# Stability for a nonlinear Fokker-Planck equation with density-dependent diffusion

#### Renjun Duan

Johann Radon Institute for Computational and Applied Mathematics (RICAM)
Austrian Academy of Sciences

This is a joint work with M. Fornasier and G. Toscani

Concentration en vitesse et en espace dans les modèles cinétiques et diffusifs (chemotaxis, gravitation, swarming)

IHP Institute Henri Poincaré, Paris, 6-7 October 2009

## **Outline**

- Background and motivations
- Main results
  - ► Stability of solutions near equilibrium
  - ► Rate of convergence
- Key ideas in the proof
  - ► Free energy functional
  - ► Hyperbolic-parabolic argument
  - ► Energy-spectrum method

# A flocking model by Cucker-Smale

Motion of m particles (e.g. birds) with  $(x_i, \xi_i) \in \mathbb{R}^n \times \mathbb{R}^n$ :

$$\begin{cases} dx_i = \xi_i dt, \\ d\xi_i = \sum_{j=1}^m U(|x_j - x_i|)(\xi_j - \xi_i) dt. \end{cases}$$

## Proposition (Cucker-Smale 07)

Let

$$U(x) = \frac{C_{n,\gamma}}{(1+|x|^2)\beta}, \quad x \in \mathbb{R}^n.$$

•  $0 \le \beta < \frac{1}{2}$ :

$$\xi_i(t) \to \frac{1}{m} \sum_{j=1}^{m} \xi_j(0), \quad 1 \le i \le m.$$
 (\*)

•  $\beta \geq \frac{1}{2}$ : (\*) holds conditionally.



### Related work

- **Ha-Tadmor:** From particle to kinetic and hydrodynamic descriptions of flocking.
- **Ha-Liu:** A simple proof of the Cucker-Smale flocking dynamics and mean-field limit.
- Carrillo-Fornasier-Rosado-Toscani: Asymptotic flocking dynamics for the kinetic Cucker-Smale model.

Our work will be more related to the study of some nonlinear Fokker-Planck equations.

- Villani: Hypocoercivity.
- **Dolbeault-Mouhot-Schmeiser:** Hypocoercivity for kinetic equations with linear relaxation terms.
- **Guo:** The Landau equation in a periodic box.



### Cucker-Smale model with non-uniform noise

We consider:  $x = (x_1, \dots, x_m), \xi = (\xi_1, \dots, \xi_m)$ ,

$$\begin{cases} dx_{i} = \xi_{i}dt, \\ d\xi_{i} = \sum_{j=1}^{m} U(|x_{j} - x_{i}|)(\xi_{j} - \xi_{i})dt + \sqrt{2\mu \sum_{j=1}^{m} U(|x_{j} - x_{i}|)}dW_{i}. \end{cases}$$

Here, for the *i*<sup>th</sup>-agent,

strength of noise 
$$\propto d_i(x) =: \sum_{j=1}^m U(|x_j - x_i|).$$

So, randomness increases as soon as particles are closer to each other.

### Remark

#### Rewrite

$$d\xi_{i} = d_{i}(x) \underbrace{\left[\sum_{j=1}^{m} w_{ij}(x)\xi_{j} - \xi_{i}\right] dt + \sqrt{2\mu d_{i}(x)} \underbrace{dW_{i}}_{\text{random}},$$
alignment

where

$$\sum_{j=1}^{m} w_{ij}(x) = 1, \quad 1 \le i \le m, \quad \forall x.$$

So, in the model we proposed, the non-uniform noise is such that

strength of alignment  $\propto$  strength of noise,

and both strengths  $\propto d_i(x)$  for each i.

! Stabilize the existing equilibrium states!

## Mean-field limit

Scale

$$U=\frac{\kappa}{m}U_0.$$

Set

$$f^{(m)}(t, x, \xi) = \frac{1}{m} \sum_{i=1}^{m} \delta(x - x_i(t)) \delta(\xi - \xi_i(t)),$$

and assume:  $\exists f(t) \in \mathcal{M}(\mathbb{R}^{2n})$ , s.t.

$$f^{(m)} \to f(t)$$
 in  $w^* - \mathcal{M}(\mathbb{R}^{2n})$  as  $m \to \infty$ .

