

Direct entropy methods

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I — Entropy

The logarithmic Sobolev inequality

Convex Sobolev inequalities — Applications

Entropy - Entropy production method

Intermediate asymptotics : heat equation

- *probability theory*: [Bakry], [Emery], [Ledoux], [Coulhon],...
- *linear diffusions*: [Toscani], [Arnold, Markowich, Toscani, Unterreiter], [Otto, Kinderlehrer, Jordan]
- *nonlinear diffusions*: [Carrillo, Toscani], [Del Pino, J.D.], [Otto], [Juengel, Markowich, Toscani], [Carrillo, Juengel, Markowich, Toscani, Unterreiter], [Markowich, Lederman], [Carrillo, Vazquez]

Heat equation:
$$\begin{cases} u_t = \Delta u & x \in \mathbb{R}^n, t \in \mathbb{R}^+ \\ u(\cdot, t=0) = u_0 \geq 0 & \int_{\mathbb{R}^n} u_0 \, dx = 1 \end{cases} \quad (1)$$

As $t \rightarrow +\infty$, $u(x, t) \sim \mathcal{U}(x, t) = \frac{e^{-|x|^2/4t}}{(4\pi t)^{n/2}}$. What is the (optimal) rate of convergence of $\|u(\cdot, t) - \mathcal{U}(\cdot, t)\|_{L^1(\mathbb{R}^n)}$?

The time dependent rescaling

$$u(x, t) = \frac{1}{R^n(t)} v \left(\xi = \frac{x}{R(t)}, \tau = \log R(t) + \tau(0) \right)$$

allows to transform this question into that of the convergence to the stationary solution $v_\infty(\xi) = (2\pi)^{n/2} e^{-|\xi|^2/2}$.

- Ansatz: $\frac{dR}{dt} = \frac{1}{R}$ $R(0) = 1$ $\tau(0) = 0$:

$$R(t) = \sqrt{1 + 2t} \quad \tau(t) = \log R(t)$$

As a consequence: $v(\tau = 0) = u_0$.

- Fokker-Planck equation:

$$\begin{cases} v_\tau = \Delta v + \nabla(\xi v) & \xi \in \mathbb{R}^n, \tau \in \mathbb{R}^+ \\ v(\cdot, \tau = 0) = u_0 \geq 0 & \int_{\mathbb{R}^n} u_0 \, dx = 1 \end{cases} \quad (2)$$

Entropy (relative to the stationary solution v_∞):

$$\Sigma[v] := \int_{\mathbf{R}^n} v \log \left(\frac{v}{v_\infty} \right) dx$$

If v is a solution of (2), then (I is the Fisher information)

$$\frac{d}{d\tau} \Sigma[v(\cdot, \tau)] = - \int_{\mathbf{R}^n} v \left| \nabla \log \left(\frac{v}{v_\infty} \right) \right|^2 dx =: -I[v(\cdot, \tau)]$$

- Euclidean logarithmic Sobolev inequality: If $\|v\|_{L^1} = 1$, then

$$\int_{\mathbf{R}^n} v \log v dx + n \left(1 + \frac{1}{2} \log(2\pi) \right) \leq \frac{1}{2} \int_{\mathbf{R}^n} \frac{|\nabla v|^2}{v} dx$$

Equality: $v(\xi) = v_\infty(\xi) = (2\pi)^{-n/2} e^{-|\xi|^2/2}$

$$\implies \Sigma[v(\cdot, \tau)] \leq \frac{1}{2} I[v(\cdot, \tau)]$$

$$\Sigma[v(\cdot, \tau)] \leq e^{-2\tau} \Sigma[u_0] = e^{-2\tau} \int_{\mathbf{R}^n} u_0 \log \left(\frac{u_0}{v_\infty} \right) dx$$

- Csiszár-Kullback inequality: Consider $v \geq 0$, $\bar{v} \geq 0$ such that $\int_{\mathbf{R}^n} v dx = \int_{\mathbf{R}^n} \bar{v} dx =: M > 0$

$$\int_{\mathbf{R}^n} v \log \left(\frac{v}{\bar{v}} \right) dx \geq \frac{1}{4M} \|v - \bar{v}\|_{L^1(\mathbf{R}^n)}^2$$

$$\implies \|v - v_\infty\|_{L^1(\mathbf{R}^n)}^2 \leq 4M \Sigma[u_0] e^{-2\tau}$$

$$\tau(t) = \log \sqrt{1 + 2t}$$

$$\|u(\cdot, t) - u_\infty(\cdot, t)\|_{L^1(\mathbf{R}^n)}^2 \leq \frac{4}{1 + 2t} \Sigma[u_0]$$

$$u_\infty(x, t) = \frac{1}{R^n(t)} v_\infty \left(\frac{x}{R(t)}, \tau(t) \right)$$

The proof of the Csiszár-Kullback inequality is given by a Taylor development at second order.

Euclidean logarithmic Sobolev inequality: other formulations

1) independent from the dimension: gaussian form [Gross, 75]

$$\int_{\mathbb{R}^n} w \log w \, d\mu(x) \leq \frac{1}{2} \int_{\mathbb{R}^n} w |\nabla \log w|^2 \, d\mu(x)$$

with $w = \frac{v}{M v_\infty}$, $\|v\|_{L^1} = M$, $d\mu(x) = v_\infty(x) dx$.

2) invariant under scaling [Weissler, 78]

$$\int_{\mathbb{R}^n} w^2 \log w^2 \, dx \leq \frac{n}{2} \log \left(\frac{2}{\pi n e} \int_{\mathbb{R}^n} |\nabla w|^2 \, dx \right)$$

for any $w \in H^1(\mathbb{R}^n)$ such that $\int w^2 \, dx = 1$

Proof: take $w = \sqrt{\frac{v}{Mv_\infty}}$ and optimize for $w_\lambda(x) = \lambda^{n/2}w(\lambda x)$

$$\begin{aligned} & \int_{\mathbb{R}^n} |\nabla w_\lambda|^2 dx - \int_{\mathbb{R}^n} w_\lambda^2 \log w_\lambda^2 dx \\ &= \lambda^2 \int_{\mathbb{R}^n} |\nabla w|^2 dx - \int_{\mathbb{R}^n} w^2 \log w^2 dx - n \log \lambda \int_{\mathbb{R}^n} w^2 dx \end{aligned}$$

□

Entropy-entropy production method: a proof of the Euclidean logarithmic Sobolev inequality:

$$\frac{d}{d\tau} (I[v(\cdot, \tau)] - 2\Sigma[v(\cdot, \tau)]) = -C \sum_{i,j=1}^n \int_{\mathbb{R}^n} \left| w_{ij} + a \frac{w_i w_j}{w} + b w \delta_{ij} \right|^2 dx < 0$$

for some $C > 0$, $a, b \in \mathbb{R}$. Here $w = \sqrt{v}$.

$$I[v(\cdot, \tau)] - 2\Sigma[v(\cdot, \tau)] \searrow I[v_\infty] - 2\Sigma[v_\infty] = 0$$

$$\implies \forall u_0, \quad I[u_0] - 2\Sigma[u_0] \geq I[v(\cdot, \tau)] - 2\Sigma[v(\cdot, \tau)] \geq 0 \text{ for } \tau > 0$$

Entropy-entropy production method: improvements of convex Sobolev inequalities

goal: large time behavior of parabolic equations:

$$\begin{cases} v_t = \operatorname{div}_x [D(x) (\nabla_x v + v \nabla_x A(x))] = \operatorname{div}[e^{-A} \nabla(v e^A)] \\ v(x, t=0) = v_0(x) \in L^1_+(\mathbb{R}^n) \end{cases} \quad t > 0, \quad x \in \mathbb{R}^n \quad (3)$$

$A(x)$... given ‘potential’

$v_\infty(x) = e^{-A(x)} \in L^1$... (unique) steady state

mass conservation: $\int_{\mathbb{R}^d} v(t) dx = \int_{\mathbb{R}^d} v_\infty dx = 1$

questions: exponential rate ? connection to logarithmic Sobolev inequalities ? [Bakry-Emery '84, Gross '75, Toscani '96, AMTU...]

