

The Keller-Segel model and the logarithmic Hardy-Littlewood-Sobolev inequality

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Kaust

A – The two-dimensional parabolic-elliptic Keller-Segel model

- Introduction: the two-dimensional parabolic-elliptic Keller-Segel model, the $M < 8\pi$ regime, scalings, etc
- The asymptotic behaviour of the solutions of the Keller-Segel model for small mass

References on the general theory of Keller-Segel systems:

[[Horstmann](#)]

Disclaimer: many references !

Introduction

The parabolic-elliptic Keller and Segel system

$$\begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \nabla v) & x \in \mathbb{R}^2, t > 0 \\ -\Delta v = u & x \in \mathbb{R}^2, t > 0 \\ u(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

We make the choice:

$$v(t, x) = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \log |x - y| u(t, y) dy$$

and observe that

$$\nabla v(t, x) = -\frac{1}{2\pi} \int_{\mathbb{R}^2} \frac{x - y}{|x - y|^2} u(t, y) dy$$

Mass conservation: $\frac{d}{dt} \int_{\mathbb{R}^2} u(t, x) dx = 0$

Blow-up

$M = \int_{\mathbb{R}^2} n_0 \, dx > 8\pi$ and $\int_{\mathbb{R}^2} |x|^2 n_0 \, dx < \infty$: blow-up in finite time
a solution u of

$$\frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \nabla v)$$

satisfies

$$\begin{aligned} & \frac{d}{dt} \int_{\mathbb{R}^2} |x|^2 u(t, x) \, dx \\ &= - \underbrace{\int_{\mathbb{R}^2} 2x \cdot \nabla u \, dx}_{-4M} + \frac{1}{2\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} \underbrace{\frac{2x \cdot (y-x)}{|x-y|^2} u(t, x) u(t, y)}_{\frac{(x-y) \cdot (y-x)}{|x-y|^2} u(t, x) u(t, y)} \, dx \, dy \\ &= 4M - \frac{M^2}{2\pi} < 0 \quad \text{if } M > 8\pi \end{aligned}$$

Existence and free energy

$M = \int_{\mathbb{R}^2} n_0 dx \leq 8\pi$: global existence [Jäger, Luckhaus], [JD, Perthame],
[Blanchet, JD, Perthame], [Blanchet, Carrillo, Masmoudi]

If u solves

$$\frac{\partial u}{\partial t} = \nabla \cdot [u (\nabla (\log u) - \nabla v)]$$

the free energy

$$F[u] := \int_{\mathbb{R}^2} u \log u dx - \frac{1}{2} \int_{\mathbb{R}^2} u v dx$$

satisfies

$$\frac{d}{dt} F[u(t, \cdot)] = - \int_{\mathbb{R}^2} u |\nabla (\log u) - \nabla v|^2 dx$$

Log HLS inequality [Carlen, Loss]: F is bounded from below if
 $M < 8\pi$

The dimension $d = 2$

- In dimension d , the norm $L^{d/2}(\mathbb{R}^d)$ is critical. If $d = 2$, the mass is critical
- Scale invariance: if (u, v) is a solution in \mathbb{R}^2 of the parabolic-elliptic Keller and Segel system, then

$$\left(\lambda^2 u(\lambda^2 t, \lambda x), v(\lambda^2 t, \lambda x) \right)$$

is also a solution

- For $M < 8\pi$, the solution vanishes as $t \rightarrow \infty$, but saying that "diffusion dominates" is not correct: to see this, study "intermediate asymptotics"

The existence setting

$$\begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \nabla v) & x \in \mathbb{R}^2, t > 0 \\ -\Delta v = u & x \in \mathbb{R}^2, t > 0 \\ u(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

Initial conditions

$$n_0 \in L^1_+(\mathbb{R}^2, (1+|x|^2) dx), \quad n_0 \log n_0 \in L^1(\mathbb{R}^2, dx), \quad M := \int_{\mathbb{R}^2} n_0(x) dx < 8\pi$$

Global existence and mass conservation: $M = \int_{\mathbb{R}^2} u(x, t) dx$ for any $t \geq 0$, see [Jäger-Luckhaus], [Blanchet, JD, Perthame]

$$v = -\frac{1}{2\pi} \log |\cdot| * u$$

Time-dependent rescaling

$$u(x, t) = \frac{1}{R^2(t)} n \left(\frac{x}{R(t)}, \tau(t) \right) \quad \text{and} \quad v(x, t) = c \left(\frac{x}{R(t)}, \tau(t) \right)$$

with $R(t) = \sqrt{1 + 2t}$ and $\tau(t) = \log R(t)$

$$\begin{cases} \frac{\partial n}{\partial t} = \Delta n - \nabla \cdot (n(\nabla c - x)) & x \in \mathbb{R}^2, t > 0 \\ c = -\frac{1}{2\pi} \log |\cdot| * n & x \in \mathbb{R}^2, t > 0 \\ n(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

[Blanchet, JD, Perthame] Convergence in self-similar variables

$$\lim_{t \rightarrow \infty} \|n(\cdot, \cdot + t) - n_\infty\|_{L^1(\mathbb{R}^2)} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \|\nabla c(\cdot, \cdot + t) - \nabla c_\infty\|_{L^2(\mathbb{R}^2)} = 0$$

means "intermediate asymptotics" in original variables:

$$\|u(x, t) - \frac{1}{R^2(t)} n_\infty \left(\frac{x}{R(t)}, \tau(t) \right)\|_{L^1(\mathbb{R}^2)} \searrow 0$$

The stationary solution in self-similar variables

$$n_\infty = M \frac{e^{c_\infty - |x|^2/2}}{\int_{\mathbb{R}^2} e^{c_\infty - |x|^2/2} dx} = -\Delta c_\infty, \quad c_\infty = -\frac{1}{2\pi} \log |\cdot| * n_\infty$$

- Radial symmetry [Naito]
- Uniqueness [Biler, Karch, Laurençot, Nadzieja]
- As $|x| \rightarrow +\infty$, n_∞ is dominated by $e^{-(1-\epsilon)|x|^2/2}$ for any $\epsilon \in (0, 1)$ [Blanchet, JD, Perthame]
- Bifurcation diagram of $\|n_\infty\|_{L^\infty(\mathbb{R}^2)}$ as a function of M :

$$\lim_{M \rightarrow 0_+} \|n_\infty\|_{L^\infty(\mathbb{R}^2)} = 0$$

[Joseph, Lundgreen] [JD, Stańczy]

The free energy in self-similar variables

$$\frac{\partial n}{\partial t} = \nabla \left[n (\log n - x + \nabla c) \right]$$

$$F[n] := \int_{\mathbb{R}^2} n \log n \, dx + \int_{\mathbb{R}^2} \frac{1}{2} |x|^2 n \, dx - \frac{1}{2} \int_{\mathbb{R}^2} n c \, dx$$

satisfies

$$\frac{d}{dt} F[n(t, \cdot)] = - \int_{\mathbb{R}^2} n |\nabla (\log n) + x - \nabla c|^2 \, dx$$