Then,

$$\partial_t f + \xi \cdot \nabla_x f + \kappa U_0 * \rho_{\xi f} \cdot \nabla_{\xi} f = \kappa U_0 * \rho_f \nabla_{\xi} \cdot (\mu \nabla_{\xi} f + \xi f),$$

with

$$\rho_f(t,x) = \int_{\mathbb{R}^n} f(t,x,\xi) d\xi, \quad \rho_{\xi f}(t,x) = \int_{\mathbb{R}^n} \xi f(t,x,\xi) d\xi.$$

# Nonlinear Fokker-Planck equation and equilibriums

We consider the Cauchy problem

$$\partial_t f + \xi \cdot \nabla_x f + U * \rho_{\xi f} \cdot \nabla_\xi f = U * \rho_f \nabla_\xi \cdot (\nabla_\xi f + \xi f),$$
  
$$f(0, x, \xi) = f_0(x, \xi).$$

Assume that U is continuous in x with

$$U(x) = U(|x|) \ge 0, \quad \int_{\mathbb{R}^n} U(x) dx = 1.$$

The existing equilibrium state is a global Maxwellian (after normalization):

$$\mathbf{M} = \mathbf{M}(\xi) = \frac{1}{(2\pi)^{n/2}} \exp(-|\xi|^2/2)$$
.

#### **Problem:**

! well-posedness and large-time behavior of solutions!



# Try to find a Lyapunov functional

One has

$$\frac{d}{dt}E(f) = -D(f),$$

where

$$E(f) = \iint_{\mathbb{R}^n \times \mathbb{R}^n} \left[ \frac{|\xi|^2}{2} + \log f \right] f dx d\xi,$$

$$D(f) = \iint_{\mathbb{R}^n \times \mathbb{R}^n} \frac{U * \rho_f}{f} |\nabla_{\xi} f + \xi f|^2 dx d\xi - \int_{\mathbb{R}^n} U * \rho_{\xi f} \cdot \rho_{\xi f} dx.$$

**Difficulty:** it is presently unknown if D(f) is non-negative!

Remark: If  $U = \delta_0$ , then  $D(f) = D_0(f)$  with

$$D_0(f) = \iint_{\mathbb{R}^n \times \mathbb{R}^n} \frac{\rho_f}{f} |\nabla_{\xi} f + \xi f|^2 dx d\xi - \int_{\mathbb{R}^n} |\rho_{\xi f}|^2 dx$$
$$= \iint_{\mathbb{R}^n \times \mathbb{R}^n} \frac{\rho_f}{f} |\nabla_{\xi} f + (\xi - \frac{\rho_{\xi f}}{\rho_f}) f|^2 dx d\xi \ge 0.$$

## Observation at the linearized level

Notice D(M) = 0, and

$$\frac{d}{d\epsilon}D(\mathbf{M}+\epsilon\phi)|_{\epsilon=0}=0, \quad \frac{d^2}{d\epsilon^2}D(\mathbf{M}+\epsilon\phi)|_{\epsilon=0}=2\mathbf{L}(\frac{\phi}{\sqrt{\mathbf{M}}}).$$

where L is a linear self-adjoint operator.

Problem: Does L have a sign over some Hilbert space?

(Yes! Non positive definite!

The density-dependent diffusion is responsible for this.)

In fact, decompose

$$L_{\xi}^2 = \mathcal{N} \oplus \mathcal{N}^{\perp}, \quad \mathcal{N} = \operatorname{Span}\{\sqrt{\mathbf{M}}, \xi\sqrt{\mathbf{M}}\},$$

and define P by

$$\mathbf{P}: L_{\xi}^{2} \to \mathcal{N}, \quad u \mapsto \mathbf{P}u \equiv \{a^{u} + b^{u} \cdot \xi\} \sqrt{\mathbf{M}}, \\ a^{u} = \langle \sqrt{\mathbf{M}}, u \rangle, \quad b^{u} = \langle \xi \sqrt{\mathbf{M}}, u \rangle.$$



# **Coercivity of the linearized operator**

Theorem (D.-Fornasier-Toscani 09)

 $-\mathbf{L}$  is coercive in the sense that  $\exists \lambda > 0$ , s.t.

