} (v - \mu) (-\Delta)^{-1} (v - \mu) dx \geq 0$$

## The two-dimensional case: (logHLS) and flows

[E. Carlen, J. Carrillo, M. Loss]

$$H_2[v] := \int_{\mathbb{R}^2} (v - \mu) (-\Delta)^{-1} (v - \mu) dx - \frac{1}{4\pi} \int_{\mathbb{R}^2} v \log \left( \frac{v}{\mu} \right) dx$$

is related to Gagliardo-Nirenberg inequalities if  $v_t = \Delta \sqrt{v}$

Alternatively, assume that  $v$  is a positive solution of

$$\frac{\partial v}{\partial t} = \Delta \log \left( \frac{v}{\mu} \right) \quad t > 0, \quad x \in \mathbb{R}^2$$

### Proposition (JD)

If  $v$  is a solution with nonnegative initial datum  $v_0$  in  $L^1(\mathbb{R}^2)$  such that  $\int_{\mathbb{R}^2} v_0 dx = 1$ ,  $v_0 \log v_0 \in L^1(\mathbb{R}^2)$  and  $v_0 \log \mu \in L^1(\mathbb{R}^2)$ , then

$$\frac{d}{dt} H_2[v(t, \cdot)] = \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla u|^2 dx - \int_{\mathbb{R}^2} (e^{\frac{u}{2}} - 1) u d\mu \geq F_2[u] - F_1[u]$$

with  $\log(v/\mu) = u/2$

## The two-dimensional case: the sphere setting

The image  $w$  of  $v$  by the inverse stereographic projection on the sphere  $\mathbb{S}^2$ , up to a scaling, solves the equation

$$\frac{\partial w}{\partial t} = \Delta_{\mathbb{S}^2} \log w \quad t > 0, \quad y \in \mathbb{S}^2$$

More precisely, if  $x = (x_1, x_2) \in \mathbb{R}^2$ , then  $u$  and  $w$  are related by

$$w(t, y) = \frac{u(t, x)}{4\pi\mu(x)}, \quad y = \left( \frac{2(x_1, x_2)}{1+|x|^2}, \frac{1-|x|^2}{1+|x|^2} \right) \in \mathbb{S}^2$$

The loss of mass of the solution of

$$\frac{\partial v}{\partial t} = \Delta \log v \quad t > 0, \quad x \in \mathbb{R}^2$$

is compensated in case of

$$\frac{\partial v}{\partial t} = \Delta \log \left( \frac{v}{\mu} \right) \quad t > 0, \quad x \in \mathbb{R}^2$$

by the source term  $-\Delta \log \mu$

# Subcritical log HLS inequalities and the subcritical Keller-Segel model

Optimal inequalities in the subcritical case provide a functional framework for large time asymptotics of the Keller-Segel model

When mass increases to its critical value ( $M = 8\pi$ ), the log HLS inequality is recovered as a singular limit



## The parabolic-elliptic Keller and Segel system

$$\begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \nabla v) & x \in \mathbb{R}^2, t > 0 \\ -\Delta v = u & x \in \mathbb{R}^2, t > 0 \\ u(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

We make the choice:

$$v(t, x) = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \log |x - y| u(t, y) dy$$

and observe that

$$\nabla v(t, x) = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \frac{x - y}{|x - y|^2} u(t, y) dy$$

Mass conservation:  $\frac{d}{dt} \int_{\mathbb{R}^2} u(t, x) dx = 0$

Finite time blow-up if  $M = \int_{\mathbb{R}^2} n_0 dx > 8\pi$

## Existence and free energy

$M = \int_{\mathbb{R}^2} n_0 dx \leq 8\pi$ : global existence [W. Jäger, S. Luckhaus],  
[JD, B. Perthame], [A. Blanchet, JD, B. Perthame], [A. Blanchet,  
J.A. Carrillo, N. Masmoudi]

If  $u$  solves

$$\frac{\partial u}{\partial t} = \nabla \cdot [u (\nabla (\log u) - \nabla v)]$$

the free energy

$$F[u] := \int_{\mathbb{R}^2} u \log u dx - \frac{1}{2} \int_{\mathbb{R}^2} u v dx$$

satisfies

$$\frac{d}{dt} F[u(t, \cdot)] = - \int_{\mathbb{R}^2} u |\nabla (\log u) - \nabla v|^2 dx$$

(log HLS) inequality [E. Carlen, M. Loss]:

