Improved Sobolev inequalities using nonlinear flows

Jean Dolbeault

http://www.ceremade.dauphine.fr/~dolbeaul

Ceremade, Université Paris-Dauphine

September 30, 2011

Toulouse *EVOL workshop* (2011, Sep 28 – 30)

A question by H. Brezis and E. Lieb

[Brezis, Lieb (1985)] Is there a natural way to bound

$$S_d \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 - \|u\|_{L^{2^*}(\mathbb{R}^d)}^2$$

from below in terms of the "distance" off from the set of optimal [Aubin-Talenti] functions when $d\geq 3$?

 \bullet [Bianchi-Egnell (1990)] There is a positive constant α such that

$$\mathsf{S}_d \|\nabla u\|_{\mathrm{L}^2(\mathbb{R}^d)}^2 - \|u\|_{\mathrm{L}^{2^*}(\mathbb{R}^d)}^2 \ge \alpha \inf_{\varphi \in \mathcal{M}} \|\nabla u - \nabla \varphi\|_{\mathrm{L}^2(\mathbb{R}^d)}^2$$

• [Cianchi, Fusco, Maggi, Pratelli (2009)] (also a version for $\|\nabla u\|_{\mathrm{L}^p(\mathbb{R}^d)}^p$) There are constants α and κ such that

$$|S_d| |\nabla u||^2_{L^2(\mathbb{R}^d)} \ge (1 + \kappa \lambda(u)^\alpha) ||u||^2_{L^{2^*}(\mathbb{R}^d)}$$

where
$$\lambda(u) = \inf_{\varphi \in \mathcal{M}} \left\{ \frac{\|u - \varphi\|_{L^{2^*}(\mathbb{R}^d)}^{2^*}}{\|u\|_{L^{2^*}(\mathbb{R}^d)}^{2^*}} : \|u\|_{L^{2^*}(\mathbb{R}^d)}^{2^*} = \|\varphi\|_{L^{2^*}(\mathbb{R}^d)}^{2^*} \right\}$$

A – Sobolev and
Hardy-Littlewood-Sobolev
inequalities:
duality, flows

Outline

Outline

- A result motivated by [Carrillo, Carlen and Loss]
- Sobolev and HLS inequalities can be related using a nonlinear flow *compatible with Legendre's duality*
- The asymptotic behaviour close to the *vanishing time* is determined by a solution with *separation of variables* based on the Aubin-Talenti solution
- The entropy H (to be defined) is negative, concave, and we can relate H(0) with H'(0) by integrating estimates on (0, T), which provides a first improvement of Sobolev's inequality if $d \geq 5$

Sobolev and HLS

As it has been noticed by E. Lieb, Sobolev's inequality in \mathbb{R}^d , $d \geq 3$,

$$||u||_{L^{2^*}(\mathbb{R}^d)}^2 \le S_d ||\nabla u||_{L^2(\mathbb{R}^d)}^2 \quad \forall \ u \in \mathcal{D}^{1,2}(\mathbb{R}^d)$$
 (1)

and the Hardy-Littlewood-Sobolev inequality

$$\mathsf{S}_d \|v\|_{\mathrm{L}^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \ge \int_{\mathbb{R}^d} v(-\Delta)^{-1} v \, dx \quad \forall \, v \in \mathrm{L}^{\frac{2d}{d+2}}(\mathbb{R}^d)$$
 (2)

are dual of each other. Here S_d is the Aubin-Talenti constant and $2^* = \frac{2\,d}{d-2}$

Using a nonlinear flow to relate Sobolev and HLS

Consider the fast diffusion equation

$$\frac{\partial v}{\partial t} = \Delta v^m \quad t > 0 \;, \quad x \in \mathbb{R}^d \tag{3}$$

If we define $H(t) := H_d[v(t, \cdot)]$, with

$$\mathsf{H}_d[v] := \int_{\mathbb{R}^d} v (-\Delta)^{-1} v \ dx - \mathsf{S}_d \|v\|_{\mathrm{L}^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2$$

then we observe that

$$\frac{1}{2} \mathsf{H}' = - \int_{\mathbb{R}^d} \mathsf{v}^{m+1} \; d\mathsf{x} + \mathsf{S}_d \left(\int_{\mathbb{R}^d} \mathsf{v}^{\frac{2d}{d+2}} \; d\mathsf{x} \right)^{\frac{1}{d}} \int_{\mathbb{R}^d} \nabla \mathsf{v}^m \cdot \nabla \mathsf{v}^{\frac{d-2}{d+2}} \; d\mathsf{x}$$

where $v = v(t, \cdot)$ is a solution of (3). With the choice $m = \frac{d-2}{d+2}$, we find that $m+1 = \frac{2d}{d+2}$

A first statement

Proposition

[J.D.] Assume that $d \ge 3$ and $\frac{d-2}{d+2}$. If v is a solution of (3) with nonnegative initial datum in $L^{2d/(d+2)}(\mathbb{R}^d)$, then

$$\frac{1}{2} \frac{d}{dt} \left[\int_{\mathbb{R}^d} v (-\Delta)^{-1} v \, dx - \mathsf{S}_d \|v\|_{\mathbf{L}^{\frac{2d}{d+2}}(\mathbb{R}^d)}^2 \right] \\
= \left(\int_{\mathbb{R}^d} v^{m+1} \, dx \right)^{\frac{2}{d}} \left[\mathsf{S}_d \|\nabla u\|_{\mathbf{L}^2(\mathbb{R}^d)}^2 - \|u\|_{\mathbf{L}^{2^*}(\mathbb{R}^d)}^2 \right] \ge 0$$