$$\int_{\mathbb{R}^n} \langle -\mathbf{L} u, u \rangle dx \geq \lambda \| \{ \mathbf{I} - \mathbf{P} \} u \|_{\nu}^2 + \frac{1}{2} \| T_{\Delta} b^u \|_{U}^2,$$

for any  $u = u(x, \xi)$ , where

$$||u||_{\nu}^{2} = \iint_{\mathbb{R}^{n} \times \mathbb{R}^{n}} |\nabla_{\xi} u(x, \xi)|^{2} + \nu(\xi) |u(x, \xi)|^{2} d\xi dx d\xi,$$

$$\nu(\xi) = 1 + |\xi|^{2},$$

$$||T_{\Delta} b^{u}||_{U}^{2} = \iint_{\mathbb{R}^{n} \times \mathbb{R}^{n}} U(|x - y|) |b^{u}(x) - b^{u}(y)|^{2} dx dy,$$

$$T_{\Delta} b^{u}(x, y) = b^{u}(x) - b^{u}(y).$$

So, it is possible to found a perturbation theory for the nonlinear Cauchy problem!

### Remarks

• Is there any  $\lambda > 0$  s.t. for any b,

$$||T_{\Delta}b||_U^2 \ge \lambda ||b||^2 ?$$

#### NO! Actually, one can prove

$$\inf_{b \in L_x^2} \|T_{\Delta}b\|_U^2 = \inf_{b \in L_x^2} \iint_{\mathbb{R}^n \times \mathbb{R}^n} U(|x-y|)|b(x) - b(y)|^2 dx dy = 0.$$

 Is there any coercivity estimate on the linearized operator L corresponding to the equation with uniform diffusion

$$\partial_t f + \xi \cdot \nabla_x f + U * \rho_{\xi f} \cdot \nabla_{\xi} f = U * \rho_f \nabla_{\xi} \cdot (\xi f) + \Delta_{\xi} f ?$$

NO! This need a long calculation.



### Reformulation

Recall

$$\partial_t f + \xi \cdot \nabla_x f + U * \rho_{\xi f} \cdot \nabla_{\xi} f = U * \rho_f \nabla_{\xi} \cdot (\nabla_{\xi} f + \xi f),$$
  
$$f(0, x, \xi) = f_0(x, \xi).$$

Set the perturbation  $u = u(t, x, \xi)$  by

$$f = \mathbf{M} + \sqrt{\mathbf{M}}u.$$

Then,

$$\partial_t u + \xi \cdot \nabla_{\mathsf{x}} u + U * \rho_{\xi \sqrt{\mathsf{M}} u} \cdot \nabla_{\xi} u = \mathbf{L} u + \Gamma(u, u),$$

where

$$\mathbf{L}u = \underbrace{\Delta_{\xi}u + \frac{1}{4}(2n - |\xi|^2)u}_{=:\mathsf{L}_{FP}u} + \underbrace{U * \rho_{\xi\sqrt{\mathsf{M}}u} \cdot \xi\sqrt{\mathbf{M}}}_{=:\mathsf{A}u},$$

$$\Gamma(u, u) = U * \rho_{\sqrt{M}u} \mathbf{L}_{FP} u + \frac{1}{2} U * \rho_{\xi\sqrt{M}u} \cdot \xi u.$$



# Nonlinear asymptotical stability near equilibrium

### Our goal is to prove

• stability:

 $\exists \delta > 0, \ C > 1 \text{ s.t. if } u_0 = \mathbf{M}^{-1/2}(f_0 - \mathbf{M}) \in B_{\delta}, \text{ where } B_{\delta}$  is a smooth neighborhood of zero, then

$$\exists ! u \in C_b([0,\infty); B_{C\delta})$$

with  $u(0) = u_0$ .

### local existence + a priori estimates

rate of convergence:

How fast for  $u(t) \to 0$  in some smooth topology?

spectral analysis + energy-spectrum method



### Main results: nonlinear case

Theorem (D.-Fornasier-Toscani 09)

Let  $n \ge 3$ ,  $N \ge 2[n/2] + 2$ . U is supposed as before. Assume

$$f_0 \equiv \mathbf{M} + \sqrt{\mathbf{M}} u_0 \ge 0, \quad \|u_0\|_{H^N_{x,\varepsilon}} \le \delta$$

for some  $\delta > 0$ .