$F$  is bounded from below if  $M < 8\pi$

## The existence setting

$$\begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \nabla v) & x \in \mathbb{R}^2, t > 0 \\ -\Delta v = u & x \in \mathbb{R}^2, t > 0 \\ u(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

Initial conditions

$$n_0 \in L^1_+(\mathbb{R}^2, (1+|x|^2) dx), \quad n_0 \log n_0 \in L^1(\mathbb{R}^2, dx), \quad M := \int_{\mathbb{R}^2} n_0(x) dx < 8\pi$$

Global existence and mass conservation:  $M = \int_{\mathbb{R}^2} u(x, t) dx$  for any  $t \geq 0$ , see [W. Jäger, S. Luckhaus], [A. Blanchet, JD, B. Perthame]

$$v = -\frac{1}{2\pi} \log |\cdot| * u$$

## Time-dependent rescaling

$$u(x, t) = \frac{1}{R^2(t)} n \left( \frac{x}{R(t)}, \tau(t) \right) \quad \text{and} \quad v(x, t) = c \left( \frac{x}{R(t)}, \tau(t) \right)$$

with  $R(t) = \sqrt{1 + 2t}$  and  $\tau(t) = \log R(t)$

$$\begin{cases} \frac{\partial n}{\partial t} = \Delta n - \nabla \cdot (n(\nabla c - x)) & x \in \mathbb{R}^2, t > 0 \\ c = -\frac{1}{2\pi} \log |\cdot| * n & x \in \mathbb{R}^2, t > 0 \\ n(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

[A. Blanchet, JD, B. Perthame] Convergence in self-similar variables

$$\lim_{t \rightarrow \infty} \|n(\cdot, \cdot + t) - n_\infty\|_{L^1(\mathbb{R}^2)} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \|\nabla c(\cdot, \cdot + t) - \nabla c_\infty\|_{L^2(\mathbb{R}^2)} = 0$$

means *intermediate asymptotics* in original variables:

$$\|u(x, t) - \frac{1}{R^2(t)} n_\infty \left( \frac{x}{R(t)}, \tau(t) \right)\|_{L^1(\mathbb{R}^2)} \searrow 0$$

## The stationary solution in self-similar variables

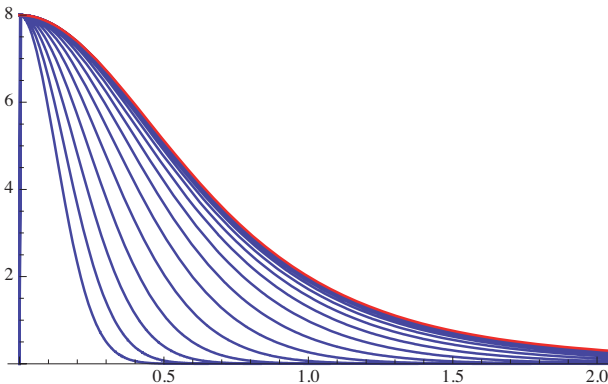
$$n_\infty = M \frac{e^{c_\infty - |x|^2/2}}{\int_{\mathbb{R}^2} e^{c_\infty - |x|^2/2} dx} = -\Delta c_\infty, \quad c_\infty = -\frac{1}{2\pi} \log |\cdot| * n_\infty$$

- Radial symmetry [Y. Naito]
- Uniqueness [P. Biler, G. Karch, P. Laurençot, T. Nadzieja]
- As  $|x| \rightarrow +\infty$ ,  $n_\infty$  is dominated by  $e^{-(1-\epsilon)|x|^2/2}$  for any  $\epsilon \in (0, 1)$  [A. Blanchet, JD, B. Perthame]
- Bifurcation diagram of  $\|n_\infty\|_{L^\infty(\mathbb{R}^2)}$  as a function of  $M$ :

$$\lim_{M \rightarrow 0_+} \|n_\infty\|_{L^\infty(\mathbb{R}^2)} = 0$$

[D.D. Joseph, T.S. Lundgren] [JD, R. Stańczy]

## The stationary solution when mass varies



**Figure:** Representation of the solution appropriately scaled so that the  $8\pi$  case appears as a limit (in red)

## The free energy in self-similar variables

$$\frac{\partial n}{\partial t} = \nabla \left[ n (\log n - x + \nabla c) \right]$$

$$F[n] := \int_{\mathbb{R}^2} n \log n \, dx + \int_{\mathbb{R}^2} \frac{1}{2} |x|^2 n \, dx - \frac{1}{2} \int_{\mathbb{R}^2} n c \, dx$$

satisfies

$$\frac{d}{dt} F[n(t, \cdot)] = - \int_{\mathbb{R}^2} n |\nabla (\log n) + x - \nabla c|^2 \, dx$$