The HLS inequality amounts to $H \leq 0$ and appears as a consequence of Sobolev, that is $H' \geq 0$ if we show that $\limsup_{t > 0} H(t) = 0$ Notice that $u = v^m$ is an optimal function for (1) if v is optimal for (2) By integrating along the flow defined by (3), we can actually obtain optimal integral remainder terms which improve on the usual Sobolev inequality (1), but only when $d \geq 5$ for integrability reasons

Theorem

[J.D.] Assume that $d \geq 5$ and let $q = \frac{d+2}{d-2}$. There exists a positive constant $\mathcal{C} \leq \left(1 + \frac{2}{d}\right) \left(1 - e^{-d/2}\right) \mathsf{S}_d$ such that

$$||S_{d}||w^{q}||_{L^{\frac{2d}{d+2}}(\mathbb{R}^{d})}^{2} - \int_{\mathbb{R}^{d}} w^{q} (-\Delta)^{-1} w^{q} dx$$

$$\leq C ||w||_{L^{2^{*}}(\mathbb{R}^{d})}^{\frac{8}{d-2}} \left[||\nabla w||_{L^{2}(\mathbb{R}^{d})}^{2} - |S_{d}||w||_{L^{2^{*}}(\mathbb{R}^{d})}^{2} \right]$$

for any $w \in \mathcal{D}^{1,2}(\mathbb{R}^d)$

Solutions with separation of variables

Consider the solution vanishing at t = T:

$$\overline{v}_T(t,x) = c (T-t)^{\alpha} (F(x))^{\frac{d+2}{d-2}} \quad \forall (t,x) \in (0,T) \times \mathbb{R}^d$$

where $\alpha = (d+2)/4$, $c^{1-m} = 4 \, m \, d$, $m = \frac{d-2}{d+2}$, p = d/(d-2) and F is the Aubin-Talenti solution of

$$-\Delta F = d(d-2) F^{(d+2)/(d-2)}$$

Let
$$||v||_* := \sup_{x \in \mathbb{R}^d} (1 + |x|^2)^{d+2} |v(x)|$$

Lemma

[M. delPino, M. Saez], [J. L. Vázquez, J. R. Esteban, A. Rodríguez] For any solution v of (3) with initial datum $v_0 \in L^{2d/(d+2)}(\mathbb{R}^d)$, $v_0 > 0$, there exists T > 0, $\lambda > 0$ and $x_0 \in \mathbb{R}^d$ such that

$$\lim_{t \to T_{-}} (T-t)^{-\frac{1}{1-m}} \|v(t,\cdot)/\overline{v}(t,\cdot) - 1\|_{*} = 0$$

with
$$\overline{v}(t,x) = \lambda^{(d+2)/2} \overline{v}_T(t,(x-x_0)/\lambda)$$

Improved inequality: proof (1/2)

$$\mathsf{J}(t) := \int_{\mathbb{R}^d} v(t,x)^{m+1} \; dx$$
 satisfies

$$J' = -(m+1) \|\nabla v^m\|_{L^2(\mathbb{R}^d)}^2 \le -\frac{m+1}{S_d} J^{1-\frac{2}{d}}$$

If $d \geq 5$, then we also have

$$J'' = 2 m(m+1) \int_{\mathbb{R}^d} v^{m-1} (\Delta v^m)^2 dx \ge 0$$

Such an estimate makes sense if $v = \overline{v}_T$. This is also true for any solution v as can be seen by rewriting the problem on \mathbb{S}^d : integrability conditions for v are exactly the same as for \overline{v}_T

Notice that

$$\frac{\mathsf{J}'}{\mathsf{J}} \le -\frac{m+1}{\mathsf{S}_d} \, \mathsf{J}^{-\frac{2}{d}} \le -\kappa \quad \text{with} \quad \kappa \, T = \frac{2\,d}{d+2} \, \frac{T}{\mathsf{S}_d} \left(\int_{\mathbb{R}^d} v_0^{m+1} \, dx \right)^{-\frac{2}{d}} \le \frac{d}{2}$$

Improved inequality: proof (2/2)

By the Cauchy-Schwarz inequality, we have

$$\frac{J'^{2}}{m+1} = \|\nabla v^{m}\|_{L^{2}(\mathbb{R}^{d})}^{4} = \left(\int_{\mathbb{R}^{d}} v^{(m-1)/2} \, \Delta v^{m} \cdot v^{(m+1)/2} \, dx\right)^{2} \\
\leq \int_{\mathbb{R}^{d}} v^{m-1} \, (\Delta v^{m})^{2} \, dx \int_{\mathbb{R}^{d}} v^{m+1} \, dx = Cst \, J'' \, J$$

so that $Q(t) := \|\nabla v^m(t,\cdot)\|_{L^2(\mathbb{R}^d)}^2 \left(\int_{\mathbb{R}^d} v^{m+1}(t,x) \ dx\right)^{-(d-2)/d}$ is monotone decreasing, and

$$H' = 2 J(S_d Q - 1), \quad H'' = \frac{J'}{J} H' + 2 J S_d Q' \le \frac{J'}{J} H' \le 0$$

$$H'' \le -\kappa H' \quad \text{with} \quad \kappa = \frac{2 d}{d + 2} \frac{1}{S_d} \left(\int_{\mathbb{R}^d} v_0^{m+1} dx \right)^{-2/d}$$

By writing that $-H(0) = H(T) - H(0) \le H'(0) (1 - e^{-\kappa T})/\kappa$ and using the estimate $\kappa T \le d/2$, the proof is completed



d = 2: Onofri's and log HLS inequalities

$$\mathsf{H}_2[v] := \int_{\mathbb{R}^2} \left(v - \mu\right) (-\Delta)^{-1} (v - \mu) \; dx - \frac{1}{4 \, \pi} \int_{\mathbb{R}^2} v \, \log \left(\frac{v}{\mu}\right) \, dx$$