(i)  $\exists ! u(t, x, \xi) \in C([0, \infty); H_{x, \xi}^N)$  with

$$f \equiv \mathbf{M} + \sqrt{\mathbf{M}}u \ge 0$$
,  $\sup_{t \ge 0} \|u(t)\|_{H^N_{x,\xi}} \le C \|u_0\|_{H^N_{x,\xi}}$ .

for some C > 1.

(ii) If  $||u_0||_{Z_1}$  with  $Z_1 = L_\xi^2(L_x^1)$  is bounded and  $||\xi \nabla_x u_0||$  is small enough, then

$$\sup_{t>0} (1+t)^{\frac{n}{4}} \|u(t)\|_{H^{N}_{x,\xi}} \leq C \left( \|u_{0}\|_{H^{N}_{x,\xi}} + \|u_{0}\|_{Z_{1}} \right),$$

for some C.



### **Iteration**

# Define an approximate solution sequence $(f^m)_{m=0}^{\infty}$ by solving Cauchy problems iteratively

$$\begin{cases} \partial_{t} f^{m+1} + \xi \cdot \nabla_{x} f^{m+1} + U * \rho_{\xi f^{m}} \cdot \nabla_{\xi} f^{m+1} \\ = U * \rho_{f^{m}} \nabla_{\xi} \cdot (\nabla_{\xi} f^{m+1} + \xi f^{m+1}), \\ f^{m+1} \equiv \mathbf{M} + \sqrt{\mathbf{M}} u^{m+1}, \\ f^{m+1}|_{t=0} = f_{0} \equiv \mathbf{M} + \sqrt{\mathbf{M}} u_{0}, \end{cases}$$

or equivalently in terms of  $u^m(t, x, \xi)$ :

$$\begin{cases} \partial_{t}u^{m+1} + \xi \cdot \nabla_{x}u^{m+1} + U * b^{u^{m}} \cdot \nabla_{\xi}u^{m+1} \\ = \mathbf{L}_{FP}u^{m+1} + \Gamma(u^{m}, u^{m+1}) + \mathbf{A}u^{m}, \\ u^{m+1}|_{t=0} = u_{0}, \end{cases}$$

where  $m \ge 0$ , and  $u^0 \equiv 0$  is set at initial step.



### Local existence

#### **Define**

$$X(0,T;M) = \left\{ v \in C([0,T]; H^{N}(\mathbb{R}^{n} \times \mathbb{R}^{n})) : \atop \sup_{0 \le t \le T} \|v(t)\|_{H^{N}_{x,\xi}} \le M, \quad \mathbf{M} + \sqrt{\mathbf{M}}v \ge 0 \right\}.$$

#### Theorem

$$\exists T_* > 0$$
,  $\epsilon_0 > 0$ ,  $M_0 > 0$  s.t. if  $u_0 \in H^N_{x,\xi}$  with

$$f_0 \equiv \mathbf{M} + \sqrt{\mathbf{M}} u_0 \ge 0, \quad \|u_0\|_{H^N_{x,\xi}} \le \epsilon_0,$$

then, for each  $m \geq 1$ ,  $u^m$  is well-defined with  $u^m \in X(0, T_*; M_0)$ . Furthermore,  $(u^m)_{m \geq 0}$  is a Cauchy sequence in  $C([0, T_*]; H^{N-1}_{X,\xi})$ , and the corresponding limit function denoted by u belongs to  $X(0, T_*; M_0)$ , and u is a solution to the Cauchy problem. Meanwhile,  $\exists$  at most one solution in  $X(0, T_*; M_0)$ .

