

Stability in Gagliardo-Nirenberg-Sobolev inequalities

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(joint work with Matteo Bonforte, Bruno Nazaret, Nikita Simonov)

Optimal constants and optimal functions are known in some functional inequalities. The next question is the stability issue: is the difference of the two terms controlling a distance to the set of optimal functions? A famous example is provided by Sobolev's inequalities: in 1991, G. Bianchi and H. Egnell proved that the difference of the two terms is bounded from below by a distance to the manifold of the Aubin-Talenti functions. They argued by contradiction and gave a very elegant although not constructive proof. Since then, estimating the stability constant and giving a constructive proof has always been a challenge.

This contribution focuses on Gagliardo-Nirenberg-Sobolev inequalities. The main tool is based on entropy methods and nonlinear flows. In our method, proving stability amounts to establish, under some constraints, a version of the entropy – entropy production inequality with an improved constant. In simple cases, for instance on the sphere, rather explicit results have been obtained by the *carré du champ* method introduced by D. Bakry and M. Emery. In the Euclidean space, results based on constructive regularity estimates for the solutions of the nonlinear flow have been obtained in a joint research project with Matteo Bonforte, Bruno Nazaret, and Nikita Simonov.

According to [1], on \mathbb{R}^d , $d \geq 3$, there is a positive constant α such that

$$(1) \quad S_d \|\nabla f\|_{L^2(\mathbb{R}^d)}^2 - \|f\|_{L^{2^*}(\mathbb{R}^d)}^2 \geq \alpha \inf_{\varphi \in \mathfrak{M}} \|\nabla f - \nabla \varphi\|_{L^2(\mathbb{R}^d)}^2,$$

where the left-hand side is the difference of the two terms in Sobolev's inequality, with optimal constant S_d , and \mathfrak{M} denotes the manifold of the optimal Aubin-Talenti functions. Various improvements as, *e.g.*, in [4, 7] have been obtained but the question of *constructive* estimates is still widely open: α is obtained by compactness estimates and contradiction arguments, and no good stability estimate is known so far at least with a strong notion of distance as in the right-hand side of (1). See [6] for a result with a weaker norm.

We consider the family of *Gagliardo-Nirenberg-Sobolev inequalities*

$$(2) \quad \|\nabla f\|_2^\theta \|f\|_{p+1}^{1-\theta} \geq \mathcal{C}_{\text{GNS}}(p) \|f\|_{2p}$$

with $\theta = \frac{d(p-1)}{(d+2-p)(d-2)p}$, $p \in (1, +\infty)$ if $d = 1$ or 2 , $p \in (1, p^*]$ if $d \geq 3$, $p^* = \frac{d}{d-2}$. It is known from [5] that equality is achieved if and only if, up to a multiplication by a constant, a translation and a scaling, $f = \mathbf{g}$ where $\mathbf{g}(x) := (1 + |x|^2)^{-1/(p-1)}$. We shall denote the corresponding manifold by \mathfrak{M} as when $p = p^*$. Sobolev's inequality when $d \geq 3$, $p = p^*$, the Euclidean Onofri inequality obtained for $d = 2$ by taking the limit as $p \rightarrow +\infty$ with $f(x) = \mathbf{g}(x) (1 + \frac{1}{2p}(h(x) - \bar{h}))$,

$$\log \left(\int_{\mathbb{R}^2} e^{h-\bar{h}} \frac{dx}{\pi(1+|x|^2)^2} \right) \leq \frac{1}{16\pi} \int_{\mathbb{R}^2} |\nabla h|^2 dx \quad \text{with} \quad \bar{h} = \int_{\mathbb{R}^2} \frac{h(x) dx}{\pi(1+|x|^2)^2}$$

and, as $p \rightarrow 1_+$, the (scale invariant) *Euclidean logarithmic Sobolev inequality*

$$\frac{d}{2} \log \left(\frac{2}{\pi d e} \int_{\mathbb{R}^d} |\nabla f|^2 dx \right) \geq \int_{\mathbb{R}^d} |f|^2 \log |f|^2 dx$$

are all limit cases of (2). Let us define the *deficit functional*

$$\delta[f] := (p-1)^2 \|\nabla f\|_2^2 + 4 \frac{d-p(d-2)}{p+1} \|f\|_{p+1}^{p+1} - \mathcal{K}_{\text{GNS}} \|f\|_2^{2p\gamma}$$

with $\mathcal{K}_{\text{GNS}} = C(p, d) \mathcal{C}_{\text{GNS}}^{2p\gamma}$, $\gamma = \frac{d+2-p(d-2)}{d-p(d-4)}$ for some explicit positive constant $C(p, d)$. A scale optimization shows that (2) is equivalent to the inequality $\delta[f] \geq 0$. Stability results for (2) with non-constructive estimates are known from [3, 10].

With $d \geq 1$, $m \in (1 - 1/d, 1)$, the *fast diffusion* equation in \mathbb{R}^d

$$(3) \quad \frac{\partial u}{\partial t} = \Delta u^m$$

with initial datum $u(t=0, x) = u_0(x) \geq 0$, $u_0 \in L^1(\mathbb{R}^d, (1+|x|^2) dx)$, can be interpreted as the gradient flow of the *entropy* $\mathbf{E} := \int_{\mathbb{R}^d} u^m dx$ with respect to Wasserstein's distance, as it is known from [8]. By the *carré du champ method* (adapted from the work of D. Bakry and M. Emery), we can relate (2) and (3), and obtain a proof of (2). The key property, inspired by the *Rényi entropy powers* of [9] is based on the one hand on the fact that

$$\mathbf{E}' = (1-m)\mathbf{I}$$

where $\mathbf{I} := \int_{\mathbb{R}^d} u |\nabla \mathbf{P}|^2 dx$ is the *Fisher information* and $\mathbf{P} = \frac{m}{m-1} u^{m-1}$ is the *pressure variable*, and on the other hand on the identity

$$-\frac{d}{dt} \log \left(\mathbf{I}^{\frac{1}{2}} \mathbf{E}^{\frac{1-\theta}{\theta(p+1)}} \right) = \int_{\mathbb{R}^d} u^m \|D^2 \mathbf{P} - \frac{1}{d} \Delta \mathbf{P} \text{Id}\|^2 dx + (m-m_1) \int_{\mathbb{R}^d} u^m |\Delta \mathbf{P} + \frac{1}{\mathbf{E}}|^2 dx.$$