A last remark on  $8\pi$  and scalings:  $n^\lambda(x) = \lambda^2 n(\lambda x)$

$$F[n^\lambda] = F[n] + \int_{\mathbb{R}^2} n \log(\lambda^2) \, dx + \int_{\mathbb{R}^2} \frac{\lambda^{-2}-1}{2} |x|^2 n \, dx + \frac{1}{4\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} n(x) n(y) \log \frac{1}{\lambda} \, dx \, dy$$

$$F[n^\lambda] - F[n] = \underbrace{\left( 2M - \frac{M^2}{4\pi} \right)}_{>0 \text{ if } M < 8\pi} \log \lambda + \frac{\lambda^{-2} - 1}{2} \int_{\mathbb{R}^2} |x|^2 n \, dx$$

## Keller-Segel with subcritical mass in self-similar variables

$$\begin{cases} \frac{\partial n}{\partial t} = \Delta n - \nabla \cdot (n(\nabla c - x)) & x \in \mathbb{R}^2, t > 0 \\ c = -\frac{1}{2\pi} \log |\cdot| * n & x \in \mathbb{R}^2, t > 0 \\ n(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

$$\lim_{t \rightarrow \infty} \|n(\cdot, \cdot + t) - n_\infty\|_{L^1(\mathbb{R}^2)} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \|\nabla c(\cdot, \cdot + t) - \nabla c_\infty\|_{L^2(\mathbb{R}^2)} = 0$$

$$n_\infty = M \frac{e^{c_\infty - |x|^2/2}}{\int_{\mathbb{R}^2} e^{c_\infty - |x|^2/2} dx} = -\Delta c_\infty, \quad c_\infty = -\frac{1}{2\pi} \log |\cdot| * n_\infty$$



# A parametrization of the solutions and the linearized operator

[J. Campos, JD]

$$-\Delta c = M \frac{e^{-\frac{1}{2}|\mathbf{x}|^2 + c}}{\int_{\mathbb{R}^2} e^{-\frac{1}{2}|\mathbf{x}|^2 + c} d\mathbf{x}}$$

Solve

$$-\phi'' - \frac{1}{r} \phi' = e^{-\frac{1}{2}r^2 + \phi}, \quad r > 0$$

with initial conditions  $\phi(0) = a$ ,  $\phi'(0) = 0$  and get with  $r = |\mathbf{x}|$

$$M(a) := 2\pi \int_{\mathbb{R}^2} e^{-\frac{1}{2}r^2 + \phi_a} d\mathbf{x}$$

$$n_a(\mathbf{x}) = M(a) \frac{e^{-\frac{1}{2}r^2 + \phi_a(r)}}{2\pi \int_{\mathbb{R}^2} r e^{-\frac{1}{2}r^2 + \phi_a} d\mathbf{x}} = e^{-\frac{1}{2}r^2 + \phi_a(r)}$$

# Mass

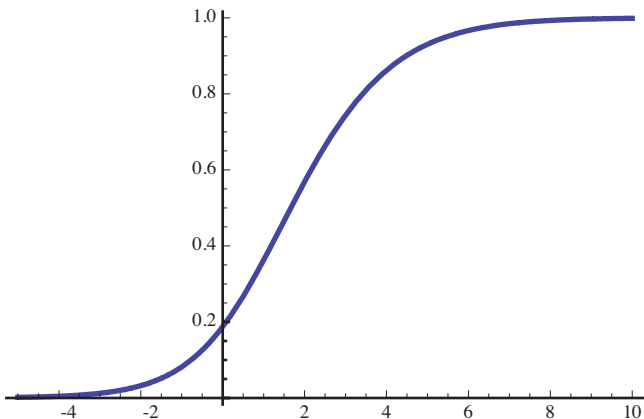
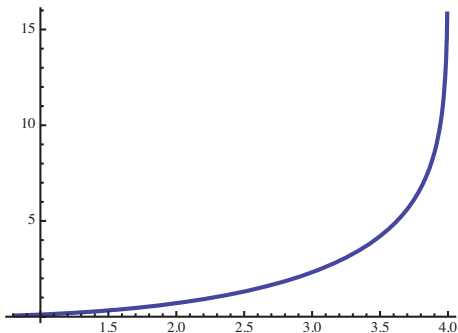