## Uniform a priori estimate-1

### Define the temporal functionals

$$\begin{split} \mathcal{E}(u(t)) &= \|u(t)\|_{L_{\xi}^{2}(H_{x}^{N})}^{2} + \kappa_{1}\mathcal{E}_{free}(u(t)) \\ &+ \kappa_{2} \sum_{1 \leq k \leq N} C_{k} \sum_{\substack{|\beta| = k \\ |\alpha| + |\beta| \leq N}} \|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \{\mathbf{I} - \mathbf{P}\} u\|^{2}, \\ \mathcal{D}(u(t)) &= \sum_{|\alpha| + |\beta| \leq N} \|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \{\mathbf{I} - \mathbf{P}\} u(t)\|_{\nu}^{2} + \sum_{|\alpha| \leq N} \|T_{\Delta} \partial_{x}^{\alpha} b^{u}(t)\|_{U}^{2} \\ &+ \|\nabla_{x} (a^{u}, b^{u})(t)\|_{H^{N-1}}^{2}, \end{split}$$

where  $0 < \kappa_2 \ll \kappa_1 \ll 1$ , and

$$|\mathcal{E}_{free}(u(t))| \leq ||u(t)||_{L^2_{\xi}(H^N_x)}^2.$$

#### Notice

$$\mathcal{E}(u(t)) \sim \|u(t)\|_{H^{N}_{v,\varepsilon}}^{2}.$$



# Uniform a priori estimate-2

One can show

$$\frac{d}{dt}\mathcal{E}(u(t)) + \lambda \mathcal{D}(u(t)) \le C\mathcal{E}(u(t))\mathcal{D}(u(t)).$$

Then, under the a priori assumption on smallness of

$$\sup_{0 \le t \le T} \mathcal{E}(u(t)),$$

one has the Lyapunov inequality

$$\frac{d}{dt}\mathcal{E}(u(t)) + \lambda \mathcal{D}(u(t)) \le 0.$$

for any  $0 \le t \le T$ .

# **Global solutions: Continuity argument**

**Define** 

$$T_* = \sup\{t : \sup_{0 \le s \le t} \|u(s)\|_{H^N_{x,\xi}}^2 \le M\}.$$

Using

local existence + uniform a priori estimates,

then,

$$\mathcal{E}(u_0) \sim \|u_0\|_{H^N_{x,\xi}}^2 \ll 1$$
,  $\exists M > 0 \Rightarrow T_* = \infty$ : global solution.

# Key point: Construction of free energy functional-1

The free energy functional is responsible to obtain the macroscopic dissipation

$$\|\nabla_{x}(a^{u},b^{u})\|_{H_{v}^{N-1}}^{2}$$
.

Make the macro-micro decomposition:

$$\begin{cases} u(t, x, \xi) = \mathbf{P}u + \{\mathbf{I} - \mathbf{P}\}u, \\ \mathbf{P}u \equiv \{a^u + b^u \cdot \xi\}\sqrt{\mathbf{M}}, \\ a^u = \langle \sqrt{\mathbf{M}}, u \rangle, \quad b^u = \langle \xi\sqrt{\mathbf{M}}, u \rangle, \end{cases}$$

**Define**  $A = (A_{ij}(\cdot))_{n \times n}$  by

$$A_{ij}(u) = \int_{\mathbb{D}^n} (\xi_i \xi_j - 1) \sqrt{\mathbf{M}} u d\xi.$$

One can obtain a series of equations satisfied by  $a^u$ ,  $b^u$ 



# **Key point: Construction of free energy functional-2**

Macro balance laws:

$$\begin{split} \partial_t a^u + \nabla_x \cdot b^u &= 0, \\ \partial_t b^u_i + \partial_i a^u - (U * b^u_i - b^u_i) + U * a^u b^u_i - U * b^u_i a^u \\ &+ \sum_{i=1}^n \partial_i A_{ij} (\{\mathbf{I} - \mathbf{P}\}u) = 0, \end{split}$$

Evolution of high-order moments:

$$\partial_t A_{ii}(\{\mathbf{I} - \mathbf{P}\}u) + \partial_i b_i^u - U * b_i^u b_i^u = A_{ii}(I + r),$$
  
$$\partial_t A_{ij}(\{\mathbf{I} - \mathbf{P}\}u) + \partial_i b_j^u + \partial_j b_i^u - U * b_i^u b_j^u - U * b_j^u b_i^u$$
  
$$= A_{ij}(I + r), i \neq j,$$

#### where

$$I = -\xi \cdot \nabla_{x} \{ \mathbf{I} - \mathbf{P} \} u + \mathbf{L}_{FP} \{ \mathbf{I} - \mathbf{P} \} u,$$
  

$$r = U * a^{u} \mathbf{L}_{FP} \{ \mathbf{I} - \mathbf{P} \} u + \frac{1}{2} U * b^{u} \cdot \xi \{ \mathbf{I} - \mathbf{P} \} u - U * b^{u} \cdot \nabla_{\xi} \{ \mathbf{I} - \mathbf{P} \} u.$$