A last remark on 8π and scalings: $n^\lambda(x) = \lambda^2 n(\lambda x)$

$$F[n^\lambda] = F[n] + \int_{\mathbb{R}^2} n \log(\lambda^2) \, dx + \int_{\mathbb{R}^2} \frac{\lambda^{-2}-1}{2} |x|^2 n \, dx + \frac{1}{4\pi} \int_{\mathbb{R}^2 \times \mathbb{R}^2} n(x) n(y) \log \frac{1}{\lambda} \, dx \, dy$$

$$F[n^\lambda] - F[n] = \underbrace{\left(2M - \frac{M^2}{4\pi} \right)}_{>0 \text{ if } M < 8\pi} \log \lambda + \frac{\lambda^{-2}-1}{2} \int_{\mathbb{R}^2} |x|^2 n \, dx$$

Theorem

There exists a positive constant M^* such that, for any initial data $n_0 \in L^2(n_\infty^{-1} dx)$ of mass $M < M^*$ satisfying the above assumptions, there is a unique solution $n \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^2)) \cap L^\infty((\tau, \infty) \times \mathbb{R}^2)$ for any $\tau > 0$

Moreover, there are two positive constants, C and δ , such that

$$\int_{\mathbb{R}^2} |n(t, x) - n_\infty(x)|^2 \frac{dx}{n_\infty} \leq C e^{-\delta t} \quad \forall t > 0$$

As a function of M , δ is such that $\lim_{M \rightarrow 0^+} \delta(M) = 1$

The condition $M \leq 8\pi$ is necessary and sufficient for the global existence of the solutions, but there are two extra smallness conditions in our proof:

- Uniform estimate: the *method of the trap*
- *Spectral gap* of a linearised operator \mathcal{L}

Proof of the first result on rates

- First step: the trap
- Second step: weighted H^1 estimates
- Third step: linearization and spectral gap
- Fourth step: collecting the estimates

The parabolic-elliptic Keller and Segel system

$$\begin{cases} \frac{\partial u}{\partial t} = \Delta u - \nabla \cdot (u \nabla v) & x \in \mathbb{R}^2, t > 0 \\ -\Delta v = u & x \in \mathbb{R}^2, t > 0 \\ u(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

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First step: the trap

Decay Estimates of $u(t)$ in $L^\infty(\mathbb{R}^2)$

Lemma

For any $M < M_1$, there exists $C = C(M)$ such that, for any solution $u \in C^0(\mathbb{R}^+, L^1(\mathbb{R}^2)) \cap L^\infty(\mathbb{R}_{\text{loc}}^+ \times \mathbb{R}^2)$

$$\|u(t)\|_{L^\infty(\mathbb{R}^2)} \leq C t^{-1} \quad \forall t > 0$$

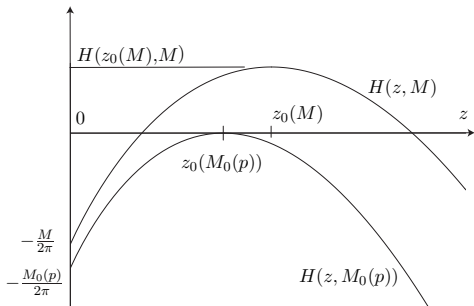
The *method of the trap*... prove that

$$H(\psi(t), M) \leq 0 \quad \text{where} \quad \psi(t) := t \|u(\cdot, t)\|_{L^\infty(\mathbb{R}^2)}$$

where $z \mapsto H(z, M)$ is a continuous function which is

- negative on $[0, z_1)$
- positive on (z_1, z_2) for some z_1, z_2 such that $0 < z_1 < z_2 < \infty$

ψ is continuous and $\psi(0) = 0 \implies \psi(t) \leq z_1 \leq z_0(M)$ for any $t \geq 0$ if $H(z_0(M), M) = \sup_{z \in [z_1, z_2]} H(z, M) \geq 0$



The *method of the trap* amounts to prove that $H(z, M) \leq 0$ implies that $z = \psi(t)$ is bounded by $z_0(M)$ as long as $H(z_0(M), M) > 0$

Duhamel's formula:

$$\begin{aligned} u(x, t_0 + t) &= \int_{\mathbb{R}^2} N(x - y, t) u(y, t_0) dy \\ &= \int_0^t \int_{\mathbb{R}^2} N(x - y, t - s) \nabla \cdot [u(y, t_0 + s) \nabla v(y, t_0 + s)] dy ds \end{aligned}$$

where $N(x, t) = \frac{1}{4\pi t} e^{-|x|^2/(4t)}$. Let $\kappa_\sigma = \|\partial N / \partial x_i(\cdot, 1)\|_{L^\sigma(\mathbb{R}^2)}$

$$\begin{aligned} \|u(\cdot, t_0 + t)\|_{L^\infty(\mathbb{R}^2)} &= \frac{1}{4\pi t} \|u(\cdot, t_0)\|_{L^1(\mathbb{R}^2)} \\ &\leq \sum_{i=1,2} \int_0^t \left\| \frac{\partial N}{\partial x_i}(\cdot, t-s) * \left[\left(u \frac{\partial v}{\partial x_i} \right)(\cdot, t_0 + s) \right] \right\|_{L^\infty(\mathbb{R}^2)} ds \\ &\leq \sum_{i=1,2} \kappa_\sigma \int_0^t (t-s)^{-(1-\frac{1}{\sigma})-\frac{1}{2}} \left\| \left(u \frac{\partial v}{\partial x_i} \right)(\cdot, t_0 + s) \right\|_{L^\rho(\mathbb{R}^2)} ds \end{aligned}$$

HLS inequality + Hölder and take $t_0 = t$

$$\begin{aligned} 2t \|u(\cdot, 2t)\|_{L^\infty(\mathbb{R}^2)} - \frac{M}{2\pi} \\ \leq \frac{2\kappa_\sigma C_{\text{HLS}}}{\pi} M^{\frac{1}{p} + \frac{1}{r}} t \int_0^t (t-s)^{\frac{1}{\sigma} - \frac{3}{2}} (t+s)^{\frac{1}{p} + \frac{1}{r} - 2} [\psi(t)]^{2 - \frac{1}{p} - \frac{1}{r}} ds \end{aligned}$$

with $\psi(t) := \sup_{0 \leq s \leq t} 2s \|u(\cdot, 2s)\|_{L^\infty(\mathbb{R}^2)}$ and

$$t \int_0^t (t-s)^{\frac{1}{\sigma} - \frac{3}{2}} (t+s)^{\frac{1}{p} + \frac{1}{r} - 2} ds = \frac{\sigma}{2 - \sigma}$$

$$\psi(t) \leq \frac{M}{2\pi} + C_0 (\psi(t))^\theta \quad \text{with} \quad C_0 = \frac{2\kappa_\sigma C_{\text{HLS}}}{\pi} M^{\frac{1}{p} + \frac{1}{r}} \frac{\sigma}{2 - \sigma}, \quad \theta = 2 - \frac{1}{p} - \frac{1}{r}$$

Choice: $H(z, M) = z - C_0 z^\theta - M/(2\pi)$

How small is the mass ?