With $\mu(x) := \frac{1}{\pi} (1 + |x|^2)^{-2}$. Assume that ν is a positive solution of

$$\frac{\partial v}{\partial t} = \Delta \log \left(\frac{v}{\mu}\right) \quad t > 0 , \quad x \in \mathbb{R}^2$$

Proposition

If $v=\mu\,e^{u/2}$ is a solution with nonnegative initial datum v_0 in $L^1(\mathbb{R}^2)$ such that $\int_{\mathbb{R}^2} v_0\ dx=1$, $v_0\log v_0\in L^1(\mathbb{R}^2)$ and $v_0\log \mu\in L^1(\mathbb{R}^2)$, then

$$\begin{split} \frac{d}{dt}\mathsf{H}_2[v(t,\cdot)] &= \frac{1}{16\,\pi} \int_{\mathbb{R}^2} |\nabla u|^2 \; dx - \int_{\mathbb{R}^2} \left(e^{\frac{u}{2}} - 1\right) u \; d\mu \\ &\geq \frac{1}{16\,\pi} \int_{\mathbb{R}^2} |\nabla u|^2 \; dx + \int_{\mathbb{R}^2} u \; d\mu - \log\left(\int_{\mathbb{R}^2} e^u \; d\mu\right) \geq 0 \end{split}$$



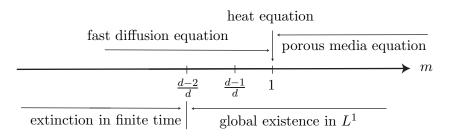
Fast diffusion equations

- entropy methods
- linearization of the entropy
- improved Gagliardo-Nirenberg inequalities

B1 – Fast diffusion equations: entropy methods

Existence, classical results

$$\begin{split} u_t &= \Delta u^m \quad x \in \mathbb{R}^d \,, \ t > 0 \\ \text{Self-similar (Barenblatt) function: } \mathcal{U}(t) &= O(t^{-d/(2-d(1-m))}) \text{ as } \\ t &\to +\infty \\ \text{[Friedmann, Kamin, 1980] } \|u(t,\cdot) - \mathcal{U}(t,\cdot)\|_{L^\infty} &= o(t^{-d/(2-d(1-m))}) \end{split}$$



Existence theory, critical values of the parameter m

Time-dependent rescaling, Free energy

Time-dependent rescaling: Take $u(\tau, y) = R^{-d}(t) v(t, y/R(\tau))$ where

$$\frac{\partial R}{\partial \tau} = R^{d(1-m)-1}$$
, $R(0) = 1$, $t = \log R$

 $lue{}$ The function v solves a Fokker-Planck type equation

$$\frac{\partial v}{\partial t} = \Delta v^m + \nabla \cdot (x \, v) \,, \quad v_{|\tau=0} = u_0$$

• [Ralston, Newman, 1984] Lyapunov functional:

Generalized entropy or **Free energy**

$$\Sigma[v] := \int_{\mathbb{R}^d} \left(\frac{v^m}{m-1} + \frac{1}{2} |x|^2 v \right) dx - \Sigma_0$$

Entropy production is measured by the **Generalized Fisher information**

$$\frac{d}{dt}\Sigma[v] = -I[v] , \quad I[v] := \int_{\mathbb{R}^d} v \left| \frac{\nabla v^m}{v} + x \right|^2 dx$$

Relative entropy and entropy production

Stationary solution: choose C such that $||v_{\infty}||_{L^1} = ||u||_{L^1} = M > 0$

$$v_{\infty}(x) := \left(C + \frac{1-m}{2m}|x|^2\right)_+^{-1/(1-m)}$$

Relative entropy: Fix Σ_0 so that $\Sigma[v_\infty] = 0$. The entropy can be put in an *m*-homogeneous form: for $m \neq 1$,

$$\Sigma[v] = \int_{\mathbb{R}^d} \psi(\frac{v}{v_{\infty}}) \ v_{\infty}^m \ dx \quad \text{with } \psi(t) = \frac{t^m - 1 - m(t - 1)}{m - 1}$$

Entropy – entropy production inequality

Theorem

$$d \geq 3, \ m \in [\frac{d-1}{d}, +\infty), \ m > \frac{1}{2}, \ m \neq 1$$

$$I[v] \geq 2\Sigma[v]$$

Corollary

A solution v with initial data $u_0 \in L^1_+(\mathbb{R}^d)$ such that $|x|^2 u_0 \in L^1(\mathbb{R}^d)$, $u_0^m \in L^1(\mathbb{R}^d)$ satisfies $\sum [v(t,\cdot)] \leq \sum [u_0] e^{-2t}$

An equivalent formulation: Gagliardo-Nirenberg inequalities

$$\sum [v] = \int_{\mathbb{R}^d} \left(\frac{v^m}{m-1} + \frac{1}{2} |x|^2 v \right) dx - \sum_0 \le \frac{1}{2} \int_{\mathbb{R}^d} v \left| \frac{\nabla v^m}{v} + x \right|^2 dx = \frac{1}{2} I[v]$$

Rewrite it with $p = \frac{1}{2m-1}$, $v = w^{2p}$, $v^m = w^{p+1}$ as

$$\frac{1}{2} \left(\frac{2m}{2m-1} \right)^2 \int_{\mathbb{R}^d} |\nabla w|^2 dx + \left(\frac{1}{1-m} - d \right) \int_{\mathbb{R}^d} |w|^{1+p} dx - K \ge 0$$

- for some γ , $K = K_0 \left(\int_{\mathbb{R}^d} v \, dx = \int_{\mathbb{R}^d} w^{2p} \, dx \right)^{\gamma}$
- $w = w_{\infty} = v_{\infty}^{1/2p}$ is optimal