[Anton Arnold, J.D.]

ENTROPY-ENTROPY PRODUCTION METHOD

[Bakry, Emery, 84]

[Arnold, Markowich, Toscani, Unterreiter, 01]

Relative entropy of $v(x)$ w.r.t. $v_\infty(x)$:

$$\Sigma[v|v_\infty] := \int_{\mathbf{R}^d} \psi\left(\frac{v}{v_\infty}\right) v_\infty \, dx \geq 0$$

with

$$\psi(w) \geq 0 \text{ for } w \geq 0, \text{ convex}$$

$$\psi(1) = \psi'(1) = 0$$

$$\text{Admissibility condition: } (\psi''')^2 \leq \frac{1}{2}\psi''\psi^{IV}$$

Examples:

$$\psi_1 = w \ln w - w + 1, \quad \Sigma_1(v|v_\infty) = \int v \ln \left(\frac{v}{v_\infty} \right) \, dx \dots \text{ physical entropy}$$

$$\psi_p = w^p - p(w-1) - 1, \quad 1 < p \leq 2, \quad \Sigma_2(v|v_\infty) = \int_{\mathbf{R}^d} (v - v_\infty)^2 v_\infty^{-1} \, dx$$

EXPONENTIAL DECAY OF ENTROPY PRODUCTION

$$I(v(t)|v_\infty) := \frac{d}{dt} \Sigma[v(t)|v_\infty] = - \int \psi''\left(\frac{v}{v_\infty}\right) |\underbrace{\nabla\left(\frac{v}{v_\infty}\right)}_{=:u}|^2 v_\infty dx \leq 0$$

Assume: $D \equiv 1$, $\underbrace{\frac{\partial^2 A}{\partial x^2}}_{\text{Hessian}} \geq \lambda_1 Id$, $\lambda_1 > 0$ ($A(x)$... unif. convex)

entropy production rate:

$$\begin{aligned} I' &= 2 \int \psi''\left(\frac{v}{v_\infty}\right) u^T \cdot \frac{\partial^2 A}{\partial x^2} \cdot u v_\infty dx + \underbrace{2 \int \text{Tr}(XY) v_\infty dx}_{\geq 0} \\ &\geq -2\lambda_1 I \end{aligned}$$

with

$$X = \begin{pmatrix} \psi''\left(\frac{v}{v_\infty}\right) & \psi'''\left(\frac{v}{v_\infty}\right) \\ \psi'''\left(\frac{v}{v_\infty}\right) & \frac{1}{2}\psi IV\left(\frac{v}{v_\infty}\right) \end{pmatrix} \geq 0$$

$$Y = \begin{pmatrix} \sum_{ij} (\frac{\partial u_i}{\partial x_j})^2 & u^T \cdot \frac{\partial u}{\partial x} \cdot u \\ u^T \cdot \frac{\partial u}{\partial x} \cdot u & |u|^4 \end{pmatrix} \geq 0$$

$$\Rightarrow |I(t)| \leq e^{-2\lambda_1 t} |I(t=0)| \quad t > 0$$

$$\forall v_0 \text{ with } |I(v_0|v_\infty)| < \infty$$

EXPONENTIAL DECAY OF RELATIVE ENTROPY

known: $I' \geq -2\lambda_1 \underbrace{I}_{=\Sigma'} , \quad \int_t^\infty \dots dt$ (4)

$$\Rightarrow \Sigma' = I \leq -2\lambda_1 \Sigma$$

Theorem 1 [Bakry, Emery], [Arnold, Markowich, Toscani, Unterreiter]

$$\frac{\partial^2 A}{\partial x^2} \geq \lambda_1 Id \quad (\text{"Bakry-Emery condition"}), \quad \Sigma[v_0|v_\infty] < \infty$$
$$\Rightarrow \Sigma[v(t)|v_\infty] \leq \Sigma[v_0|v_\infty] e^{-2\lambda_1 t}, \quad t > 0$$

$$\|v(t) - v_\infty\|_{L^1}^2 \leq C \Sigma[v(t)|v_\infty] \dots \text{Csiszár-Kullback}$$

CONVEX SOBOLEV INEQUALITIES

Entropy–entropy production estimate (4) for $A(x) = -\ln v_\infty$ (uniformly convex):

$$\Sigma[v|v_\infty] \leq \frac{1}{2\lambda_1} |I(v|v_\infty)|$$

Example 1: logarithmic entropy $\psi_1(w) = w \ln w - w + 1$

$$\int v \ln \left(\frac{v}{v_\infty} \right) dx \leq \frac{1}{2\lambda_1} \int v \left| \nabla \ln \left(\frac{v}{v_\infty} \right) \right|^2 dx$$

$$\forall v, v_\infty \in L^1_+(\mathbb{R}^n), \int v dx = \int v_\infty dx = 1$$

logarithmic Sobolev inequality – “entropy version”

Set $f^2 = \frac{v}{v_\infty} \Rightarrow$

$$\int f^2 \ln f^2 dv_\infty \leq \frac{2}{\lambda_1} \int |\nabla f|^2 dv_\infty$$

$$\forall f \in L^2(\mathbb{R}^n, dv_\infty), \int f^2 dv_\infty = 1$$

logarithmic Sobolev inequality– dv_∞ measure version [Gross '75]

Example 2: non-logarithmic entropies:

$$\psi_p(w) = w^p - p(w - 1) - 1, \quad 1 < p \leq 2$$

$$(B_p) \quad \frac{p}{p-1} \left[\int f^2 dv_\infty - \left(\int |f|^{\frac{2}{p}} dv_\infty \right)^p \right] \leq \frac{2}{\lambda_1} \int |\nabla f|^2 dv_\infty$$

from (4) with $\frac{v}{v_\infty} = \frac{|f|^{\frac{2}{p}}}{\int |f|^{\frac{2}{p}} dv_\infty}$

$$\forall f \in L^{\frac{2}{p}}(\mathbb{R}^n, v_\infty dx)$$

Poincaré-type inequality [Beckner '89], $(B_p) \Rightarrow (B_2)$, $1 < p \leq 2$

REFINED CONVEX SOBOLEV INEQUALITIES

Estimate of entropy production rate / entropy production:

$$\begin{aligned} I' &= 2 \int \psi'' \left(\frac{v}{v_\infty} \right) u^T \cdot \frac{\partial^2 A}{\partial x^2} \cdot uv_\infty dx + \underbrace{2 \int \text{Tr}(XY)v_\infty dx}_{\geq 0} \\ &\geq -2\lambda_1 I \end{aligned}$$

[Arnold, J.D.]: Observation for $\psi_p(w) = w^p - p(w-1) - 1$, $1 < p < 2$:

$$X = \begin{pmatrix} \psi'' \left(\frac{v}{v_\infty} \right) & \psi''' \left(\frac{v}{v_\infty} \right) \\ \psi''' \left(\frac{v}{v_\infty} \right) & \frac{1}{2} \psi^{IV} \left(\frac{v}{v_\infty} \right) \end{pmatrix} > 0$$