The exponents σ , ρ , p , q and r are related by

$$\left\{ \begin{array}{ll} \frac{1}{\sigma} + \frac{1}{\rho} = 1, & 1 < \sigma < 2 \\ \frac{1}{p} + \frac{1}{q} = \frac{1}{\rho}, & p, q > 2 \\ \frac{1}{r} - \frac{1}{q} = \frac{1}{2}, & r > 1 \end{array} \right.$$

For the choice $r = 4/3$, $q = 4$, $C_{\text{HLS}} = 2\sqrt{\pi}$

$$C_0 = \frac{4\kappa_\sigma}{\sqrt{\pi}} M^{\frac{1}{p} + \frac{1}{4}} \frac{\sigma}{2 - \sigma} \text{ with } \sigma = \frac{4p}{3p-4}$$

... there exists $M_0(p)$ such that $H(z_0(M), M) > 0$ if and only if $M < M_0(p)$ and $\sup_{p \in (4, +\infty)} M_0(p) = \lim_{p \rightarrow +\infty} M_0(p) \approx 0.822663 < 8\pi \approx 25.1327$

Corollary

For any mass $M < M_1$ and all $p \in [1, \infty]$, there exists a positive constant $C = C(p, M)$ with $\lim_{M \rightarrow 0^+} C(p, M) = 0$, such that

$$\|u(t)\|_{L^p(\mathbb{R}^2)} \leq C t^{-(1-\frac{1}{p})} \quad \forall t > 0$$

Remark The above rates are optimal as can easily be checked using the self-similar solutions (n_∞, c_∞)

Second step: weighted H^1 estimates

L^p and H^1 estimates in the self-similar variables

Consider the solution of

$$\begin{cases} \frac{\partial n}{\partial t} = \Delta n - \nabla \cdot (n(\nabla c - x)) & x \in \mathbb{R}^2, t > 0 \\ c = -\frac{1}{2\pi} \log |\cdot| * n & x \in \mathbb{R}^2, t > 0 \\ n(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

For any $p \in (1, \infty]$

$$\|n(t)\|_{L^p(\mathbb{R}^2)} \leq C_1 \quad \forall t > 0$$

for some positive constant C_1 , and for $p > 2$

$$2\pi \|\nabla c(t)\|_{L^\infty} \leq \underbrace{\sup_{x \in \mathbb{R}^2} \int_{|x-y| \geq 1} \frac{n(t, y)}{|x-y|} dy}_{\leq M} + \underbrace{\sup_{x \in \mathbb{R}^2} \int_{|x-y| \leq 1} \frac{n(t, y)}{|x-y|} dy}_{\leq (2\pi)^{\frac{p-1}{p-2}} \|n\|_{L^p(\mathbb{R}^2)}^{\frac{p}{p-1}}}$$

$$\|n(t)\|_{L^p(\mathbb{R}^2)} \leq C_1 \quad \text{and} \quad \|\nabla c(t)\|_{L^\infty(\mathbb{R}^2)} \leq C_2 \quad \forall t > 0$$

Lemma

The constants C_1 and C_2 depend on M and are such that

$$\lim_{M \rightarrow 0_+} C_i(M) = 0 \quad i = 1, 2$$

Exponential weights

With $K = K(x) = e^{|\cdot|^2/2}$, let us rewrite the equation for n as

$$\frac{\partial n}{\partial t} - \frac{1}{K} \nabla \cdot (K \nabla n) = -\nabla c \cdot \nabla n + 2n + n^2$$

Proposition

For any mass $M \in (0, M_1)$, there is a positive constant C such that

$$\|n(t)\|_{H^1(K)} \leq C \quad \forall t > 0$$

First ingredient [M. Escobedo and O. Kavian]: for any $q > 2$ and $\varepsilon > 0$, there exists a positive constant $C(\varepsilon, q)$ such that

$$\int_{\mathbb{R}^2} n^2 K dx \leq \varepsilon \int_{\mathbb{R}^2} |\nabla n|^2 K dx + C(\varepsilon, q) \|n\|_{L^q(\mathbb{R}^2)}^2$$

$L^2(K)$ estimate

Multiply the equation by nK and integrate by parts

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |n|^2 K \, dx + \int_{\mathbb{R}^2} |\nabla n|^2 K \, dx \\ = - \int_{\mathbb{R}^2} n \nabla c \cdot \nabla n K \, dx + 2 \int_{\mathbb{R}^2} n^2 K \, dx + \int_{\mathbb{R}^2} n^3 K \, dx \\ \leq \varepsilon \int_{\mathbb{R}^2} |\nabla n|^2 K \, dx + C \end{aligned}$$

and so

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |n|^2 K \, dx + \underbrace{(1 - \varepsilon) \int_{\mathbb{R}^2} |\nabla n|^2 K \, dx}_{\geq \int_{\mathbb{R}^2} |n|^2 K \, dx} \leq C$$

(expand the square in $\int_{\mathbb{R}^2} |\nabla(nK)|^2 K^{-1} \, dx \geq 0$)

$H^1(K)$ estimate (1/2)

Let $S(t)$ be the semi-group generated by $-K^{-1} \nabla \cdot (K \nabla \cdot)$ on $L^2(K)$

$$n(t, x) = S(t) n_0(x) - \int_0^t S(t-s) (\nabla c \cdot \nabla n)(s) ds + \int_0^t S(t-s) (2n + n^2)(s) ds$$

$$\begin{aligned} & \|n(t)\|_{H^1(K)} - \|S(t) n_0\|_{H^1(K)} \\ & \leq \int_0^t \|S(t-s) (\nabla c \cdot \nabla n)(s)\|_{H^1(K)} ds + \int_0^t \|S(t-s) (2n + n^2)(s)\|_{H^1(K)} ds \end{aligned}$$

Second ingredient: $\|S(t) h\|_{H^1(K)} \leq \kappa (1 + t^{-1/2}) \|h\|_{L^2(K)}$

$$\begin{aligned} & \frac{1}{\kappa} (\|n(t)\|_{H^1(K)} - \|S(t) n_0\|_{H^1(K)}) \\ & \leq C_2 \int_0^t \left(1 + \frac{1}{\sqrt{t-s}}\right) \|\nabla n(s)\|_{L^2(K)} ds + (2 + C_1) \int_0^t \left(1 + \frac{1}{\sqrt{t-s}}\right) \|n(s)\|_{L^2(K)} ds \end{aligned}$$