Theorem

[Del Pino, J.D.] With $1 (fast diffusion case) and <math>d \ge 3$

$$\|w\|_{L^{2p}(\mathbb{R}^d)} \le A \|\nabla w\|_{L^2(\mathbb{R}^d)}^{\theta} \|w\|_{L^{p+1}(\mathbb{R}^d)}^{1-\theta}$$

$$A = \left(\frac{y(p-1)^2}{2\pi d}\right)^{\frac{\theta}{2}} \left(\frac{2y-d}{2y}\right)^{\frac{1}{2p}} \left(\frac{\Gamma(y)}{\Gamma(y-\frac{d}{2})}\right)^{\frac{\theta}{d}} , \quad \theta = \frac{d(p-1)}{p(d+2-(d-2)p)} , \quad y = \frac{p+1}{p-1}$$

20 d G

...the Bakry-Emery method

Consider the generalized Fisher information

$$I[v] := \int_{\mathbb{R}^d} v |Z|^2 dx$$
 with $Z := \frac{\nabla v^m}{v} + x$

and compute

$$\frac{d}{dt} I[v(t,\cdot)] + 2 I[v(t,\cdot)] = -2 (m-1) \int_{\mathbb{R}^d} u^m (\operatorname{div} Z)^2 dx - 2 \sum_{i,j=1}^d \int_{\mathbb{R}^d} u^m (\partial_i Z^j)^2 dx$$

• the Fisher information decays exponentially:

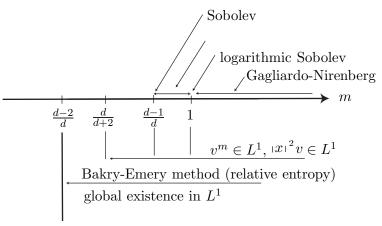
$$I[v(t,\cdot)] \leq I[u_0] e^{-2t}$$

- $\lim_{t\to\infty} I[v(t,\cdot)] = 0$ and $\lim_{t\to\infty} \Sigma[v(t,\cdot)] = 0$
- $\frac{d}{dt} \left(I[v(t,\cdot)] 2 \Sigma[v(t,\cdot)] \right) \le 0 \text{ means } I[v] \ge 2 \Sigma[v]$

[Carrillo, Toscani], [Juengel, Markowich, Toscani], [Carrillo, Juengel, Markowich, Toscani, Unterreiter], [Carrillo, Vázquez]

Fast diffusion: finite mass regime

Inequalities...



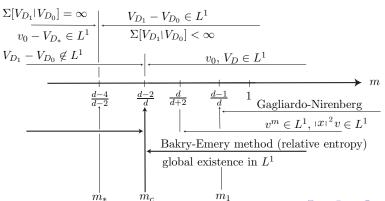
... existence of solutions of
$$u_t = \Delta u^m$$

B2 – Fast diffusion equations: the infinite mass regime Linearization of the entropy

Extension to the infinite mass regime, finite time vanishing

- If $m > m_c := \frac{d-2}{d} \le m < m_1$, solutions globally exist in $L^1(\mathbb{R}^d)$ and the Barenblatt self-similar solution has finite mass.
- \bullet For $m \leq m_c$, the Barenblatt self-similar solution has infinite mass

Extension to $m \le m_c$? Work in relative variables!



Entropy methods and linearization: intermediate asymptotics, vanishing

[A. Blanchet, M. Bonforte, J.D., G. Grillo, J.L. Vázquez]

$$\frac{\partial u}{\partial \tau} = -\nabla \cdot (u \nabla u^{m-1}) = \frac{1-m}{m} \Delta u^m \tag{4}$$

• $m_c < m < 1, \ T = +\infty$: intermediate asymptotics, $\tau \to +\infty$

$$R(\tau) := (T + \tau)^{\frac{1}{d(m-m_c)}}$$

• $0 < m < m_c$, $T < +\infty$: vanishing in finite time $\lim_{\tau \nearrow T} u(\tau, y) = 0$

$$R(\tau) := (T - \tau)^{-\frac{1}{d(m_c - m)}}$$

Self-similar $Barenblatt\ type\ solutions$ exists for any m Rescaling: time-dependent change of variables

$$t := \frac{1-m}{2} \log \left(\frac{R(\tau)}{R(0)} \right)$$
 and $x := \sqrt{\frac{1}{2d |m-m_c|}} \frac{y}{R(\tau)}$

Generalized Barenblatt profiles: $V_D(x) := (D + |x|^2)^{\frac{1}{m-1}}$

Sharp rates of convergence

Assumptions on the initial datum v_0

(H1)
$$V_{D_0} \le v_0 \le V_{D_1}$$
 for some $D_0 > D_1 > 0$

(H2) if $d \ge 3$ and $m \le m_*$, $(v_0 - V_D)$ is integrable for a suitable $D \in [D_1, D_0]$

Theorem

[Bonforte, J.D., Grillo, Vázquez] Under Assumptions (H1)-(H2), if m < 1 and $m \neq m_* := \frac{d-4}{d-2}$, the entropy decays according to

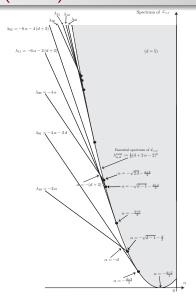
$$\mathcal{E}[v(t,\cdot)] \le C e^{-2(1-m)\Lambda_{\alpha,d}t} \quad \forall \ t \ge 0$$