- Assume $\frac{\partial A^2}{\partial x^2} \geq \lambda_1 Id \Rightarrow \Sigma'' \geq -2\lambda_1 \Sigma' + \kappa \frac{|\Sigma'|^2}{1+\Sigma}$, $\kappa = \frac{2-p}{p} < 1$
- $$\Rightarrow k(\Sigma[v|v_\infty]) \leq \frac{1}{2\lambda_1} |\Sigma'| = \frac{1}{2\lambda_1} \int \psi''\left(\frac{v}{v_\infty}\right) |\nabla \frac{v}{v_\infty}|^2 dv_\infty$$
- “refined convex Sobolev inequality” with $x \leq k(x) = \frac{1+x-(1+x)^\kappa}{1-\kappa}$
- Set $v/v_\infty = |f|^{\frac{2}{p}} / \int |f|^{\frac{2}{p}} dv_\infty \Rightarrow$

Theorem 2

$$\begin{aligned} \frac{1}{2} \left(\frac{p}{p-1} \right)^2 & \left[\int f^2 dv_\infty - \left(\int |f|^{\frac{2}{p}} dv_\infty \right)^{2(p-1)} \left(\int f^2 dv_\infty \right)^{\frac{2-p}{p}} \right] \\ & \leq \frac{2}{\lambda_1} \int |\nabla f|^2 dv_\infty \quad \forall f \in L^{\frac{2}{p}}(\mathbb{R}^n, dv_\infty) \end{aligned}$$

“refined Beckner inequality” [Arnold, J.D. '00]
 $(rB_p) \Rightarrow (rB_2) = (B_2)$, $1 < p \leq 2$

Long Time Behavior of the Quantum Fokker-Planck equation

[A. Arnold, J. L. Lopez, P. A. Markowich, J. Soler], [C. Sparber, J. A. Carrillo, J. D., P. A. Markowich]

Heat equation with a source term

[G. Karch, J.D.]

An example of application: the flashing ratchet. Long time behavior and dynamical systems interpretation

[M. Chipot, D. Heath, D. Kinderlehrer, M. Kowalczyk, N. Walkington,...]
[J.D., David Kinderlehrer, Michał Kowalczyk]

Flashing ratchet: a simple model for a molecular motor (Brownian motors, molecular ratchets, or Brownian ratchets)

Diffusion tends to spread and dissipate density / transport concentrates density at specific sites determined by the energy landscape: unidirectional transport of mass.

Fokker-Planck type problem

$$\begin{aligned} u_t &= (u_x + \psi_x u)_x & (x, t) \in \Omega \times (0, \infty) \\ u_x + \psi_x u &= 0 & (x, t) \in \partial\Omega \times (0, \infty) \\ u(x, 0) &= u_0(x) & x \in \Omega \end{aligned} \tag{5}$$

$$u_0 > 0, \int_{\Omega} u_0 = 1, \psi = \psi(x, t)$$

PERIODIC STATE AND ASYMPTOTIC BEHAVIOUR

Theorem 1 Let $\psi \in L^\infty([0, T) \times \Omega)$ be a T -periodic potential and assume that there exists a finite partition of $[0, T)$ into intervals $[T_i, T_{i+1})$, $i = 0, \dots, n$ with $T_0 = 0$, $T_n = T$ such that $\psi_{[T_i, T_{i+1}]} \in L^\infty([T_i, T_{i+1}), W^{1,\infty}(\Omega))$. Then there exists a unique nonnegative T -periodic solution U to (5) such that $\int_\Omega U(x, t) dx = 1$ for any $t \in [0, T)$.

Entropy and entropy production : $\sigma_q(u) = \begin{cases} \frac{u^q - 1}{q-1} & \text{if } q > 1, \\ u \ln u & \text{if } q = 1. \end{cases}$

$$\Sigma_q[u|v] = \int_\Omega \left[\sigma_q\left(\frac{u}{v}\right) - \sigma'_q(1) \left(\frac{u}{v} - 1\right) \right] v \, dx$$

$$I_q[u|v] = \int_\Omega \sigma''_q\left(\frac{u}{v}\right) \left|\nabla\left(\frac{u}{v}\right)\right|^2 v \, dx,$$

Theorem 2 Let u_1, u_2 be any two solutions to (5).

$$\Sigma_q[u_1(t)|u_2(t)] \leq e^{-C_q t} \Sigma_q[u_1(0)|u_2(0)]$$

Proposition 3 Ω is a bounded domain in \mathbb{R}^d with C^1 boundary. Let u and v be two nonnegative functions in $L^1 \cap L^q(\Omega)$ if $q \in (1, 2]$ and in $L^1(\Omega)$ with $u \log u$ and $u \log v$ in $L^1(\Omega)$ ($q = 1$).

$$\Sigma_q[u|v] \geq 2^{-2/q} q \left[\max \left(\|u\|_{L^q(\Omega)}^{2-q}, \|v\|_{L^q(\Omega)}^{2-q} \right) \right]^{-1} \|u - v\|_{L^q(\Omega)}^2$$

Corollary 4 Let $q \in [1, 2]$. Any solution of (5) with initial data $u_0 \in L^1 \cap L^q(0, 1)$ $u_0 \log u_0 \in L^1(0, 1)$ if $q = 1$, converges to $\|u_0\|_{L^1} U(x, t)$, (periodic solution):

$$\|u(x, t) - \|u_0\|_{L^1} U(x, t)\|_{L^q(0, 1; dx)} \leq K e^{-C_{q,\psi} t} \quad \forall t \geq 0^{21}$$

Let $\textcolor{violet}{u}_\psi := \|u_0\|_{L^1} \frac{e^{-\psi}}{\int_\Omega e^{-\psi} dx}$.

$$\begin{aligned}\frac{d}{dt} \Sigma_1[u|u_\psi] &= \int_\Omega \left[1 + \log \left(\frac{u}{u_\psi} \right) \right] u_t \, dx - \int_\Omega \frac{u}{u_\psi} u_{\psi,t} \, dx \\ &= -I_1[u|u_\psi] - \int_\Omega \frac{u}{u_\psi} u_{\psi,t} \, dx\end{aligned}$$

Lemma 5 Let $u \geq 0$ be a solution to (5) such that $\int_\Omega u \, dx = 1$. With the above notations, the following estimate holds:

$$\frac{d}{dt} \Sigma_1[u|u_\psi] \leq -C_\psi \Sigma_1[u|u_\psi] + K_\psi.$$

Fixed-point for the map $\mathcal{T}(u(\cdot, 0)) = u(\cdot, T)$ in

$$\mathcal{Y} = \{u \in H^1(\Omega) \mid u \geq 0, \|u\|_{L^1(\Omega)} = 1, \Sigma_1[u|u_0(\cdot, 0)] \leq K_\psi/C_\psi\}.$$

Flashing potentials: same on each time interval.

Let \mathcal{X} be the set of bounded nonnegative functions u in $L^1 \cap L^q(\Omega)$ (resp. in $L^1(\Omega)$ with $u \log u$ in $L^1(\Omega)$) if $q \in (1, 2]$ (resp. if $q = 1$) such that $\int_{\Omega} u dx = 1$.

Theorem 6 Assume that $v \in \mathcal{X}$ with $0 < m := \inf_{\Omega} v \leq v \leq \sup_{\Omega} v =: M < \infty$. For any $q \in [1, 2]$

$$\mathcal{J} = \frac{q}{q-1} \inf_{\substack{u \in \mathcal{X} \\ u \neq v \text{ a.e.}}} \frac{I_q[u|v]}{\Sigma_q[u|v]} \quad \text{if } q > 1, \quad \mathcal{J} = \inf_{\substack{u \in \mathcal{X} \\ u \neq v \text{ a.e.}}} \frac{I_1[u|v]}{\Sigma_1[u|v]} \quad \text{if } q = 1 \tag{6}$$

can be estimated by

$$\mathcal{J} \geq 4 \lambda_1(\Omega) \frac{m}{M} \tag{7}$$

where $\lambda_1(\Omega)$ is Poincaré's constant of Ω (with weight 1).