$H^1(K)$ estimate (2/2)

$$\frac{1}{\kappa} \|n(t+\tau)\|_{H^1(K)} \leq \left(1 + \frac{1}{\sqrt{t}}\right) C_1 + C_3 \int_0^t \left(1 + \frac{1}{\sqrt{t-s}}\right) \|n(s+\tau)\|_{H^1(K)} ds$$

Let $H(T) = \sup_{t \in (0, T)} \int_0^t \left(1 + \frac{1}{\sqrt{t-s}}\right) \|n(s+\tau)\|_{H^1(K)} ds$ and choose

$$T > 0 \text{ such that } \frac{1}{2\kappa} = C_3 \int_0^T \left(1 + \frac{1}{\sqrt{T-s}}\right) ds = C_3(T + 2\sqrt{T})$$

$$\frac{1}{\kappa} H(T) \leq (\pi + 4\sqrt{T} + T) C_1 + \frac{1}{2\kappa} H(T) \implies H(T) \leq 2(\pi + 4\sqrt{T} + T) \kappa C_1$$

For any $t \in (0, T)$

$$\frac{1}{\kappa} \|n(t+\tau)\|_{H^1(K)} \leq \left(1 + \frac{1}{\sqrt{t}}\right) C_1 + C_3 H(T) \leq \left(1 + \frac{1}{\sqrt{t}}\right) C_1 + 2(\pi + 4\sqrt{T} + T) \kappa C_3$$

Conclusion: bound in $H^1(K)$

The estimate

$$\frac{1}{\kappa} \|n(t+\tau)\|_{H^1(K)} \leq \left(1 + \frac{1}{\sqrt{t}}\right) C_1 + C_3 H(T) \leq \left(1 + \frac{1}{\sqrt{t}}\right) C_1 + 2 \left(\pi + 4\sqrt{T} + T\right)$$

for any $t \in (0, T)$ gives a bound on $\|n(T + \tau)\|_{H^1(K)}$ for any $\tau > 0$

Lemma

$$\|n(t)\|_{H^1(K)} \leq C \max \left\{ 1, \frac{\sqrt{T}}{\sqrt{t}} \right\} \quad \forall t > 0$$

Actually $n(t)$ can be bounded also in $H^1(n_\infty^{-1})$ but further estimates are needed...

Third step: linearization and spectral gap

A linearized operator

Introduce f and g defined by

$$n(x, t) = n_\infty(x)(1 + f(x, t)) \quad \text{and} \quad c(x, t) = c_\infty(x)(1 + g(x, t))$$

(f, g) is solution of the non-linear problem

$$\begin{cases} \frac{\partial f}{\partial t} - \mathcal{L}(t, x, f, g) = -\frac{1}{n_\infty} \nabla \cdot [f n_\infty \nabla (g c_\infty)] & x \in \mathbb{R}^2, t > 0 \\ -\Delta(c_\infty g) = f n_\infty & x \in \mathbb{R}^2, t > 0 \end{cases}$$

where \mathcal{L} is the linear operator given by

$$\mathcal{L}(t, x, f, g) = \frac{1}{n_\infty} \nabla \cdot [n_\infty \nabla (f - g c_\infty)]$$

The conservation of mass is replaced here by $\int_{\mathbb{R}^2} f n_\infty dx = 0$

A spectral gap estimate

Proposition

For any $M \in (0, M_2)$, for any $f \in H^1(n_\infty dx)$ such that

$$\int_{\mathbb{R}^2} f n_\infty dx = 0 \implies \int_{\mathbb{R}^2} |\nabla f|^2 n_\infty dx \geq \Lambda(M) \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

for some $\Lambda(M) > 0$ and $\lim_{M \rightarrow 0^+} \Lambda(M) = 1$

Let $h = \sqrt{n_\infty} f = \sqrt{\lambda} e^{-|x|^2/4 + c_\infty/2} f$

$$\lambda |\nabla f|^2 n_\infty = |\nabla h|^2 + \frac{|x|^2}{4} h^2 + \frac{1}{4} |\nabla c_\infty|^2 h^2 + h \nabla h \cdot (x - \nabla c_\infty) - \frac{1}{2} x \cdot \nabla c_\infty h^2$$

(integrations by parts)

$$\begin{aligned} \int_{\mathbb{R}^2} h \nabla h \cdot x dx &= - \int_{\mathbb{R}^2} h^2 dx \\ \int_{\mathbb{R}^2} h \nabla h \cdot \nabla c_\infty dx &= \frac{1}{2} \int_{\mathbb{R}^2} h^2 (-\Delta c_\infty) dx \leq \frac{1}{2} \|n_\infty\|_{L^\infty(\mathbb{R}^2)} \int_{\mathbb{R}^2} h^2 dx \\ \frac{1}{2} \int_{\mathbb{R}^2} x \cdot \nabla c_\infty h^2 dx &\leq \frac{\sigma^2 - 1}{\sigma^2} \int_{\mathbb{R}^2} \frac{|x|^2}{4} h^2 dx + \frac{1}{4} \frac{\sigma^2}{\sigma^2 - 1} \int_{\mathbb{R}^2} |\nabla c_\infty|^2 h^2 dx \end{aligned}$$

$H^1(n_\infty^{-1})$ estimate

Assume that $n_0/n_\infty \in L^2(n_\infty)$

There exists a constant $C > 0$ such that

$$|x| > 1 \implies \left| c_\infty + M/(2\pi) \log |x| \right| \leq C$$

$n_\infty K = e^{c_\infty}$ behaves like $O(|x|^{-M/(2\pi)})$ as $|x| \rightarrow \infty$

$$\frac{\partial n}{\partial t} - n_\infty \nabla \cdot \left(\frac{1}{n_\infty} \nabla n \right) = (\nabla c_\infty - \nabla c) \cdot \nabla n + 2n + n^2$$

Corollary

If $M < M_2$, then any solution n is bounded in

$$L^\infty(\mathbb{R}^+, L^2(n_\infty^{-1} dx)) \cap L^\infty((\tau, \infty), H^1(n_\infty^{-1} dx))$$

for any $\tau > 0$

Fourth step: collecting the estimates

Proof of the exponential rate of convergence

$$\begin{cases} \frac{\partial f}{\partial t} - \mathcal{L}(t, x, f, g) = -\frac{1}{n_\infty} \nabla \cdot [f n_\infty \nabla (g c_\infty)] & x \in \mathbb{R}^2, t > 0 \\ -\Delta(c_\infty g) = f n_\infty & x \in \mathbb{R}^2, t > 0 \end{cases}$$

Multiply by $f n_\infty$ and integrate by parts

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |f|^2 n_\infty dx + \int_{\mathbb{R}^2} |\nabla f|^2 n_\infty dx \\ = \underbrace{\int_{\mathbb{R}^2} \nabla f \cdot \nabla (g c_\infty) n_\infty dx}_{=I} + \underbrace{\int_{\mathbb{R}^2} \nabla f \cdot \nabla (g c_\infty) f n_\infty dx}_{=II} \end{aligned}$$