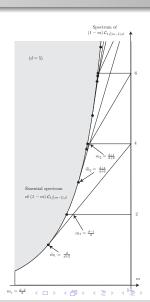
where $\Lambda_{\alpha,d} > 0$ is the best constant in the Hardy–Poincaré inequality

$$\Lambda_{\alpha,d} \int_{\mathbb{R}^d} |f|^2 d\mu_{\alpha-1} \leq \int_{\mathbb{R}^d} |\nabla f|^2 d\mu_{\alpha} \quad \forall \ f \in H^1(d\mu_{\alpha})$$

with $\alpha := 1/(m-1) < 0$, $d\mu_{\alpha} := h_{\alpha} dx$, $h_{\alpha}(x) := (1+|x|^2)^{\alpha}$

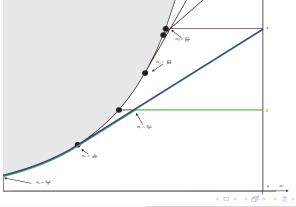
Plots (d = 5)





Improved asymptotic rates

[Bonforte, J.D., Grillo, Vázquez] Assume that $m \in (m_1, 1)$, $d \ge 3$. Under Assumption (H1), if v is a solution of the fast diffusion equation with initial datum v_0 such that $\int_{\mathbb{R}^d} x \, v_0 \, dx = 0$, then the asymptotic convergence holds with an improved rate corresponding to the improved spectral gap.



Higher order matching asymptotics

[J.D., G. Toscani] For some $m \in (m_c, 1)$ with $m_c := (d-2)/d$, we consider on \mathbb{R}^d the fast diffusion equation

$$\frac{\partial u}{\partial \tau} + \nabla \cdot \left(u \, \nabla u^{m-1} \right) = 0$$

The strategy is easy to understand using a time-dependent rescaling and the relative entropy formalism. Define the function ν such that

$$u(\tau, y + x_0) = R^{-d} v(t, x)$$
, $R = R(\tau)$, $t = \frac{1}{2} \log R$, $x = \frac{y}{R}$

Then v has to be a solution of

$$\frac{\partial v}{\partial t} + \nabla \cdot \left[v \left(\sigma^{\frac{d}{2}(m - m_c)} \nabla v^{m-1} - 2x \right) \right] = 0 \quad t > 0 , \quad x \in \mathbb{R}^d$$

with (as long as we make no assumption on R)

$$2\,\sigma^{-\frac{d}{2}(m-m_c)} = R^{1-d\,(1-m)}\,\frac{dR}{d\tau}$$

Refined relative entropy

Consider the family of the Barenblatt profiles

$$B_{\sigma}(x) := \sigma^{-\frac{d}{2}} \left(C_{M} + \frac{1}{\sigma} |x|^{2} \right)^{\frac{1}{m-1}} \quad \forall \ x \in \mathbb{R}^{d}$$
 (5)

Note that σ is a function of t: as long as $\frac{d\sigma}{dt} \neq 0$, the Barenblatt profile B_{σ} is not a solution but we may still consider the relative entropy

$$\mathcal{F}_{\sigma}[v] := \frac{1}{m-1} \int_{\mathbb{R}^d} \left[v^m - B_{\sigma}^m - m B_{\sigma}^{m-1} \left(v - B_{\sigma} \right) \right] dx$$

Let us briefly sketch the strategy of our method before giving all details

The time derivative of this relative entropy is

$$\frac{d}{dt}\mathcal{F}_{\sigma(t)}[v(t,\cdot)] = \underbrace{\frac{d\sigma}{dt}\left(\frac{d}{d\sigma}\mathcal{F}_{\sigma}[v]\right)_{|\sigma=\sigma(t)}}_{\text{choose it}} + \frac{m}{m-1}\int_{\mathbb{R}^d} \left(v^{m-1} - B_{\sigma(t)}^{m-1}\right) \frac{\partial v}{\partial t} dx$$

$$\iff \text{Minimize } \mathcal{F}_{\sigma}[v] \text{ w.r.t. } \sigma \iff \int_{\mathbb{R}^d} |x|^2 B_{\sigma} \ dx = \int_{\mathbb{R}^d} |x|^2 v \ dx$$

The entropy / entropy production estimate

According to the definition of B_{σ} , we know that

$$2x = \sigma^{\frac{d}{2}(m-m_c)} \nabla B_{\sigma}^{m-1}$$

Using the new change of variables, we know that

$$\frac{d}{dt}\mathcal{F}_{\sigma(t)}[v(t,\cdot)] = -\frac{m\,\sigma(t)^{\frac{d}{2}(m-m_c)}}{1-m}\int_{\mathbb{R}^d}v\left|\nabla\left[v^{m-1}-B^{m-1}_{\sigma(t)}\right]\right|^2\,dx$$

Let $w := v/B_{\sigma}$ and observe that the relative entropy can be written as

$$\mathcal{F}_{\sigma}[v] = \frac{m}{1-m} \int_{\mathbb{R}^d} \left[w - 1 - \frac{1}{m} (w^m - 1) \right] B_{\sigma}^m dx$$

(Repeating) define the relative Fisher information by

$$\mathcal{I}_{\sigma}[v] := \int_{\mathbb{R}^d} \left| \frac{1}{m-1} \nabla \left[\left(w^{m-1} - 1 \right) B_{\sigma}^{m-1} \right] \right|^2 B_{\sigma} w \ dx$$

so that
$$\frac{d}{dt}\mathcal{F}_{\sigma(t)}[v(t,\cdot)] = -m(1-m)\frac{\sigma(t)}{\sigma(t)}\mathcal{I}_{\sigma(t)}[v(t,\cdot)] \quad \forall \ t>0$$

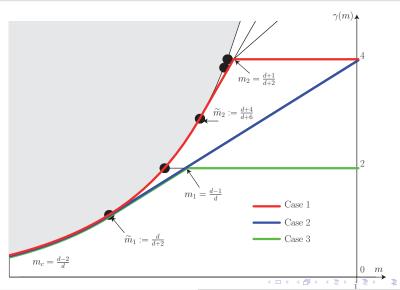
When linearizing, one more mode is killed and $\sigma(t)$ scales out