Hence $\mathbf{I}^{\frac{1}{2}} \mathbf{E}^{\frac{1-\theta}{\theta(p+1)}}$ is monotone with a limit given by self-similar *Barenblatt functions* or, equivalently by \mathbf{g}^{2p} . With the relations $f = u^{m-1/2}$, $p = 1/(2m-1)$, so that $\int_{\mathbb{R}^d} u dx = \int_{\mathbb{R}^d} f^{2p} dx$, $\mathbf{E} = \int_{\mathbb{R}^d} f^{p+1} dx$ and $\mathbf{I} = (p+1)^2 \int_{\mathbb{R}^d} |\nabla f|^2 dx$, the role of (2) is easily recovered. A rigorous proof goes through the time-dependent rescaling

$$u(t, x) = \frac{1}{\kappa^d R^d} v \left(\tau, \frac{x}{\kappa R} \right) \quad \text{where} \quad \frac{dR}{dt} = R^{1-\mu}, \quad \tau(t) := \frac{1}{2} \log R(t)$$

so that (3) is changed, with same initial datum u_0 and the choice $R(0) = 1$, into the *Fokker-Planck type equation*

$$(4) \quad \frac{\partial v}{\partial \tau} + \nabla \cdot \left[v (\nabla v^{m-1} - 2x) \right] = 0.$$

With now $f = v^{m-1/2}$ such that $\int_{\mathbb{R}^d} f^{2p} dx = \int_{\mathbb{R}^d} \mathbf{g}^{2p} dx$ and again $p = \frac{1}{2m-1}$, with $\mathcal{Q}[v] := \mathcal{I}[v]/\mathcal{F}[v]$, $\delta[f] \geq 0$ is equivalent to $\mathcal{Q}[v] \geq 4$, where $\mathcal{B} := \mathbf{g}^{2p}$ and

$$\mathcal{F}[v] := \int_{\mathbb{R}^d} (\mathcal{B}^{m-1} (v - \mathcal{B}) - \frac{1}{m} (v^m - \mathcal{B}^m)) dx, \quad \mathcal{I}[v] := \int_{\mathbb{R}^d} v |\nabla v^{m-1} + 2x|^2 dx.$$

For any $m \in (1 - 1/d, 1)$, the main result in [2] is the fact that

$$(5) \quad \mathcal{Q}[v(\tau, \cdot)] > 4 + \eta \quad \forall \tau > 0,$$

for some $\eta > 0$, under the conditions that $\int_{\mathbb{R}^d} v_0(1, x) dx = \int_{\mathbb{R}^d} \mathcal{B}(1, x) dx$ and

$$\mathbf{A}[v_0] := \sup_{r>0} r^{\frac{d-p(d-4)}{p-1}} \int_{|x|>r} v_0 dx$$

is finite, with an explicit dependence of η on $\mathbf{A}[v_0]$ and $\mathcal{F}[v_0]$. The method involves:

- (i) An *initial time layer* property: if $\mathcal{Q}[v(T, \cdot)] \geq 4 + \eta$ for some $\eta > 0$ and $T > 0$, then $\mathcal{Q}[v(\tau, \cdot)] \geq 4 + 4\eta e^{-4T}/(4 + \eta - \eta e^{-4T})$ for any $\tau \in [0, T]$.
- (ii) A *threshold time*. Based on a *global Harnack Principle*, there exists some $T > 0$ such that $(1 - \varepsilon)\mathcal{B} \leq v(\tau, \cdot) \leq (1 + \varepsilon)\mathcal{B}$ for any $\tau > T$.
- (iii) An *asymptotic time layer* property: as a consequence of an improved *Hardy-Poincaré* (spectral gap) inequality, (5) holds for $\eta = \eta(\varepsilon)$, for any $\tau \geq T$.

Rewritten in terms of f , the improved entropy – entropy production inequality $\mathcal{Q}[v] > 4 + \eta$ is a stability result.

Theorem 1. *Let $d \geq 1$ and $p \in (1, p^*)$. There is an explicit $\mathcal{C} = \mathcal{C}[f]$ such that, for any $f \in L^{2p}(\mathbb{R}^d, (1 + |x|^2) dx)$ such that $\nabla f \in L^2(\mathbb{R}^d)$ and $\mathbf{A}[f^{2p}] < \infty$,*

$$\delta[f] \geq \mathcal{C}[f] \inf_{\varphi \in \mathfrak{M}} \int_{\mathbb{R}^d} |(p-1)\nabla f + f^p \nabla \varphi^{1-p}|^2 dx.$$

The main point is that the dependence of $\mathcal{C}[f]$ on $\mathbf{A}[f^{2p}]$ and $\mathcal{F}[f^{2p}]$ is explicit and does not degenerate if $f \in \mathfrak{M}$. The critical case $p = p^*$ can also be covered up to an additional scaling. A major open issue is of course to remove the condition $\mathbf{A}[f^{2p}] < \infty$, which requires a new approach.

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