Cauchy-Schwarz' inequality

$$I = \int_{\mathbb{R}^2} \nabla f \cdot \nabla (g c_\infty) n_\infty dx \leq \|\nabla f\|_{L^2(n_\infty dx)} \|\nabla (g c_\infty)\|_{L^2(n_\infty dx)}$$

Hölder's inequality (with $q > 2$)

$$\|\nabla(g c_\infty)\|_{L^2(n_\infty dx)} \leq M^{1/2-1/q} \|n_\infty\|_{L^\infty(\mathbb{R}^2)}^{1/q} \|\nabla(g c_\infty)\|_{L^q(\mathbb{R}^2)}$$

HLS inequality (with $1/p = 1/2 + 1/q$)

$$\|\nabla(g c_\infty)\|_{L^q(\mathbb{R}^2)} \leq \frac{1}{2\pi} \left(\int_{\mathbb{R}^2} \left| (f n_\infty) * \frac{1}{|\cdot|} \right|^q dx \right)^{\frac{1}{q}} \leq \frac{C_{\text{HLS}}}{2\pi} \|f n_\infty\|_{L^p(\mathbb{R}^2)}$$

Hölder's inequality: $\|f n_\infty\|_{L^p(\mathbb{R}^2)} \leq \|f\|_{L^2(n_\infty dx)} \|n_\infty\|_{L^{q/2}(\mathbb{R}^2)}^{1/2}$

$$I = \int_{\mathbb{R}^2} \nabla f \cdot \nabla(g c_\infty) f n_\infty dx \leq C_*(M) \|f\|_{L^2(n_\infty dx)} \|\nabla f\|_{L^2(n_\infty dx)}$$

$$C_*(M) := C_{\text{HLS}} (2\pi)^{-1} M^{1/2-1/q} \|n_\infty\|_{L^{q/2}(\mathbb{R}^2)}^{1/2} \|n_\infty\|_{L^\infty(\mathbb{R}^2)}^{1/q} \rightarrow 0 \text{ as } M \rightarrow 0$$

second term and conclusion

Use $g c_\infty = c - c_\infty$ and the Cauchy-Schwarz inequality

$$\int_{\mathbb{R}^2} \nabla f \cdot \nabla (g c_\infty) f n_\infty dx \leq \underbrace{\|\nabla c - \nabla c_\infty\|_{L^\infty(\mathbb{R}^2)}}_{\leq 2 C_2(M) \searrow 0} \|f\|_{L^2(n_\infty dx)} \|\nabla f\|_{L^2(n_\infty dx)}$$

Spectral gap estimate

$$\underbrace{\sqrt{\Lambda(M)}}_{\rightarrow 1} \|f\|_{L^2(n_\infty dx)} \leq \|\nabla f\|_{L^2(n_\infty dx)}$$

With $\gamma(M) := \frac{C_*(M) + 2 C_2(M)}{\sqrt{\Lambda(M)}} \searrow 0$,

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |f|^2 n_\infty dx \leq - [1 - \gamma(M)] \int_{\mathbb{R}^2} |\nabla f|^2 n_\infty dx$$

If n_1 and n_2 are two solutions in $C^0(\mathbb{R}^+, L^1(\mathbb{R}^2)) \cap L^\infty((\tau, \infty) \times \mathbb{R}^2)$ for any $\tau > 0$, with $f = (n_2 - n_1)/n_\infty$ we also get

$$\frac{1}{2} \frac{d}{dt} \int_{\mathbb{R}^2} |f|^2 n_\infty dx \leq - [1 - \gamma(M)] \Lambda(M) \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

As a consequence, if the initial condition is the same, then $n_1 = n_2$

B – Sobolev and
Hardy-Littlewood-Sobolev
inequalities:
duality, flows

Gagliardo-Nirenberg inequalities

Consider the following sub-family of Gagliardo-Nirenberg inequalities

$$\|f\|_{L^{2p}(\mathbb{R}^d)} \leq C_{p,d} \|\nabla f\|_{L^2(\mathbb{R}^d)}^\theta \|f\|_{L^{p+1}(\mathbb{R}^d)}^{1-\theta}$$

with $\theta = \theta(p) := \frac{p-1}{p} \frac{d}{d+2-p(d-2)}$

- $1 < p \leq \frac{d}{d-2}$ if $d \geq 3$
- $1 < p < \infty$ if $d = 2$

[M. del Pino, J.D.] equality holds in if $f = F_p$ with

$$F_p(x) = (1 + |x|^2)^{-\frac{1}{p-1}} \quad \forall x \in \mathbb{R}^d$$

and that all extremal functions are equal to F_p up to a multiplication by a constant, a translation and a scaling.

- If $d \geq 3$, the limit case $p = d/(d-2)$ corresponds to Sobolev's inequality [T. Aubin, G. Talenti]
- When $p \rightarrow 1$, we recover the euclidean logarithmic Sobolev inequality in optimal scale invariant form [F. Weissler]
- If $d = 2$ and $p \rightarrow \infty$...

Onofri's inequality as a limit case

When $d = 2$, Onofri's inequality can be seen as an endpoint case of the family of the Gagliardo-Nirenberg inequalities [J.D.]

Proposition

[J.D.] Assume that $g \in \mathcal{D}(\mathbb{R}^d)$ is such that $\int_{\mathbb{R}^2} g \, d\mu = 0$ and let

$$f_p := F_p \left(1 + \frac{g}{2p} \right)$$

With $\mu(x) := \frac{1}{\pi} (1 + |x|^2)^{-2}$, and $d\mu(x) := \mu(x) \, dx$, we have

$$1 \leq \lim_{p \rightarrow \infty} C_{p,2} \frac{\|\nabla f\|_{L^2(\mathbb{R}^2)}^{\theta(p)} \|f\|_{L^{p+1}(\mathbb{R}^2)}^{1-\theta(p)}}{\|f\|_{L^{2p}(\mathbb{R}^2)}} = \frac{e^{\frac{1}{16\pi}} \int_{\mathbb{R}^2} |\nabla g|^2 \, dx}{\int_{\mathbb{R}^2} e^g \, d\mu}$$

The standard form of the euclidean version of Onofri's inequality is

$$\log \left(\int_{\mathbb{R}^2} e^g \, d\mu \right) - \int_{\mathbb{R}^2} g \, d\mu \leq \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla g|^2 \, dx$$

Legendre duality: Onofri and log HLS

Legendre's duality: $F^*[v] := \sup \left(\int_{\mathbb{R}^d} u v \, dx - F[u] \right)$

$$F_1[u] := \log \left(\int_{\mathbb{R}^2} e^u \, d\mu \right) \quad \text{and} \quad F_2[u] := \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla u|^2 \, dx + \int_{\mathbb{R}^2} u \mu \, dx$$

Onofri's inequality amounts to $F_1[u] \leq F_2[u]$ with $d\mu(x) := \mu(x) \, dx$,
 $\mu(x) := \frac{1}{\pi(1+|x|^2)^2}$