Theorem (J.D., G. Toscani)

$$\text{Let } m \in (\widetilde{m}_1, 1), \ d \geq 2, \ v_0 \in L^1_+(\mathbb{R}^d) \ \text{such that } v_0^m, \ |y|^2 \ v_0 \in L^1(\mathbb{R}^d)$$

$$\mathcal{E}[v(t, \cdot)] \leq C \ e^{-2\gamma(m)\,t} \quad \forall \ t \geq 0$$
 where
$$\gamma(m) = \begin{cases} \frac{((d-2)\,m - (d-4))^2}{4\,(1-m)} & \text{if } m \in (\widetilde{m}_1, \widetilde{m}_2] \\ 4\,(d+2)\,m - 4\,d & \text{if } m \in [\widetilde{m}_2, m_2] \\ 4 & \text{if } m \in [m_2, 1) \end{cases}$$

Spectral gaps and best constants



Gagliardo-Nirenberg and Sobolev inequalities: improvements

[J.D., G. Toscani]

Best matching Barenblatt profiles

(Repeating) Consider the fast diffusion equation

$$\frac{\partial u}{\partial t} + \nabla \cdot \left[u \left(\sigma^{\frac{d}{2}(m - m_c)} \nabla u^{m-1} - 2x \right) \right] = 0 \quad t > 0 , \quad x \in \mathbb{R}^d$$

with a nonlocal, time-dependent diffusion coefficient

$$\sigma(t) = \frac{1}{K_M} \int_{\mathbb{R}^d} |x|^2 u(x,t) dx, \quad K_M := \int_{\mathbb{R}^d} |x|^2 B_1(x) dx$$

where

$$B_{\lambda}(x) := \lambda^{-\frac{d}{2}} \left(C_M + \frac{1}{\lambda} |x|^2 \right)^{\frac{1}{m-1}} \quad \forall \ x \in \mathbb{R}^d$$

and define the relative entropy

$$\mathcal{F}_{\lambda}[u] := \frac{1}{m-1} \int_{\mathbb{R}^d} \left[u^m - B_{\lambda}^m - m B_{\lambda}^{m-1} \left(u - B_{\lambda} \right) \right] dx$$



Three ingredients for global improvements

① $\inf_{\lambda>0} \mathcal{F}_{\lambda}[u(x,t)] = \mathcal{F}_{\sigma(t)}[u(x,t)]$ so that

$$\frac{d}{dt}\mathcal{F}_{\sigma(t)}[u(x,t)] = -\mathcal{J}_{\sigma(t)}[u(\cdot,t)]$$

where the relative Fisher information is

$$\mathcal{J}_{\lambda}[u] := \lambda^{\frac{d}{2}(m-m_c)} \frac{m}{1-m} \int_{\mathbb{R}^d} u \left| \nabla u^{m-1} - \nabla B_{\lambda}^{m-1} \right|^2 dx$$

② In the Bakry-Emery method, there is an additional (good) term

$$4\left[1+2C_{m,d}\frac{\mathcal{F}_{\sigma(t)}[u(\cdot,t)]}{M^{\gamma}\sigma_{0}^{\frac{d}{2}(1-m)}}\right]\frac{d}{dt}\left(\mathcal{F}_{\sigma(t)}[u(\cdot,t)]\right)\geq \frac{d}{dt}\left(\mathcal{J}_{\sigma(t)}[u(\cdot,t)]\right)$$

The Csiszár-Kullback inequality is also improved

$$\mathcal{F}_{\sigma}[u] \geq \frac{m}{8 \int_{\mathbb{R}^d} B_1^m dx} C_M^2 \|u - B_{\sigma}\|_{\mathrm{L}^1(\mathbb{R}^d)}^2$$



improved decay for the relative entropy

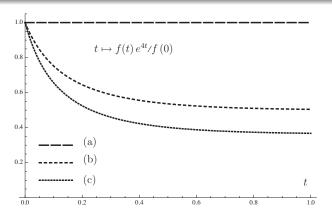


Figure: Upper bounds on the decay of the relative entropy: $t\mapsto f(t)\,e^{4t}/f(0)$

- (a): estimate given by the entropy-entropy production method
- (b): exact solution of a simplified equation
- (c): numerical solution (found by a shooting method)



A Csiszár-Kullback(-Pinsker) inequality

Let $m \in (\widetilde{m}_1, 1)$ with $\widetilde{m}_1 = \frac{d}{d+2}$ and consider the relative entropy

$$\mathcal{F}_{\sigma}[u] := \frac{1}{m-1} \int_{\mathbb{R}^d} \left[u^m - B_{\sigma}^m - m B_{\sigma}^{m-1} \left(u - B_{\sigma} \right) \right] dx$$

$\mathsf{Theorem}$

Let $d \geq 1$, $m \in (\widetilde{m}_1,1)$ and assume that u is a nonnegative function in $L^1(\mathbb{R}^d)$ such that u^m and $x \mapsto |x|^2 u$ are both integrable on \mathbb{R}^d . If $\|u\|_{L^1(\mathbb{R}^d)} = M$ and $\int_{\mathbb{R}^d} |x|^2 u \ dx = \int_{\mathbb{R}^d} |x|^2 B_\sigma \ dx$, then

$$\frac{\mathcal{F}_{\sigma}[u]}{\sigma^{\frac{d}{2}(1-m)}} \geq \frac{m}{8 \int_{\mathbb{R}^d} B_1^m \ dx} \left(C_M \|u - B_{\sigma}\|_{\mathrm{L}^1(\mathbb{R}^d)} + \frac{1}{\sigma} \int_{\mathbb{R}^d} |x|^2 |u - B_{\sigma}| \ dx \right)^2$$