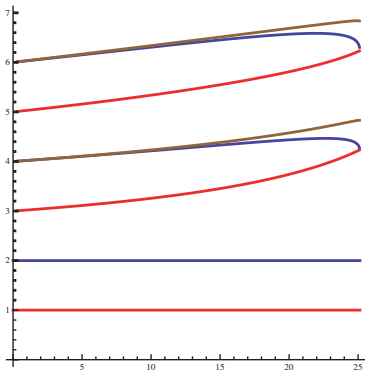
Figure: The mass can be computed as  $M(a) = 2\pi \int_0^\infty n_a(r) r dr$ . Plot of  $a \mapsto M(a)/8\pi$

## Bifurcation diagram



**Figure:** *The bifurcation diagram can be parametrized by  $a \mapsto (\frac{1}{2\pi} M(a), \|c_a\|_{L^\infty(\mathbb{R}^d)})$  with  $\|c_a\|_{L^\infty(\mathbb{R}^d)} = c_a(0) = a - b(a)$  (cf. Keller-Segel system in a ball with no flux boundary conditions)*

## Spectrum of $\mathcal{L}$ (lowest eigenvalues only)



**Figure:** The lowest eigenvalues of  $-\mathcal{L} = (-\Delta)^{-1}(n_a f)$  (shown as a function of the mass) are 0, 1 and 2, thus establishing that the spectral gap of  $-\mathcal{L}$  is 1

[A. Blanchet, JD, M. Escobedo, J. Fernández], [J. Campos, JD],  
[V. Calvez, J.A. Carrillo], [J. Bedrossian, N. Masmoudi]

# Spectral analysis in the functional framework determined by the relative entropy method

## Simple eigenfunctions

**Kernel** Let  $f_0 = \frac{\partial}{\partial M} c_\infty$  be the solution of

$$-\Delta f_0 = n_\infty f_0$$

and observe that  $g_0 = f_0/c_\infty$  is such that

$$\frac{1}{n_\infty} \nabla \cdot (n_\infty \nabla (f_0 - c_\infty g_0)) =: \mathcal{L} f_0 = 0$$

**Lowest non-zero eigenvalues**  $f_1 := \frac{1}{n_\infty} \frac{\partial n_\infty}{\partial x_1}$  associated with  $g_1 = \frac{1}{c_\infty} \frac{\partial c_\infty}{\partial x_1}$  is an eigenfunction of  $\mathcal{L}$ , such that  $-\mathcal{L} f_1 = f_1$

With  $D := x \cdot \nabla$ , let  $f_2 = 1 + \frac{1}{2} D \log n_\infty = 1 + \frac{1}{2n_\infty} D n_\infty$ . Then

$$-\Delta (D c_\infty) + 2 \Delta c_\infty = D n_\infty = 2 (f_2 - 1) n_\infty$$

and so  $g_2 := \frac{1}{c_\infty} (-\Delta)^{-1} (n_\infty f_2)$  is such that  $-\mathcal{L} f_2 = 2 f_2$

## Functional setting...

Lemma (A. Blanchet, JD, B. Perthame)

*Sub-critical HLS inequality*

$$F[n] := \int_{\mathbb{R}^2} n \log \left( \frac{n}{n_\infty} \right) dx - \frac{1}{2} \int_{\mathbb{R}^2} (n - n_\infty)(c - c_\infty) dx \geq 0$$

achieves its minimum for  $n = n_\infty$

$$Q_1[f] = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^2} F[n_\infty(1 + \varepsilon f)] \geq 0$$

if  $\int_{\mathbb{R}^2} f n_\infty dx = 0$ . Notice that  $f_0$  generates the kernel of  $Q_1$

Lemma (J. Campos, JD)

*HLS-Poincaré type inequality* For any  $f \in H^1(\mathbb{R}^2, n_\infty dx)$  such that  $\int_{\mathbb{R}^2} f n_\infty dx = 0$ , we have

$$\int_{\mathbb{R}^2} |\nabla(-\Delta)^{-1}(f n_\infty)|^2 n_\infty dx = \int_{\mathbb{R}^2} |\nabla(g c_\infty)|^2 n_\infty dx \leq \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

## ... and eigenvalues

With  $g$  such that  $-\Delta(g c_\infty) = f n_\infty$ ,  $Q_1$  determines a scalar product

$$\langle f_1, f_2 \rangle := \int_{\mathbb{R}^2} f_1 f_2 n_\infty dx - \int_{\mathbb{R}^2} f_1 n_\infty (g_2 c_\infty) dx$$

on the orthogonal space to  $f_0$  in  $L^2(n_\infty dx)$

$$Q_2[f] := \int_{\mathbb{R}^2} |\nabla(f - g c_\infty)|^2 n_\infty dx \quad \text{with} \quad g = -\frac{1}{c_\infty} \frac{1}{2\pi} \log|\cdot| * (f n_\infty)$$

is a positive quadratic form, whose polar operator is the self-adjoint operator  $\mathcal{L}$

$$\langle f, \mathcal{L} f \rangle = Q_2[f] \quad \forall f \in \mathcal{D}(L_2)$$