# **Key point: Construction of free energy functional-3**

**Define**  $\mathcal{E}_{free}(u(t))$  by

$$\begin{split} \mathcal{E}_{free}(u(t)) &= 3 \sum_{|\alpha| \leq N-1} \sum_{j} \sum_{i \neq j} \int_{\mathbb{R}^{n}} A_{ii} (\partial_{x}^{\alpha} \partial_{j} \{\mathbf{I} - \mathbf{P}\} u) \partial_{x}^{\alpha} b_{j}^{u} dx \\ &- 3 \sum_{|\alpha| \leq N-1} \sum_{ij} \int_{\mathbb{R}^{n}} A_{ij} (\partial_{x}^{\alpha} \partial_{i} \{\mathbf{I} - \mathbf{P}\} u) \partial_{x}^{\alpha} b_{j}^{u} dx \\ &+ \sum_{|\alpha| \leq N-1} \int_{\mathbb{R}^{n}} \partial_{x}^{\alpha} \nabla_{x} a^{u} \cdot \partial_{x}^{\alpha} b^{u} dx. \end{split}$$

Then,

$$\frac{d}{dt} \mathcal{E}_{free}(u(t)) + \lambda \|\nabla_{x}(a^{u}, b^{u})\|_{H_{x}^{N-1}}^{2} 
\leq C \sum_{|\alpha| \leq N} (\|T_{\Delta}\partial_{x}^{\alpha}b^{u}\|_{U}^{2} + \|\partial_{x}^{\alpha}\{\mathbf{I} - \mathbf{P}\}u\|^{2}) 
+ C \|(a^{u}, b^{u})\|_{H_{x}^{N}}^{2} (\|\nabla_{x}(a^{u}, b^{u})\|_{H_{x}^{N-1}}^{2} + \sum_{|\alpha| \leq N} \|\partial_{x}^{\alpha}\{\mathbf{I} - \mathbf{P}\}u\|^{2}).$$

# Representation of solutions

Consider the Cauchy problem of the linearized equation with a nonhomogeneous source:

$$\begin{cases} \partial_t u = \mathbf{B}u + h, & t > 0, x \in \mathbb{R}^n, \\ u|_{t=0} = u_0, & x \in \mathbb{R}^n, \end{cases}$$

where  $n \ge 1$  is the spatial dimension,  $h = h(t, x, \xi)$  and  $u_0 = u_0(x, \xi)$  are given, and the linear operator B is defined by

$$\mathbf{B} = -\xi \cdot \nabla_{\mathsf{X}} + \mathbf{L}, \quad \mathbf{L} = \mathbf{L}_{\mathsf{FP}} + \mathbf{A}.$$

Formally, the solution can be written as the Duhamel formula

$$u(t) = e^{Bt}u_0 + \int_0^t e^{B(t-s)}h(s)ds.$$

## Main result: Linearized case

Set

$$\sigma_{q,m} = \frac{n}{2} \left( \frac{1}{q} - \frac{1}{2} \right) + \frac{m}{2}.$$

### Theorem

Let  $1 \le q \le 2$ ,  $n \ge 1$ , and let  $Z_q = L_{\mathcal{E}}^2(L_x^1)$ .

(i) 
$$\|\partial_x^{\alpha} e^{Bt} u_0\| \le C(1+t)^{-\sigma_{q,m}} (\|\partial_x^{\alpha'} u_0\|_{Z_q} + \|\partial_x^{\alpha} u_0\|),$$

with  $m = |\alpha - \alpha'|$ .

(ii) If  $\mathbf{P}h = 0$ , i.e. h is microscopic, then

$$\begin{split} \left\| \partial_{x}^{\alpha} \int_{0}^{t} e^{\mathsf{B}(t-s)} h(s) ds \right\|^{2} \\ & \leq C \int_{0}^{t} (1+t-s)^{-2\sigma_{q,m}} (\|\nu^{-1/2} \partial_{x}^{\alpha'} h(s)\|_{Z_{q}}^{2} + \|\nu^{-1/2} \partial_{x}^{\alpha} h(s)\|^{2}) ds, \end{split}$$

$$J_0$$

with  $m = |\alpha - \alpha'|$ .