Csiszár-Kullback(-Pinsker): proof (1/2)

Let
$$v:=u/B_\sigma$$
 and $d\mu_\sigma:=B_\sigma^m\,dx$

$$\int_{\mathbb{R}^{d}} (v-1) d\mu_{\sigma} = \int_{\mathbb{R}^{d}} B_{\sigma}^{m-1} (u-B_{\sigma}) dx$$

$$= \sigma^{\frac{d}{2}(1-m)} C_{M} \int_{\mathbb{R}^{d}} (u-B_{\sigma}) dx + \sigma^{\frac{d}{2}(m_{c}-m)} \int_{\mathbb{R}^{d}} |x|^{2} (u-B_{\sigma}) dx = 0$$

$$\int_{\mathbb{R}^{d}} (v-1) d\mu_{\sigma} = \int_{v>1} (v-1) d\mu_{\sigma} - \int_{v<1} (1-v) d\mu_{\sigma} = 0$$

$$\int_{\mathbb{R}^{d}} |v-1| d\mu_{\sigma} = \int_{v>1} (v-1) d\mu_{\sigma} + \int_{v<1} (1-v) d\mu_{\sigma}$$

$$\int_{\mathbb{R}^{d}} |u-B_{\sigma}| B_{\sigma}^{m-1} dx = \int_{\mathbb{R}^{d}} |v-1| d\mu_{\sigma} = 2 \int_{v<1} |v-1| d\mu_{\sigma}$$

Csiszár-Kullback(-Pinsker): proof (2/2)

A Taylor expansion shows that

$$\mathcal{F}_{\sigma}[u] = \frac{1}{m-1} \int_{\mathbb{R}^d} \left[v^m - 1 - m(v-1) \right] d\mu_{\sigma} = \frac{m}{2} \int_{\mathbb{R}^d} \xi^{m-2} |v-1|^2 d\mu_{\sigma}$$
$$\geq \frac{m}{2} \int_{v<1} |v-1|^2 d\mu_{\sigma}$$

Using the Cauchy-Schwarz inequality, we get

$$\left(\int_{v<1} |v-1| \ d\mu_{\sigma}\right)^{2} = \left(\int_{v<1} |v-1| \ B_{\sigma}^{\frac{m}{2}} \ B_{\sigma}^{\frac{m}{2}} \ dx\right)^{2} \leq \int_{v<1} |v-1|^{2} \ d\mu_{\sigma} \int_{\mathbb{R}^{d}} B_{\sigma}^{m} \ dx$$

and finally obtain that

$$\mathcal{F}_{\sigma}[u] \geq \frac{m}{2} \frac{\left(\int_{v < 1} |v - 1| \ d\mu_{\sigma}\right)^2}{\int_{\mathbb{R}^d} B_{\sigma}^m \ dx} = \frac{m}{8} \frac{\left(\int_{\mathbb{R}^d} |u - B_{\sigma}| \ B_{\sigma}^{m-1} \ dx\right)^2}{\int_{\mathbb{R}^d} B_{\sigma}^m \ dx}$$

An improved Gagliardo-Nirenberg inequality: the setting

The inequality

$$\|f\|_{\mathrm{L}^{2p}(\mathbb{R}^d)} \leq \mathcal{C}_{p,d}^{\mathrm{GN}} \, \|\nabla f\|_{\mathrm{L}^2(\mathbb{R}^d)}^{\theta} \, \|f\|_{\mathrm{L}^{p+1}(\mathbb{R}^d)}^{1-\theta}$$

with $\theta = \theta(p) := \frac{p-1}{p} \frac{d}{d+2-p(d-2)}$, $1 if <math>d \ge 3$ and 1 if <math>d = 2, can be rewritten, in a non-scale invariant form, as

$$\int_{\mathbb{R}^d} |\nabla f|^2 dx + \int_{\mathbb{R}^d} |f|^{p+1} dx \ge \mathsf{K}_{p,d} \left(\int_{\mathbb{R}^d} |f|^{2p} dx \right)^{\gamma}$$

with $\gamma = \gamma(p, d) := \frac{d+2-p(d-2)}{d-p(d-4)}$. Optimal function are given by

$$f_{M,y,\sigma}(x) = \frac{1}{\sigma^{\frac{d}{2}}} \left(C_M + \frac{|x-y|^2}{\sigma} \right)^{-\frac{1}{p-1}} \quad \forall \ x \in \mathbb{R}^d$$

where C_M is determined by $\int_{\mathbb{R}^d} f_{M,\gamma,\sigma}^{2p} dx = M$

$$\mathfrak{M}_d := \left\{ f_{M,y,\sigma} : (M,y,\sigma) \in \mathcal{M}_d := (0,\infty) \times \mathbb{R}^d \times (0,\infty) \right\}$$

An improved Gagliardo-Nirenberg inequality (1/2)

Relative entropy functional

$$\mathcal{R}^{(p)}[f] := \inf_{g \in \mathfrak{M}_d^{(p)}} \int_{\mathbb{R}^d} \left[g^{1-p} \left(|f|^{2p} - g^{2p} \right) - \frac{2p}{p+1} \left(|f|^{p+1} - g^{p+1} \right) \right] dx$$