Relation between entropy and entropy production: exponential decay of the relative entropy.

II — Entropy methods for nonlinear models involving diffusions

A model for traffic flow

[R. Illner, A. Klar, T. Materne, 2002], [Reinhard Illner, J.D., 2002] $f = f(t, v)$ is an homogeneous distribution function, with velocities ranging in $(0, 1)$:

$$f_t = (-B(t, v) f + D(t, v) f')', \quad (t, v) \in \mathbb{R}^+ \times (0, 1)$$

where $f_t = \partial f / \partial t$, $f' = \partial f / \partial v$. Let $C(t, v) := -\int_0^v \frac{B(t, w)}{D(t, w)} dw$

$$g(t, v) = \rho \frac{e^{-C(t, v)}}{\int_0^1 e^{-C(t, w)} dw} \quad \text{is a local equilibrium}$$

Zero flux: $-B(t, v) g + D(t, v) g' = 0$ but $g_t \equiv 0$ is not granted.

Relative entropy:

$$\begin{aligned} e[t, f] := & \int_0^1 (f \log f - g \log g + C(t, v)(f - g)) dv \\ & - \iint_{(0,1) \times (0,t)} C_t(s, v)(f - g)(s, v) dv ds \end{aligned}$$

Then $\frac{d}{dt} e[t, f(t, .)] = -\int_0^1 D(t, v) f \left| \frac{f'}{f} - \frac{g'}{g} \right|^2 dv \dots$ but we dont have a lower bound for $e[t, f(t, .)]$.

Density: $\rho = \int_0^1 f(t, v) dv$ does not depend on t

Mean velocity: $u(t) = \frac{1}{\rho} \int_0^1 v f(t, v) dv$

Braking term:

$$B(t, v) = \begin{cases} -C_B |v - u(t)|^2 \rho \left(1 - \left|\frac{v-u(t)}{1-u(t)}\right|^\delta\right) & \text{if } v > u(t) \\ C_A |v - u(t)|^2 (1 - \rho) & \text{if } v \leq u(t) \end{cases}$$

Diffusion term: $D(t, v) = \sigma m_1(\rho) m_2(u(t)) |v - u(t)|^\gamma$

Proposition 7 [Illner-Klar-Materne02] *Any stationary solution is uniquely determined by ρ and its average velocity u . The set $(\rho, u[\rho])$ is in general multivalued. For any $\rho \in (0, 1]$.*

Example. *The Maxwellian case.*

CONVEX ENTROPIES

Relative entropy of f w.r.t. g by $E[f|g] = \int_0^1 \Phi\left(\frac{f}{g}\right) g \, dv$

“Standard” example: $\Phi_\alpha(x) = (x^\alpha - x)/(\alpha - 1)$ for some $\alpha > 1$,
 $\Phi(x) = x \log x$ if “ $\alpha = 1$ ”

$$\begin{cases} f_t = \left[D(t, v) f \left(\frac{f'}{f} - \frac{g'}{g} \right) \right]' = \left[D(t, v) g \left(\frac{f}{g} \right)' \right]' & \forall (t, v) \in \mathbb{R}^+ \times (0, 1) \\ \left(\frac{f}{g} \right)' (t, v) = 0 & \forall t \in \mathbb{R}^+, v = 0, 1 \end{cases}$$

$$g(t, v) := \kappa(t) e^{-C(t, v)} \text{ for some } \kappa(t) \neq 0.$$

$$\begin{aligned} \frac{d}{dt} E[f(t, \cdot) | g(t, \cdot)] &= \int_0^1 \Phi'\left(\frac{f}{g}\right) f_t \, dv + \int_0^1 \underbrace{\left[\Phi\left(\frac{f}{g}\right) - \frac{f}{g} \Phi'\left(\frac{f}{g}\right) \right]}_{\Psi\left(\frac{f}{g}\right)} g_t \, dv \\ &= 0 \quad \text{if} \quad \dot{\kappa} = \kappa \frac{\int_0^1 \Psi\left(\frac{f}{g}\right) g C_t(t, v) \, dv}{\int_0^1 \Psi\left(\frac{f}{g}\right) g \, dv}, \quad \kappa(0) = 1 \end{aligned}$$

with $\Psi(x) := \Phi(x) - x\Phi'(x) < 0$

CONVERGENCE TO A STATIONARY SOLUTION

$$\limsup_{t \rightarrow +\infty} \kappa(t) < +\infty .$$

Theorem 8 Let $\Phi = \Phi_\alpha(x) = (x^\alpha - x)/(\alpha - 1)$, f be a smooth global in time solution and assume that $E[f|g]$ is well defined and C^1 in t . If $\exists \varepsilon \in (0, \frac{1}{2})$ s.t. $\varepsilon < u(t) = \frac{1}{\rho} \int_0^1 v f(t, v) dv < 1 - \varepsilon$ $\forall t > 0$, then, as $t \rightarrow +\infty$, $f(t, \cdot)$ converges a.e. to a stationary solution f_∞ .

Coupling with a Poisson equation

[AMT], [P. Biler, J.D.] Nernst-Planck / Debye-Hückel drift-diffusion

$$\begin{aligned} u_t &= \nabla \cdot (\nabla u + u \nabla \phi) \\ v_t &= \nabla \cdot (\nabla v - v \nabla \phi) \\ \Delta \phi &= v - u \end{aligned}$$

in a bounded domain $\Omega \subset \mathbb{R}^d$, $d \leq 3$.

No-flux boundary conditions: $\frac{\partial u}{\partial \nu} + u \frac{\partial \phi}{\partial \nu} = 0$, $\frac{\partial v}{\partial \nu} - v \frac{\partial \phi}{\partial \nu} = 0$ on $\partial \Omega$
and either **conducting boundary conditions**

$$\phi = 0 \quad \text{on } \partial \Omega$$

or “**free**” boundary conditions (this corresponds to a container immersed in a medium with the same dielectric constant as the solute)

$$\phi = E_d * (v - u)$$

where E_d is the fundamental solution of the Laplacian in \mathbb{R}^d .

INTERMEDIATE ASYMPTOTICS: $\Omega = \mathbb{R}^d$

$$\frac{1}{M_u} u_{as}(x, t) = \frac{1}{M_v} v_{as}(x, t) = (2\pi(2t+1))^{d/2} \exp\left(-\frac{|x|^2}{2(2t+1)}\right)$$

Entropy functional :

$$\Sigma = \int u \log\left(\frac{u}{u_{as}}\right) dx + \int v \log\left(\frac{v}{v_{as}}\right) dx + \frac{1}{2} \|\nabla \phi(t)\|_2^2$$

Theorem 9 *There exists a constant C such that*

$$L(t), \quad \|u(t) - u_{as}(t)\|_1^2 + \|v(t) - v_{as}(t)\|_1^2 + \|\nabla \phi(t)\|_2^2 \leq C H(t)$$

$$\text{where } H(t) = \begin{cases} (2t+1)^{-1/2}, & d = 3 \\ (2t+1)^{-1}(\log(2t+1) + 1), & d = 4 \\ (2t+1)^{-1}, & d > 4 \end{cases}$$

Moreover if $M_u = M_v$, then $H(t) = (2t+1)^{-1}$ for any $d \geq 3$.

CONVERGENCE TO A STATIONARY SOLUTION Assume that Ω is bounded. Entropy functional

$$\begin{aligned} W(t) &= \int u \log u \, dx - \int U(x) \log U(x) \, dx \\ &\quad + \int v \log v \, dx - \int V(x) \log V(x) \, dx \\ &\quad + \frac{1}{2} \int (u - v)\phi \, dx - \frac{1}{2} \int (U - V)\Phi \, dx, \end{aligned}$$

where U, V, Φ is the unique stationary solution such that $\int U \, dx = \int u \, dx, \int V \, dx = \int v \, dx$.