Lemma (J. Campos, JD)

$\mathcal{L}$  has pure discrete spectrum and its lowest eigenvalue is 1



## Linearized Keller-Segel theory



$$\mathcal{L}f = \frac{1}{n_\infty} \nabla \cdot (n_\infty \nabla (f - c_\infty g))$$

Corollary (J. Campos, JD)

$$\langle f, f \rangle \leq \langle \mathcal{L}f, f \rangle$$

The linearized problem takes the form

$$\frac{\partial f}{\partial t} = \mathcal{L}f$$

where  $\mathcal{L}$  is a self-adjoint operator on the orthogonal of  $f_0$  equipped with  $\langle \cdot, \cdot \rangle$ . A solution of

$$\frac{d}{dt} \langle f, f \rangle = -2 \langle \mathcal{L}f, f \rangle$$

has therefore exponential decay

## More functional inequalities

## A new Onofri type inequality

By Legendreduality, sub-critical HLS and HLS-Poincaré type inequalities become

Theorem (J. Campos, JD)

For any  $M \in (0, 8\pi)$ , if  $n_\infty = M \frac{e^{c_\infty - \frac{1}{2}|x|^2}}{\int_{\mathbb{R}^2} e^{c_\infty - \frac{1}{2}|x|^2} dx}$  with  $c_\infty = (-\Delta)^{-1} n_\infty$ ,  $d\mu_M = \frac{1}{M} n_\infty dx$ , we have the inequality

$$\log \left( \int_{\mathbb{R}^2} e^\phi d\mu_M \right) - \int_{\mathbb{R}^2} \phi d\mu_M \leq \frac{1}{2M} \int_{\mathbb{R}^2} |\nabla \phi|^2 dx \quad \forall \phi \in \mathcal{D}_0^{1,2}(\mathbb{R}^2)$$

Corollary (J. Campos, JD)

The following *Poincaré* inequality holds

$$\int_{\mathbb{R}^2} |\psi - \bar{\psi}|^2 n_\infty dx \leq \int_{\mathbb{R}^2} |\nabla \psi|^2 dx \quad \text{where} \quad \bar{\psi} = \int_{\mathbb{R}^2} \psi d\mu_M$$

## An improved interpolation inequality (coercivity estimate)

Lemma (J. Campos, JD)

For any  $f \in L^2(\mathbb{R}^2, n_\infty dx)$  such that  $\int_{\mathbb{R}^2} f f_0 n_\infty dx = 0$  holds, we have

$$-\frac{1}{2\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) n_\infty(x) \log|x-y| f(y) n_\infty(y) dx dy \leq (1-\varepsilon) \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

for some  $\varepsilon > 0$ , where  $g c_\infty = G_2 * (f n_\infty)$  and, if  $\int_{\mathbb{R}^2} f n_\infty dx = 0$  holds,

$$\int_{\mathbb{R}^2} |\nabla(g c_\infty)|^2 dx \leq (1-\varepsilon) \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

An application:

Rates of convergence for the subcritical Keller-Segel model  
(self-similar variables)

## Exponential convergence for any mass $M \leq 8\pi$



If  $n_{0,*}(\sigma)$  stands for the symmetrized function associated to  $n_0$ , assume that for any  $s \geq 0$

$$(H) \quad \exists \varepsilon \in (0, 8\pi - M) \quad \text{such that} \quad \int_0^s n_{0,*}(\sigma) d\sigma \leq \int_{B(0, \sqrt{s/\pi})} n_{\infty, M+\varepsilon}(x) dx$$

### Theorem (J. Campos, JD)

*Under the above assumption, if  $n_0 \in L^2_+(n_\infty^{-1} dx)$  and  $M := \int_{\mathbb{R}^2} n_0 dx < 8\pi$ , then any solution with initial datum  $n_0$  is such that*

$$\int_{\mathbb{R}^2} |n(t, x) - n_\infty(x)|^2 \frac{dx}{n_\infty(x)} \leq C e^{-2t} \quad \forall t \geq 0$$

*for some positive constant  $C$ , where  $n_\infty$  is the unique stationary solution with mass  $M$*

Thank you for your attention !