Proposition

For any $v \in L^1_+(\mathbb{R}^2)$ with $\int_{\mathbb{R}^2} v \, dx = 1$, such that $v \log v$ and $(1 + \log |x|^2) v \in L^1(\mathbb{R}^2)$, we have

$$F_1^*[v] - F_2^*[v] = \int_{\mathbb{R}^2} v \log \left(\frac{v}{\mu} \right) \, dx - 4\pi \int_{\mathbb{R}^2} (v - \mu) (-\Delta)^{-1} (v - \mu) \, dx \geq 0$$

[Carlen-Loss, Beckner, Calvez-Corrias]

A puzzling result of Carlen, Carrillo and Loss ($d \geq 3$)

[E. Carlen, J.A. Carrillo and M. Loss] The fast diffusion equation

$$\frac{\partial v}{\partial t} = \Delta v^m \quad t > 0, \quad x \in \mathbb{R}^d$$

with exponent $m = d/(d+2)$, when $d \geq 3$, is such that

$$H_d[v] := \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2$$

obeys to

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} H_d[v(t, \cdot)] &= \frac{1}{2} \frac{d}{dt} \left[\int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \right] \\ &= \frac{d(d-2)}{(d-1)^2} S_d \|u\|_{L^{q+1}(\mathbb{R}^d)}^{4/(d-1)} \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 - \|u\|_{L^{2q}(\mathbb{R}^d)}^{2q} \end{aligned}$$

with $u = v^{(d-1)/(d+2)}$ and $q = \frac{d+1}{d-1}$. If $\frac{d(d-2)}{(d-1)^2} S_d = (C_{q,d})^{2q}$, the r.h.s. is nonnegative. Optimality is achieved simultaneously in both functionals (Barenblatt regime): the Hardy-Littlewood-Sobolev inequalities can be improved by an integral remainder term

... and the two-dimensional case

Recall that $(-\Delta)^{-1}v = G_d * v$ with

- $G_d(x) = \frac{1}{d-2} |\mathbb{S}^{d-1}|^{-1} |x|^{2-d}$ if $d \geq 3$
- $G_2(x) = \frac{1}{2\pi} \log|x|$ if $d = 2$

Same computation in dimension $d = 2$ with $m = 1/2$ gives

$$\begin{aligned} \frac{\|v\|_{L^1(\mathbb{R}^2)}}{8} \frac{d}{dt} \left[\frac{4\pi}{\|v\|_{L^1(\mathbb{R}^2)}} \int_{\mathbb{R}^2} v (-\Delta)^{-1} v \, dx - \int_{\mathbb{R}^2} v \log v \, dx \right] \\ = \|u\|_{L^4(\mathbb{R}^2)}^4 \|\nabla u\|_{L^2(\mathbb{R}^2)}^2 - \pi \|v\|_{L^6(\mathbb{R}^2)}^6 \end{aligned}$$

The r.h.s. is one of the Gagliardo-Nirenberg inequalities ($d = 2$, $q = 3$): $\pi (C_{3,2})^6 = 1$

The l.h.s. is bounded from below by the logarithmic Hardy-Littlewood-Sobolev inequality and achieves its minimum if $v = \mu$ with

$$\mu(x) := \frac{1}{\pi(1+|x|^2)^2} \quad \forall x \in \mathbb{R}^2$$

As it has been noticed by E. Lieb, Sobolev's inequality in \mathbb{R}^d , $d \geq 3$,

$$\|u\|_{L^{2^*}(\mathbb{R}^d)}^2 \leq S_d \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 \quad \forall u \in \mathcal{D}^{1,2}(\mathbb{R}^d) \quad (1)$$

and the Hardy-Littlewood-Sobolev inequality

$$S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \geq \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx \quad \forall v \in L^{\frac{2d}{d+2}}(\mathbb{R}^d) \quad (2)$$

are **dual** of each other. Here S_d is the Aubin-Talenti constant and $2^* = \frac{2d}{d-2}$. Can we recover this using a nonlinear flow approach? Can we improve it?

Keller-Segel model: another motivation [Carrillo, Carlen and Loss] and [Blanchet, Carlen and Carrillo]

Using a nonlinear flow to relate Sobolev and HLS

Consider the *fast diffusion* equation

$$\frac{\partial v}{\partial t} = \Delta v^m \quad t > 0, \quad x \in \mathbb{R}^d \quad (3)$$

If we define $H(t) := H_d[v(t, \cdot)]$, with

$$H_d[v] := \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2$$

then we observe that

$$\frac{1}{2} H' = - \int_{\mathbb{R}^d} v^{m+1} \, dx + S_d \left(\int_{\mathbb{R}^d} v^{\frac{2d}{d+2}} \, dx \right)^{\frac{2}{d}} \int_{\mathbb{R}^d} \nabla v^m \cdot \nabla v^{\frac{d-2}{d+2}} \, dx$$

where $v = v(t, \cdot)$ is a solution of (??). With the choice $m = \frac{d-2}{d+2}$, we find that $m + 1 = \frac{2d}{d+2}$

A first statement

Proposition

[J.D.] Assume that $d \geq 3$ and $m = \frac{d-2}{d+2}$. If v is a solution of (??) with nonnegative initial datum in $L^{2d/(d+2)}(\mathbb{R}^d)$, then

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \left[\int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - S_d \|v\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \right] \\ = \left(\int_{\mathbb{R}^d} v^{m+1} \, dx \right)^{\frac{2}{d}} \left[S_d \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 - \|u\|_{L^{2^*}(\mathbb{R}^d)}^2 \right] \geq 0 \end{aligned}$$

The HLS inequality amounts to $H \leq 0$ and appears as a consequence of Sobolev, that is $H' \geq 0$ if we show that $\limsup_{t \rightarrow 0} H(t) = 0$

Notice that $u = v^m$ is an optimal function for (??) if v is optimal for (??)