Theorem 10 If $d \geq 2$, then there exist two positive constants λ and C such that

$$W(t) \leq W(0) e^{-\lambda t}$$

$$\|u(t) - U\|_1^2 + \|v(t) - V\|_1^2 + \|\nabla(\phi - \Phi)\|_2^2 \leq C e^{-\lambda t}$$

Streater's model

[P. Biler, J.D., M. Esteban, G. Karch]

Nonlinear diffusions coupled with a Poisson equation

[P. Biler, J.D., P. Markowich]

Entropy methods for nonlinear diffusions

- [1. Comparison techniques]
- 2. Entropy-Entropy production methods
- 3. Variational approach: [Gross], [Aubin], [Talenti], [Weissler], [Carlen],
[Carlen-Loss], [Beckner], [Del Pino, JD]
- [4. Mass transportation methods] [Cordero-Erausquin, Gangbo, Houdré],
[Cordero-Erausquin, Nazaret, Villani], [Agueh, Ghoussoub, Kang]

OPTIMAL CONSTANTS FOR GAGLIARDO-NIRENBERG INEQ.

[Del Pino, J.D.]

Theorem 11 $1 < p < n$, $1 < a \leq \frac{p(n-1)}{n-p}$, $b = p \frac{a-1}{p-1}$

$$\begin{aligned} \|w\|_b &\leq S \|\nabla w\|_p^\theta \|w\|_a^{1-\theta} && \text{if } a > p \\ \|w\|_a &\leq S \|\nabla w\|_p^\theta \|w\|_b^{1-\theta} && \text{if } a < p \end{aligned}$$

Equality if $w(x) = A(1 + B|x|^{\frac{p}{p-1}})_+^{-\frac{p-1}{a-p}}$

$$a > p: \theta = \frac{(q-p)n}{(q-1)(np-(n-p)q)}$$

$$a < p: \theta = \frac{(p-q)n}{q(n(p-q)+p(q-1))}$$

Proof based on [Serrin, Tang]

THE OPTIMAL L^p -EUCLIDEAN LOGARITHMIC SOBOLEV INEQUALITY

[Del Pino, J.D., 2001], [Gentil 2002], [Cordero-Erausquin, Gangbo, Houdré, 2002]

Theorem 12 *If $\|u\|_{L^p} = 1$, then*

$$\int |u|^p \log |u| dx \leq \frac{n}{p^2} \log \left[\mathcal{L}_p \int |\nabla u|^p dx \right]$$
$$\mathcal{L}_p = \frac{p}{n} \left(\frac{p-1}{e} \right)^{p-1} \pi^{-\frac{p}{2}} \left[\frac{\Gamma(\frac{n}{2}+1)}{\Gamma(n \frac{p-1}{p} + 1)} \right]^{\frac{p}{n}}$$

$$Equality: u(x) = \left(\pi^{\frac{n}{2}} \left(\frac{\sigma}{p} \right)^{\frac{n}{p^*}} \frac{\Gamma(\frac{n}{p^*}+1)}{\Gamma(\frac{n}{2}+1)} \right)^{-1/p} e^{-\frac{1}{\sigma}|x-\bar{x}|^{p^*}}$$

$p = 2$: Gross' logarithmic Sobolev inequality [Gross, 75], [Weissler, 78]

$p = 1$: [Ledoux 96], [Beckner, 99]

Remark: The three following identities are equivalent:

(i) For any $w \in W^{1,p}(\mathbb{R}^n)$ with $\int |w|^p dx = 1$,

$$\int |w|^p \log |w| dx \leq \frac{n}{p^2} \log \left[\mathcal{L}_p \int |\nabla w|^p dx \right]$$

(ii) Let P_t^p be the semigroup associated $\textcolor{red}{u_t} = \Delta_p(u^{1/(p-1)})$:

$$\|P_t^p f\|_\beta \leq \|f\|_\alpha A(n, p, \alpha, \beta) t^{-\frac{n}{p} \frac{\beta-\alpha}{\alpha\beta}}$$

(iii) Let Q_t^p be the semigroup associated to $\textcolor{red}{v_t} + \frac{1}{p} |\nabla v|^p = 0$:

$$\|e^{Q_t^p g}\|_\beta \leq \|e^g\|_\alpha B(n, p, \alpha, \beta) t^{-\frac{n}{p} \frac{\beta-\alpha}{\alpha\beta}}$$

Case $a = p$ ($q = 1$): it is more convenient to use this inequality in a non homogeneous form:

$$\inf_{\mu > 0} \left[\frac{n}{p} \log \left(\frac{n}{p\mu} \right) + \mu \frac{\|\nabla w\|_p^p}{\|w\|_p^p} \right] = n \log \left(\frac{\|\nabla w\|_p}{\|w\|_p} \right) + \frac{n}{p}.$$

Corollary 13 For any $w \in W^{1,p}(\mathbb{R}^n)$, $w \neq 0$, for any $\mu > 0$,

$$p \int |w|^p \log \left(\frac{|w|}{\|w\|_p} \right) dx + \frac{n}{p} \log \left(\frac{p\mu e}{n \mathcal{L}_p} \right) \int |w|^p dx \leq \mu \int |\nabla w|^p dx.$$

Case $a \neq p$ ($q \neq 1$): ... $\Sigma \leq C I$

INTERMEDIATE ASYMPTOTICS FOR NONLINEAR DIFFUSIONS

[Manuel Del Pino, J.D.]

$$u_t = \Delta_p u^m$$

Convergence to a stationary solution for:

$$v_t = \Delta_p v^m + \nabla(x v)$$

Let $q = 1 + m - (p - 1)^{-1}$. Whether q is bigger or smaller than 1 determines two different regimes like for $m = 1$.

For $q > 0$, define the *entropy* by

$$\Sigma[v] = \int [\sigma(v) - \sigma(v_\infty) - \sigma'(v_\infty)(v - v_\infty)] dx$$

$$\sigma(s) = \frac{s^q - s}{q-1} \text{ if } q \neq 1$$

$$\sigma(s) = s \log s \text{ if } q = 1 \ (p = 2)$$

[Del Pino, J.D.] Intermediate asymptotics of $u_t = \Delta_p u^m$

Theorem 14 $n \geq 2$, $1 < p < n$, $\frac{n-(p-1)}{n(p-1)} \leq m \leq \frac{p}{p-1}$ and $q = 1 + m - \frac{1}{p-1}$

$$(i) \quad \|u(t, \cdot) - u_\infty(t, \cdot)\|_q \leq K R^{-(\frac{\alpha}{2} + n(1 - \frac{1}{q}))}$$

$$(ii) \quad \|u^q(t, \cdot) - u_\infty^q(t, \cdot)\|_{1/q} \leq K R^{-\frac{\alpha}{2}}$$

$$(i): \frac{1}{p-1} \leq m \leq \frac{p}{p-1}$$

$$(ii): \frac{n-(p-1)}{n(p-1)} \leq m \leq \frac{1}{p-1}$$

$$\alpha = (1 - \frac{1}{p}(p-1)^{\frac{p-1}{p}}) \frac{p}{p-1}, \quad R = (1 + \gamma t)^{1/\gamma}, \quad \gamma = (mn+1)(p-1) - (n-1)$$

$$u_\infty(t, x) = \frac{1}{R^n} v_\infty(\log R, \frac{x}{R})$$

$$v_\infty(x) = (C - \frac{p-1}{mp} (q-1) |x|^{\frac{p}{p-1}})_+^{1/(q-1)} \text{ if } m \neq \frac{1}{p-1}$$

$$v_\infty(x) = C e^{-(p-1)^2 |x|^{p/(p-1)}/p} \text{ if } m = (p-1)^{-1}.$$

Use $v_t = \Delta_p v^m + \nabla \cdot (x v)$

$$w = v^{(mp+q-(m+1))/p}, \quad a = b q = p \frac{m(p-1)+p-2}{mp(p-1)-1}.$$

Intermediate asymptotics in L^1 for general nonlinear diffusion equations

[P. Biler, J.D., M. Esteban]

$$u_t = \Delta f(u) \quad (8)$$

If $f(u) = u^m$, Equation (8) can be transformed via a suitable space-time rescaling into an autonomous equation

$$v_\tau = \Delta f(v) + \nabla \cdot (v \nabla V), \quad (9)$$

with $V(x) = \frac{1}{2}|x|^2$. Here we simply assume that f *behaves like a power law* in a neighborhood of 0.