$\mathsf{Theorem}$

Let $d \ge 2$, p > 1 and assume that p < d/(d-2) if $d \ge 3$. If

$$\frac{\int_{\mathbb{R}^d} |x|^2 |f|^{2p} \ dx}{\left(\int_{\mathbb{R}^d} |f|^{2p} \ dx\right)^{\gamma}} = \frac{d(p-1) \sigma_* M_*^{\gamma-1}}{d+2-p(d-2)}, \ \sigma_*(p) := \left(4 \frac{d+2-p(d-2)}{(p-1)^2(p+1)}\right)^{\frac{4p}{d-p(d-4)}}$$

for any $f \in \mathrm{L}^{p+1} \cap \mathcal{D}^{1,2}(\mathbb{R}^d)$, then we have

$$\int_{\mathbb{R}^{d}} |\nabla f|^{2} dx + \int_{\mathbb{R}^{d}} |f|^{p+1} dx - \mathsf{K}_{p,d} \left(\int_{\mathbb{R}^{d}} |f|^{2p} dx \right)^{\gamma} \ge \mathsf{C}_{p,d} \frac{\left(\mathcal{R}^{(p)}[f] \right)^{2}}{\left(\int_{\mathbb{R}^{d}} |f|^{2p} dx \right)^{\gamma}}$$

Improved Sobolev inequalities

An improved Gagliardo-Nirenberg inequality (2/2)

A Csiszár-Kullback inequality

$$\mathcal{R}^{(p)}[f] \ge \mathsf{C}_{\mathrm{CK}} \|f\|_{\mathrm{L}^{2p}(\mathbb{R}^d)}^{2p(\gamma-2)} \inf_{g \in \mathfrak{M}_d^{(p)}} \||f|^{2p} - g^{2p}\|_{\mathrm{L}^1(\mathbb{R}^d)}^2$$

with
$$C_{\text{CK}} = \frac{p-1}{p+1} \frac{d+2-p(d-2)}{32p} \sigma_*^{d \frac{p-1}{4p}} M_*^{1-\gamma}$$
. Let

$$\mathfrak{C}_{p,d} := \mathsf{C}_{d,p} \, \mathsf{C}_{\mathrm{CK}}^{2}$$

Corollary

Under previous assumptions, we have

$$\int_{\mathbb{R}^{d}} |\nabla f|^{2} dx + \int_{\mathbb{R}^{d}} |f|^{p+1} dx - \mathsf{K}_{p,d} \left(\int_{\mathbb{R}^{d}} |f|^{2p} dx \right)^{\gamma} \\
\geq \mathfrak{C}_{p,d} \|f\|_{L^{2p}(\mathbb{R}^{d})}^{2p(\gamma-4)} \inf_{g \in \mathfrak{M}_{d}(p)} \||f|^{2p} - g^{2p}\|_{L^{1}(\mathbb{R}^{d})}^{4}$$

Conclusion 1: improved inequalities

- We have found an improvement of an optimal Gagliardo-Nirenberg inequality, which provides an explicit measure of the distance to the manifold of optimal functions.
- The method is based on the nonlinear flow
- $\ \, \square$ The explicit improvement gives (is equivalent to) an improved entropy entropy production inequality

Conclusion 2: improved rates

If $m \in (m_1, 1)$, with

$$f(t) := \mathcal{F}_{\sigma(t)}[u(\cdot, t)]$$

$$\sigma(t) = \frac{1}{K_M} \int_{\mathbb{R}^d} |x|^2 u(x, t) dx$$

$$j(t) := \mathcal{J}_{\sigma(t)}[u(\cdot, t)]$$

$$\mathcal{J}_{\sigma}[u] := \frac{m \, \sigma^{\frac{d}{2}(m-m_c)}}{1-m} \int_{\mathbb{R}^d} u \, \left| \nabla u^{m-1} - \nabla \mathfrak{B}_{\sigma}^{m-1} \right|^2 \, dx$$

we can write a system of coupled ODEs

$$\begin{cases}
f' = -j \leq 0 \\
\sigma' = -2 d \frac{(1-m)^2}{m K_M} \sigma^{\frac{d}{2}(m-m_c)} f \leq 0 \\
j' + 4 j = \frac{d}{2} (m - m_c) \left[\frac{j}{\sigma} + 4 d (1-m) \frac{f}{\sigma} \right] \sigma' - r
\end{cases}$$
(6)

In the rescaled variables, we have found an *improved decay* (algebraic rate) of the relative entropy. This is a new nonlinear effect, which matters for the initial time layer

Conclusion 3: Best matching Barenblatt profiles are delayed

Let *u* be such that

$$v(\tau, x) = \frac{\mu^d}{R(D\tau)^d} u\left(\frac{1}{2}\log R(D\tau), \frac{\mu x}{R(D\tau)}\right)$$

with $\tau \mapsto R(\tau)$ given as the solution to

$$\frac{1}{R} \frac{dR}{d\tau} = \left(\frac{\mu^2}{K_M} \int_{\mathbb{R}^d} |x|^2 \, v(\tau, x) \, dx \right)^{-\frac{\sigma}{2}(m - m_c)}, \quad R(0) = 1$$

Then

$$\frac{1}{R}\frac{dR}{d\tau} = \left[R^2(\tau)\,\sigma\left(\frac{1}{2}\log R(D\,\tau)\right)\right]^{-\frac{d}{2}(m-m_c)}$$

that is $R(\tau) = R_0(\tau) \le R_0(\tau)$ where $\frac{1}{R} \frac{dR_0}{d\tau} = (R_0^2(\tau) \sigma(0))^{-\frac{d}{2}(m-m_c)}$ and asymptotically as $\tau \to \infty$, $R(\tau) = R_0(\tau - \delta)$ for some delay $\delta > 0$

A: Sobolev and Hardy-Littlewood-Sobolev inequalities: duality, flows B1: Fast diffusion equations: entropy methods B2: Fast diffusion equations: linearization of the entropy Gagliardo-Nirenberg inequalities: improvements

Thank you for your attention!