By integrating along the flow defined by (??), we can actually obtain optimal integral remainder terms which improve on the usual Sobolev inequality (??), but only when $d \geq 5$ for integrability reasons

Theorem

[J.D.] Assume that $d \geq 5$ and let $q = \frac{d+2}{d-2}$. There exists a positive constant $\mathcal{C} \leq (1 + \frac{2}{d}) (1 - e^{-d/2}) S_d$ such that

$$\begin{aligned} S_d \|w^q\|_{L^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 - \int_{\mathbb{R}^d} w^q (-\Delta)^{-1} w^q dx \\ \leq \mathcal{C} \|w\|_{L^{2^*}(\mathbb{R}^d)}^{\frac{8}{d-2}} \left[\|\nabla w\|_{L^2(\mathbb{R}^d)}^2 - S_d \|w\|_{L^{2^*}(\mathbb{R}^d)}^2 \right] \end{aligned}$$

for any $w \in \mathcal{D}^{1,2}(\mathbb{R}^d)$

Solutions with *separation of variables*

Consider the solution of $\frac{\partial v}{\partial t} = \Delta v^m$ vanishing at $t = T$:

$$\bar{v}_T(t, x) = c (T - t)^\alpha (F(x))^{\frac{d+2}{d-2}}$$

where F is the Aubin-Talenti solution of

$$-\Delta F = d(d-2) F^{(d+2)/(d-2)}$$

Let $\|v\|_* := \sup_{x \in \mathbb{R}^d} (1 + |x|^2)^{d+2} |v(x)|$

Lemma

[M. delPino, M. Saez], [J. L. Vázquez, J. R. Esteban, A. Rodríguez]
For any solution v with initial datum $v_0 \in L^{2d/(d+2)}(\mathbb{R}^d)$, $v_0 > 0$, there exists $T > 0$, $\lambda > 0$ and $x_0 \in \mathbb{R}^d$ such that

$$\lim_{t \rightarrow T^-} (T - t)^{-\frac{1}{1-m}} \|v(t, \cdot) / \bar{v}(t, \cdot) - 1\|_* = 0$$

with $\bar{v}(t, x) = \lambda^{(d+2)/2} \bar{v}_T(t, (x - x_0)/\lambda)$

Improved inequality: proof (1/2)

$J(t) := \int_{\mathbb{R}^d} v(t, x)^{m+1} dx$ satisfies

$$J' = -(m+1) \|\nabla v^m\|_{L^2(\mathbb{R}^d)}^2 \leq -\frac{m+1}{S_d} J^{1-\frac{2}{d}}$$

If $d \geq 5$, then we also have

$$J'' = 2m(m+1) \int_{\mathbb{R}^d} v^{m-1} (\Delta v^m)^2 dx \geq 0$$

Notice that

$$\frac{J'}{J} \leq -\frac{m+1}{S_d} J^{-\frac{2}{d}} \leq -\kappa \quad \text{with} \quad \kappa T = \frac{2d}{d+2} \frac{T}{S_d} \left(\int_{\mathbb{R}^d} v_0^{m+1} dx \right)^{-\frac{2}{d}} \leq \frac{d}{2}$$

Improved inequality: proof (2/2)

By the **Cauchy-Schwarz inequality**, we have

$$\begin{aligned} \frac{J'^2}{(m+1)^2} &= \|\nabla v^m\|_{L^2(\mathbb{R}^d)}^4 = \left(\int_{\mathbb{R}^d} v^{(m-1)/2} \Delta v^m \cdot v^{(m+1)/2} dx \right)^2 \\ &\leq \int_{\mathbb{R}^d} v^{m-1} (\Delta v^m)^2 dx \int_{\mathbb{R}^d} v^{m+1} dx = \text{Cst } J'' J \end{aligned}$$

so that $Q(t) := \|\nabla v^m(t, \cdot)\|_{L^2(\mathbb{R}^d)}^2 \left(\int_{\mathbb{R}^d} v^{m+1}(t, x) dx \right)^{-(d-2)/d}$ is **monotone decreasing**, and

$$H' = 2J(S_d Q - 1), \quad H'' = \frac{J'}{J} H' + 2JS_d Q' \leq \frac{J'}{J} H' \leq 0$$

$$H'' \leq -\kappa H' \quad \text{with} \quad \kappa = \frac{2d}{d+2} \frac{1}{S_d} \left(\int_{\mathbb{R}^d} v_0^{m+1} dx \right)^{-2/d}$$

By writing that $-H(0) = H(T) - H(0) \leq H'(0)(1 - e^{-\kappa T})/\kappa$ and using the estimate $\kappa T \leq d/2$, the proof is completed \square

$$H_2[v] := \int_{\mathbb{R}^2} (v - \mu) (-\Delta)^{-1} (v - \mu) dx - \frac{1}{4\pi} \int_{\mathbb{R}^2} v \log \left(\frac{v}{\mu} \right) dx$$

With $\mu(x) := \frac{1}{\pi} (1 + |x|^2)^{-2}$. Assume that v is a positive solution of

$$\frac{\partial v}{\partial t} = \Delta \log(v/\mu) \quad t > 0, \quad x \in \mathbb{R}^2$$

Proposition

If $v = \mu e^{u/2}$ is a solution with nonnegative initial datum v_0 in $L^1(\mathbb{R}^2)$ such that $\int_{\mathbb{R}^2} v_0 dx = 1$, $v_0 \log v_0 \in L^1(\mathbb{R}^2)$ and $v_0 \log \mu \in L^1(\mathbb{R}^2)$, then

$$\begin{aligned} \frac{d}{dt} H_2[v(t, \cdot)] &= \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla u|^2 dx - \int_{\mathbb{R}^2} (e^{\frac{u}{2}} - 1) u d\mu \\ &\geq \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla u|^2 dx + \int_{\mathbb{R}^2} u d\mu - \log \left(\int_{\mathbb{R}^2} e^u d\mu \right) \geq 0 \end{aligned}$$

C: Keller-Segel model

- 1 Small mass results
- 2 Spectral analysis
- 3 Collecting estimates: towards exponential convergence

Keller-Segel with subcritical mass in self-similar variables

$$\begin{cases} \frac{\partial n}{\partial t} = \Delta n - \nabla \cdot (n(\nabla c - x)) & x \in \mathbb{R}^2, t > 0 \\ c = -\frac{1}{2\pi} \log |\cdot| * n & x \in \mathbb{R}^2, t > 0 \\ n(\cdot, t = 0) = n_0 \geq 0 & x \in \mathbb{R}^2 \end{cases}$$

$$\lim_{t \rightarrow \infty} \|n(\cdot, \cdot + t) - n_\infty\|_{L^1(\mathbb{R}^2)} = 0 \quad \text{and} \quad \lim_{t \rightarrow \infty} \|\nabla c(\cdot, \cdot + t) - \nabla c_\infty\|_{L^2(\mathbb{R}^2)} = 0$$

$$n_\infty = M \frac{e^{c_\infty - |x|^2/2}}{\int_{\mathbb{R}^2} e^{c_\infty - |x|^2/2} dx} = -\Delta c_\infty, \quad c_\infty = -\frac{1}{2\pi} \log |\cdot| * n_\infty$$

A parametrization of the solutions and the linearized operator

[Campos, JD]

$$-\Delta c = M \frac{e^{-\frac{1}{2}|x|^2+c}}{\int_{\mathbb{R}^2} e^{-\frac{1}{2}|x|^2+c} dx}$$

Solve

$$-\phi'' - \frac{1}{r} \phi' = e^{-\frac{1}{2}r^2+\phi}, \quad r > 0$$

with initial conditions $\phi(0) = a$, $\phi'(0) = 0$ and get

$$M(a) := 2\pi \int_{\mathbb{R}^2} e^{-\frac{1}{2}r^2+\phi_a} dx$$

$$n_a(x) = M(a) \frac{e^{-\frac{1}{2}r^2+\phi_a(r)}}{2\pi \int_{\mathbb{R}^2} r e^{-\frac{1}{2}r^2+\phi_a} dx} = e^{-\frac{1}{2}r^2+\phi_a(r)}$$