General case: Time-dependent rescaling $u(t, x) = R^{-d}(t) v\left(\tau(t), \frac{x}{R(t)}\right)$,
 $\tau(t) = \log(R(t))$, $\frac{dR}{dt} = R^{(1-m)d-1}$

$$v_\tau = e^{md\tau} \Delta f\left(e^{-d\tau} v\right) + \nabla \cdot (xv) \quad (10)$$

“Local” stationary solution: $v_\infty^R(x) = R^d \mathbf{g} \left(R^{-(m-1)d} (\alpha_\infty^R(M) - \frac{1}{2}|x|^2) \right)$

such that $\|v_\infty^R\|_{L^1(\mathbb{R}^d)} = \|v\|_{L^1(\mathbb{R}^d)}$.

Example: $f(u) = u^m$, $m \neq 1$, $v_\infty(x) = \mathbf{g}(\alpha_\infty(M) - \frac{1}{2}|x|^2)$,
 $\mathbf{g}(x) = ((m-1)x/m)_+^{1/(m-1)}$

Assumption:

- (i) $f(s) = s^m \varphi(s)$ with $\varphi \in C^0([0, +\infty)) \cap C^1(0, +\infty)$, $\varphi > 0$, $\varphi(0) = 1$, $\varphi'(s) = \mathcal{O}(s^k)$ ($k > -1$) near $s = 0^+$
- (ii) there exists a function H with $H''(s) = \frac{f'(s)}{s}$, $H(0) = 0$, such that $(m-1)H(s) \leq f(s) \quad \forall s > 0$
- (iii) $f \in C^3(0, +\infty) \cap C^0[0, +\infty)$, $f' > 0$, $f'/s \in L^1_{loc}(0, +\infty)$ and $f(u) \leq \frac{d}{d-1}uf'(u)$ for any $u > 0$ ($d \geq 3$). If $m \leq 1$, $\exists s_0 > 0$ such that $f''_{|[0,s_0]} \geq 0$ and $\mathbf{g}(\alpha - \frac{1}{2}|\cdot|^2) \in L^1(\mathbb{R}^d)$.

Relative entropy :

$$\Sigma[\tau, v] = \int_{\mathbb{R}^d} \left(e^{md\tau} H(e^{-d\tau} v) - e^{md\tau} H(e^{-d\tau} v_\infty^R) + \frac{1}{2}|x|^2(v - v_\infty^R) \right) dx$$

Theorem 15 Under these assumptions, then for any solution v of (10), there exist two constants $\beta := \min\{2, d(k+1)\} > 0$, $K > 0$ such that for all $\tau > 0$,

$$0 \leq \Sigma[\tau, v] \leq K e^{-\beta \tau}.$$

Rates are as good as in the power-law case as soon as $k \geq -(d-2)/d$ (if the function φ is C^1 at the origin, then $\beta = 2$).

Theorem 16 If moreover $\ell_1 := \liminf_{s \rightarrow +\infty} f'(s)s^{m-1} > 0$ if $m \in (1, 2)$; $\ell_2 := \liminf_{s \rightarrow +\infty} \frac{s H''(s)}{|H'(s)|^3} > 0$ if $m \in \left(\frac{d-1}{d}, 1\right)$, then as $t \rightarrow +\infty$,

$$\|u(t, \cdot) - u_\infty(t, \cdot)\|_{L^1(\mathbf{R}^d, u_\infty^{m-1}(t, x) dx)} \leq C t^{-\frac{d(m-1)+\beta/2}{2+(m-1)d}} \quad \text{if } 1 < m \leq 2$$

$$\|H(u)(t, \cdot) - H(u_\infty)(t, \cdot)\|_{L^1(\mathbf{R}^d, dx)} \leq C t^{-\frac{d(m-1)+\beta/2}{2+(m-1)d}} \quad \text{if } (d-1)/d \leq m$$

III — L^1 Intermediate asymptotics for scalar conservation laws

J.D., Miguel Escobedo

Introduction

Consider for some $q > 1$ a nonnegative entropy solution of

$$\begin{cases} U_\tau + (U^q)_\xi = 0, \quad \xi \in \mathbb{R}, \quad \tau > 0, \\ U(\tau = 0, \cdot) = U_0. \end{cases} \quad (11)$$

T.-P. Liu & M. Pierre (1984): for every $p \in [1, +\infty)$,

$$\lim_{\tau \rightarrow \infty} \tau^{\frac{1}{q}(1 - \frac{1}{p})} \|U(\tau) - U_\infty(\tau)\|_p = 0.$$

U_∞ is the self-similar solution defined by

$$U_\infty(\tau, \xi) = \begin{cases} \left(\frac{|\xi|}{q\tau} \right)^{\frac{1}{q-1}} & 0 \leq \xi \leq c(\tau) := q (\|u_0\|_1/(q-1))^{(q-1)/q} \tau^{1/q} \\ 0 & \text{elsewhere} \end{cases}$$

P. Lax (1957): if U is the unique entropy solution to

$$U_\tau + f(U)_\xi = 0, \quad U(0, \xi) = U_0(\xi)$$

with $f \in C^2$, $f(0) = f'(0) = 0$ and $f'' > 0$, and if $U_0 \geq 0$ has compact support in (s_-, s_+) , then the following estimate holds:

$$\|U(\tau, \cdot) - W_\infty(\tau, \cdot - s_-)\|_1 = O(\tau^{-1/2}) \quad \text{as } \tau \rightarrow \infty,$$

where $W_\infty(\tau, \xi) = \frac{\xi}{f''(0)} \tau^{-1}$ if $0 < \xi < -s_- + s_+ + \sqrt{2 \|u_0\|_1 f''(0)} \tau^{-1/2}$, and 0 elsewhere.

Y.-J. Kim (2003): an alternative approach is based on a detailed study of special self-similar solutions. Same kind of results.

MAIN RESULTS

Theorem 17 [ENTROPY ESTIMATE] *Let U be a global, piecewise C^1 entropy solution of (11) corresponding to a nonnegative initial data U_0 in $L^1 \cap L^\infty(\mathbb{R})$ which is compactly supported in $(\xi_0, +\infty)$ for some $\xi_0 \in \mathbb{R}$ and such that*

$$\liminf_{\substack{\xi \rightarrow \xi_0 \\ \xi > \xi_0}} \frac{U_0(\xi)}{|\xi - \xi_0|^{1/(q-1)}} > 0.$$

Then, for any $\alpha \in (0, \frac{q}{q-1})$ and $\epsilon > 0$,

$$\limsup_{\tau \rightarrow +\infty} \tau^{\alpha-\epsilon} \int_{\mathbb{R}} |U(\tau, \xi) - U_\infty(\tau, \xi - \xi_0)| \frac{d\xi}{|\xi - \xi_0|^\alpha} = 0.$$

$$\frac{1}{2} < \alpha(1 - 1/q) \iff q/(2(q-1)) < \alpha$$

Corollary 18 [L¹ ESTIMATE] *For any $\beta < 1$, there exists a constant C_β such that*

$$\|U(\tau, \cdot) - U_\infty(\tau, \cdot - \xi_0)\|_1 \leq C_\beta \tau^{-\beta}$$

Theorem 19 [UNIFORM ESTIMATE]

$$\lim_{\tau \rightarrow +\infty} \sup_{\xi \in \text{supp}(U(\tau, \cdot))} \tau^{1/q} \left| U(\tau, \xi) - \left| \frac{\xi - \xi_0}{\tau q} \right|^{1/(q-1)} \right| = 0$$

$$\lim_{\tau \rightarrow +\infty} (1 + q \tau)^{-1/q} \max [\text{supp}(U(\tau, \cdot))] = \left(\frac{q \int_{\mathbf{R}} U_0 d\xi}{q - 1} \right)^{(q-1)/q} =: c_M$$

$$\max [\text{supp}(U(\tau, \cdot))] \geq (1 + q \tau)^{1/q} c_M (1 + O(\tau^{-1}))$$

Two main tools:

1. A *time-dependent rescaling* which preserves the initial data and replaces the characterization of the intermediate asymptotics by the convergence to a stationary solution.
2. A time-decreasing functional which plays the role of an *entropy*, in the sense that it captures global informations on the evolution of the solution and controls its large time asymptotics.

NOTIONS OF SOLUTION, TIME-DEPENDENT RESCALING, ENTROPY

Proposition 20 Let $U \geq 0$ be a piecewise C^1 entropy solution of (11), whose points of discontinuity are $\xi_1(\tau) < \xi_2(\tau) < \dots < \xi_n(\tau)$. The rescaled function

$$u(t, x) = e^t U \left(\frac{1}{q} (e^{qt} - 1), e^t x \right)$$

is a piecewise C^1 function, whose points of discontinuity are given by $s_i(t) \equiv e^{-t} \xi_i((e^{qt} - 1)/q)$ such that

$$s'_i(t) = \frac{(u_i^+)^q - s_i(t) u_i^+ - (u_i^-)^q + s_i(t) u_i^-}{u_i^+ - u_i^-}.$$

Out of the curves $x = s_i(t)$ the function u is a classical solution of

$$u_t = (x u - u^q)_x \tag{12}$$

and across these curves it satisfies

$$u_i^- := \lim_{\substack{x \rightarrow s_i(t) \\ x < s_i(t)}} u(t, x) > \lim_{\substack{x \rightarrow s_i(t) \\ x > s_i(t)}} u(t, x) := u_i^+.$$

$U_0 := U(0, \cdot) = u(0, \cdot) =: u_0$. If $U_0 \in L^1(\mathbb{R})$: $\|u(t)\|_1 = \|U_0\|_1$ $\forall t > 0$.

DEFINITION: u is a **solution** of (12) iff it is a piecewise C^1 function whose discontinuity points are given by $\{s_i(t)\}_{i=1}^n$ satisfying (20), which solves (12) out of these curves and such that () holds across them. With this definition, if v is a solution of (12) then U is an entropy solution of (11).

For every $c > 0$, let u_∞^c be the **stationary solution** of (12) defined by

$$u_\infty^c(x) = \begin{cases} x^{1/(q-1)} & 0 \leq x \leq c \\ 0 & \text{if } x < 0 \text{ or } x > c \end{cases} \quad (13)$$

If $c = c_M := (qM/(q-1))^{(q-1)/q}$, $u_\infty := u_\infty^{c_M}$. $\|u_\infty\|_1 = M$.

COMPARISON RESULTS Consider two solutions U, V of (11) with initial data U_0, V_0 such that $0 \leq U_0 \leq A_0 V_0$ a.e. for some $A_0 > 0$. Then

$$U(\tau, \cdot) \leq A_0 V(A_0^{q-1}\tau, \cdot) \text{ a.e. } \forall \tau \in \mathbb{R}^+.$$

Scaling argument and comparison principle for entropy solutions.

$$\begin{aligned} u_0 \leq A_0 u_\infty^c \text{ a.e.} \implies u(t, x) \leq A(t) u_\infty^{c(t)}(x) \text{ a.e. } \forall t \in \mathbb{R}^+ \\ \text{with } A(t) = \frac{A_0 e^{qt/(q-1)}}{\left[1 + A_0^{q-1}(e^{qt} - 1)\right]^{1/(q-1)}} \text{ and } c(t) = c \left(\frac{A_0}{A(t)}\right)^{(q-1)/q}. \end{aligned}$$

A GRAPH CONVERGENCE RESULT

Proposition 21 Let $\epsilon > 0$, $M = \int u_0 dx$ and $c_M = (\frac{qM}{q-1})^{(q-1)/q}$. Assume that $u_0 \geq 0$ is a piecewise C^1 initial data with support in $[0, +\infty)$, such that $\liminf_{\substack{x \rightarrow 0 \\ x > 0}} x^{1/(1-q)} u_0(x) > 0$.

Then there exists $T > 0$ such that, for any $t > T$:

- (i) the support of $u(\cdot, t)$ is an interval $[0, s(t)]$.
- (ii) $\inf_{x \in [0, s(t)]} x^{1/(1-q)} u(x, t) > 0$.
- (iii) there exists a constant $A_0 > 0$ such that $u \leq A(t) u_\infty^{s(t)}$
- (iv) for an arbitrarily small $\epsilon > 0$, there exists a constant κ such that $u \geq (1 - \kappa e^{-qt}) u_\infty^{c_M - \epsilon}$.

FIGURES

Theorem 22 Let $u_0 \geq 0$ be a piecewise C^1 initial data with compact support in $[0, +\infty)$ such that $\liminf_{\substack{x \rightarrow 0 \\ x > 0}} x^{1/(1-q)} u_0(x) > 0$.

Then

$$\lim_{t \rightarrow \infty} \sup_{x \in (0, s(t))} |u(t, x) - u_\infty^{s(t)}| = 0,$$

where $[0, s(t)] := \text{Supp}(u(\cdot, t))$ for $t > 0$ large enough. Moreover

$$\lim_{t \rightarrow +\infty} s(t) = c_M =: \left(\frac{qM}{q-1}\right)^{(q-1)/q} \quad \text{with } M = \int_{\mathbf{R}} u_0 \, dx$$

and $s(t) \geq c_M - O(e^{-qt})$.

1st step: A CONSEQUENCE OF THE RANKINE-HUGONIOT RELATION

Let $h(t) := \lim_{x \rightarrow s(t), x < s(t)} u(x, t)$. Then

$$\begin{cases} \frac{ds}{dt} = h^{q-1} - s \\ \frac{dh}{dt} = h \left(1 - \lim_{\substack{x \rightarrow s(t) \\ x < s(t)}} (u^{q-1})_x(x, t) \right) \end{cases} \quad (14)$$

2nd step: ENTROPY SOLUTION If U is an entropy solution of (11) then it satisfies the *entropy inequality* $(U^{q-1})_\xi \leq \frac{1}{q\tau}$ and

$$(u^{q-1})_x \leq (1 - e^{-qt})^{-1}$$

holds in the distribution sense.