With $-\Delta \varphi_f = n_a f$, consider the operator defined by

$$\mathcal{L}f := \frac{1}{n_a} \nabla \cdot (n_a (\nabla (f - \varphi_f))), \quad x \in \mathbb{R}^2$$

Spectrum of \mathcal{L} (lowest eigenvalues only)

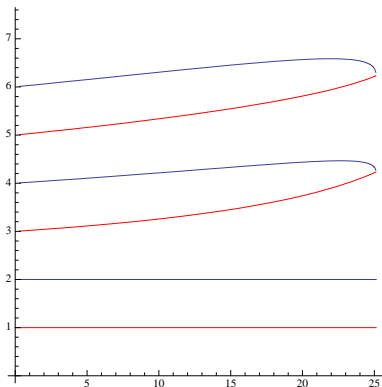


Figure: The lowest eigenvalues of $-\mathcal{L}$ (shown as a function of the mass) are 0, 1 and 2, thus establishing that the spectral gap of $-\mathcal{L}$ is 1

Simple eigenfunctions

Kernel Let $f_0 = \frac{\partial}{\partial M} c_\infty$ be the solution of

$$-\Delta f_0 = n_\infty f_0$$

and observe that $g_0 = f_0/c_\infty$ is such that

$$\frac{1}{n_\infty} \nabla \cdot (n_\infty \nabla (f_0 - c_\infty g_0)) =: \mathcal{L} f_0 = 0$$

Lowest non-zero eigenvalues $f_1 := \frac{1}{n_\infty} \frac{\partial n_\infty}{\partial x_1}$ associated with $g_1 = \frac{1}{c_\infty} \frac{\partial c_\infty}{\partial x_1}$ is an eigenfunction of \mathcal{L} , such that $-\mathcal{L} f_1 = f_1$

With $D := x \cdot \nabla$, let $f_2 = 1 + \frac{1}{2} D \log n_\infty = 1 + \frac{1}{2 n_\infty} D n_\infty$. Then

$$-\Delta (D c_\infty) + 2 \Delta c_\infty = D n_\infty = 2 (f_2 - 1) n_\infty$$

and so $g_2 := \frac{1}{c_\infty} (-\Delta)^{-1} (n_\infty f_2)$ is such that $-\mathcal{L} f_2 = 2 f_2$

$$F[n] := \int_{\mathbb{R}^2} n \log \left(\frac{n}{n_\infty} \right) dx - \frac{1}{2} \int_{\mathbb{R}^2} (n - n_\infty)(c - c_\infty) dx$$

achieves its minimum for $n = n_\infty$ according to log HLS and

$$Q_1[f] = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon^2} F[n_\infty(1 + \varepsilon f)] \geq 0$$

if $\int_{\mathbb{R}^2} f n_\infty dx = 0$. Notice that f_0 generates the kernel of Q_1

Lemma

For any $f \in H^1(\mathbb{R}^2, n_\infty dx)$ such that $\int_{\mathbb{R}^2} f n_\infty dx = 0$, we have

$$\int_{\mathbb{R}^2} |\nabla(g c_\infty)|^2 n_\infty dx \leq \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

... and eigenvalues

With g such that $-\Delta(g c_\infty) = f n_\infty$, Q_1 determines a scalar product

$$\langle f_1, f_2 \rangle := \int_{\mathbb{R}^2} f_1 f_2 n_\infty dx - \int_{\mathbb{R}^2} f_1 n_\infty (g_2 c_\infty) dx$$

on the orthogonal to f_0 in $L^2(n_\infty dx)$ and with $G_2(x) := -\frac{1}{2\pi} \log|x|$

$$Q_2[f] := \int_{\mathbb{R}^2} |\nabla(f - g c_\infty)|^2 n_\infty dx \quad \text{with} \quad g = \frac{1}{c_\infty} G_2 * (f n_\infty)$$

is a positive quadratic form, whose polar operator is the self-adjoint operator \mathcal{L}

$$\langle f, \mathcal{L} f \rangle = Q_2[f] \quad \forall f \in \mathcal{D}(\mathcal{L}_2)$$

Lemma

In this setting, \mathcal{L} has pure discrete spectrum and its lowest eigenvalue is positive

An interpolation inequality induced by log HLS

Lemma

For any $f \in L^2(\mathbb{R}^2, n_\infty dx)$ such that $\int_{\mathbb{R}^2} f f_0 n_\infty dx = 0$ holds, we have

$$-\frac{1}{2\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) n_\infty(x) \log|x-y| f(y) n_\infty(y) dx dy \leq \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

where $g c_\infty = G_2 * (f n_\infty)$ and, if $\int_{\mathbb{R}^2} f n_\infty dx = 0$ holds,

$$\int_{\mathbb{R}^2} |\nabla(g c_\infty)|^2 dx \leq \int_{\mathbb{R}^2} |f|^2 n_\infty dx$$

Equalities in the above inequality then holds if and only if $f = 0$

A new Onofri type inequality

- The spectral gap inequality of \mathcal{L} is a refined version of

Theorem (Onofri type inequality)

For any $M \in (0, 8\pi)$, if $n_\infty = M \frac{e^{c_\infty - \frac{1}{2}|x|^2}}{\int_{\mathbb{R}^2} e^{c_\infty - \frac{1}{2}|x|^2} dx}$ with $c_\infty = (-\Delta)^{-1} n_\infty$, $d\mu_M = \frac{1}{M} n_\infty dx$, we have the inequality

$$\log \left(\int_{\mathbb{R}^2} e^\phi d\mu_M \right) - \int_{\mathbb{R}^2} \phi d\mu_M \leq \frac{1}{2M} \int_{\mathbb{R}^2} |\nabla \phi|^2 dx \quad \forall \phi \in \mathcal{D}_0^{1,2}(\mathbb{R}^2)$$

Back to Keller-Segel: exponential convergence for any mass $M \leq 8\pi$

- [Campos, JD] Uniform convergence of $n(t, \cdot)$ to n_∞ can be established for any $M \in (0, 8\pi)$ by an adaptation of the symmetrization techniques of [Diaz, Nagai, Rakotoson]
- Spectral gap of \mathcal{L} can be established (Persson's lemma or concentration-compactness methods). An improved interpolation inequality also holds:

$$\begin{aligned} -\frac{1}{2\pi} \iint_{\mathbb{R}^2 \times \mathbb{R}^2} f(x) n_\infty(x) \log|x-y| f(y) n_\infty(y) dx dy \\ \leq (1-\varepsilon) \int_{\mathbb{R}^2} |f|^2 n_\infty dx \end{aligned}$$

- Exponential convergence of the relative entropy follows [Campos, JD, work in progress] but estimates are delicate

Thank you for your attention !