3rd step: BY CONTRADICTION Let $u_\infty(x) := x^{1/(q-1)}$. Then

(i) For any $\epsilon > 0$, there exists $t_1 > 0$ such that

$$s(t) \geq C_M - \epsilon \quad \forall t > t_1.$$

(ii) For any $\epsilon > 0$, $\delta \in (0, 1)$, $t_0 > 0$, there exists $t_1 > t_0$ such that

$$h(t_1) \geq (1 - \delta) u_\infty(s(t_1)).$$

4th step: COMBINE STEP 1 AND 2

$$\frac{dh}{dt} \geq -\frac{h}{e^{qt} - 1}.$$

Assume that $h(t_0) = h_0 > 0$ for some $t_0 > 0$. Then

$$h(t) \geq h_0 (1 - e^{-qt_0})^{1/q} \quad \forall t > t_0.$$

5th step: AN ESTIMATE BASED ON THE MASS As $t \rightarrow +\infty$, $s(t)$ converges to C_M and for any $\eta \in (0, 1)$, there exists a $t_0 \geq 0$ such that

$$u(\cdot, t) \geq (1 - \eta) u_\infty^{s(t)} \quad \forall t \geq t_0 .$$

Proof. This follows from

$$u(x, t) \geq \left((h(t))^{q-1} - \frac{s(t) - x}{1 - e^{-qt}} \right)_+^{1/(q-1)} =: v(x, t) \quad \text{on } (C_M - \epsilon, s(t)) . \quad (15)$$

and for $s(t) > c_M$ from

$$M \geq (1 - \eta) \int_0^{C_M - \epsilon} u_\infty \, dx + \int_{C_M - \epsilon}^{s(t)} v \, dx .$$

□

L^1 INTERMEDIATE ASYMPTOTICS

For any positive constants c and c' , define the **entropy** by

$$\Sigma(t) = \int_0^{c'} \mu(x) |u(t, x) - u_\infty^c(x)| dx = \int_0^c \mu |u - u_\infty^c| dx + \int_c^{c'} \mu u dx \quad (16)$$

for some nonincreasing function μ , which is continuous on $(0, +\infty)$. The special choice $c' > \sup_{t \in \mathbb{R}^+} \max_{x \in \mathbb{R}} \{\text{supp } u(t, x)\}$, $\mu(x) := |x|^{-\alpha}$ and $c = c_M$ will provide exponential rates of convergence. Let $f(v) := v - v^q$.

Proposition 23 [DECAY OF THE ENTROPY] *Let u_0 be an initial data with compact support in $[0, +\infty)$, such that $0 \leq u_0(x) \leq A_0 x^{1/(q-1)}$ for some $A_0 > 0$. Assume that $\lim_{x \rightarrow 0} \mu(x) u_\infty^q(x) = 0$. Let $c' > 0$. Suppose that $\mu' u_\infty^q$ and μu_∞ are integrable on $(0, c')$.*

Then for every fixed $c \in (0, c')$, for $t \geq 0$ a.e.

$$\begin{aligned} \frac{d\Sigma}{dt} &\leq \int_0^c \mu'(u_\infty^c)^q \left| f\left(\frac{u}{u_\infty^c}\right) \right| dx - \int_c^{c'} \mu'(u_\infty^{c'})^q f\left(\frac{u}{u_\infty^{c'}}\right) dx \\ &\quad - \mu(c) c^{q/(q-1)} \left\{ f\left(\frac{u^+(c)}{c^{\frac{1}{q-1}}}\right) + \left| f\left(\frac{u^-(c)}{c^{\frac{1}{q-1}}}\right) \right| \right\} + \mu(c') (c')^{q/(q-1)} f\left(\frac{u^-(c')}{(c')^{\frac{1}{q-1}}}\right) \end{aligned}$$

where $u^\pm(c) := \lim_{\substack{x \rightarrow c \\ \pm(x-c) > 0}} u(x)$.

If $c = c'$, then $\frac{d\Sigma}{dt} \leq \int_0^c \mu'(u_\infty^c)^q \left| f\left(\frac{u}{u_\infty^c}\right) \right| dx \leq 0$.

Proof.

$$\frac{d\Sigma}{dt} = \int_0^c \mu u_t \left[\mathbb{1}_{u>u_\infty^c} - \mathbb{1}_{u<u_\infty^c} \right] dx + \int_c^{c'} \mu u_t dx + [\mu(s)|u - u_\infty^c(s)| \cdot s'(t)]_{u^+}^{u^-}$$

Using the equation and an integration by parts, we get

$$\begin{aligned} \frac{d\Sigma}{dt} &\leq \int_0^c \mu' (u_\infty^c)^q \left| f\left(\frac{u}{u_\infty^c}\right) \right| dx + \mu(s) (u_\infty^c(s))^q \Psi(v^-, v^+) \\ &\quad - \mu(c) c^{\frac{q}{q-1}} \left[\left| f\left(c^{-\frac{1}{q-1}} u^-(t, c)\right) \right| + f\left(c^{-\frac{1}{q-1}} u^+(t, c)\right) \right] \\ &\quad - \int_c^{c'} \mu' (u_\infty^{c'})^q f\left(\frac{u}{u_\infty^{c'}}\right) dx + \mu(c') (c')^{q/(q-1)} f\left(\frac{u^-(c')}{(c')^{\frac{1}{q-1}}}\right) \end{aligned}$$

where

$$\Psi(v^-, v^+) = [f(v^+) - f(v^-)] \cdot \frac{|v^+ - 1| - |v^- - 1|}{v^+ - v^-} + |f(v^+)| - |f(v^-)|$$

Since $v^+ < v^-$, we have to distinguish three cases:

(i) $1 \leq v^+ \leq v^-$: $f(v^-) \leq f(v^+) \leq 0$ and $\Psi(v^-, v^+) = 0$.

(ii) $v^+ < 1 \leq v^-$: $f(v^-) \leq 0 < f(v^+)$ and by concavity of $f(v) = v - v^q$

$$\begin{aligned}\frac{1}{2} \Psi(v^-, v^+) &= \frac{v^- - 1}{v^- - v^+} f(v^+) + \frac{1 - v^+}{v^- - v^+} f(v^-) \\ &\leq f\left(\frac{v^- - 1}{v^- - v^+} v^+ + \frac{1 - v^+}{v^- - v^+} v^-\right) = f(1) = 0.\end{aligned}$$

(iii) $v^+ < v^- \leq 1$: $f(v^-) \geq 0$ and $f(v^+) \geq 0$, $\Psi(v^-, v^+) = 0$.

Rates of decay

Proposition 24 Assume that $c \leq c_M$, $c \leq c'$ and $c = c_M$ if $c' > c$. If $\mu(x) = x^{-\alpha} \quad \forall x > 0$ for some $\alpha \in (0, \frac{q}{q-1})$, then $\lim_{t \rightarrow +\infty} \Sigma_\alpha(t) = 0$ and

$$\frac{d\Sigma_\alpha}{dt} + (q-1)\alpha \Sigma_\alpha(t) - \alpha \int_c^{c'} x^{-\alpha} u \, dx - r(c') = o(\Sigma_\alpha(t)) \quad \text{as } t \rightarrow +\infty$$

with $r(c') = \mu(c') (c')^{q/(q-1)} f((c')^{-1/(q-1)} u^-(c'))$.

Proof. With $c(t) := \min(s(t), c)$,

$$\frac{d\Sigma_\alpha}{dt} + (q-1)\alpha \Sigma_\alpha(t) - r(c') \leq C_q \int_0^{c(t)} x^{-\alpha + \frac{1}{q-1}} \left(\frac{u}{u_\infty^c} - 1 \right)^2 \, dx + q\alpha \int_c^{c'} x^{-\alpha} u$$

where $\chi(t) \equiv 0$ if $s(t) > c(t)$, $\chi(t) := \int_{s(t)}^c \mu u_\infty^c \, dx$ if $s(t) \leq c(t)$.