# Spectral analysis-1

**Define** 

$$\mathcal{E}^{I}(\widehat{u}(t,k)) = \|\widehat{u}(t,k)\|_{L_{\xi}^{2}}^{2} + \kappa \operatorname{Re} \mathcal{E}^{I}_{free}(\widehat{u}(t,k))$$

for a small constant  $\kappa > 0$ , where

$$\mathcal{E}_{free}^{I}(\widehat{u}(t,k)) = 3\sum_{j} \sum_{i \neq j} \frac{\mathrm{i}k_{j}}{1 + |k|^{2}} (A_{ii}(\{\mathbf{I} - \mathbf{P}\}\widehat{u}) \mid \widehat{b_{j}^{u}})$$

$$-3\sum_{ij} \frac{\mathrm{i}k_{i}}{1 + |k|^{2}} (A_{ij}(\{\mathbf{I} - \mathbf{P}\}\widehat{u}) \mid \widehat{b_{j}^{u}})$$

$$-\frac{\mathrm{i}k}{1 + |k|^{2}} \cdot (\widehat{b^{u}} \mid \widehat{a^{u}}).$$

Notice  $|\mathcal{E}_{free}^{l}(\widehat{u}(t,k))| \leq C \|\widehat{u}(t,k)\|_{L_{\varepsilon}^{2}}^{2}$ . Then,

$$\mathcal{E}^{I}(\widehat{u}(t,k)) \sim \|\widehat{u}(t,k)\|_{L_{\varepsilon}^{2}}^{2}$$

# Spectral analysis-2

#### One can show

$$\frac{\partial}{\partial t} \mathcal{E}^{I}(\widehat{u}(t,k)) + \frac{\lambda |k|^2}{1 + |k|^2} \mathcal{E}^{I}(\widehat{u}(t,k)) \leq C \|\nu^{-1/2} \widehat{h}(t,k)\|_{L_{\xi}^2}^2,$$

### Key idea: Make the Fourier analysis on

$$\begin{split} &\partial_t a^u + \nabla_x \cdot b^u = 0, \\ &\partial_t b^u_i + \partial_i a^u - (U * b^u_i - b^u_i) + \sum_{j=1}^n \partial_j A_{ij} (\{\mathbf{I} - \mathbf{P}\}u) = 0, \\ &\partial_t A_{ii} (\{\mathbf{I} - \mathbf{P}\}u) + \partial_i b^u_i = A_{ii} (I + h), \\ &\partial_t A_{ij} (\{\mathbf{I} - \mathbf{P}\}u) + \partial_i b^u_j + \partial_j b^u_i = A_{ij} (I + h), \quad i \neq j, \end{split}$$

and use Kawashima's hyperbolic-parabolic argument. Then, one has



# Spectral analysis-3

$$\begin{split} &\frac{\partial}{\partial t} \text{Re } \mathcal{E}_{free}^{I}(\widehat{u}(t,k)) + \frac{|k|^{2}}{4(1+|k|^{2})} (|\widehat{a^{u}}|^{2} + |\widehat{b^{u}}|^{2}) \\ &\leq \frac{1 - \text{Re } \widehat{U}}{1+|k|^{2}} |\widehat{b^{u}}|^{2} + \frac{C}{1+|k|^{2}} ||\nu^{-1/2}\widehat{h}||_{L_{\xi}^{2}}^{2} + C ||\{\mathbf{I} - \mathbf{P}\}\widehat{u}||_{L_{\xi}^{2}}^{2}. \end{split}$$

#### Combined with

$$\frac{1}{2} \frac{\partial}{\partial t} \|\widehat{u}(t,k)\|_{L_{\xi}^{2}} + \lambda |\{\mathbf{I} - \mathbf{P}\}\widehat{u}|_{\nu}^{2} + (1 - \operatorname{Re} \widehat{U})|\widehat{b^{u}}|^{2} \\
\leq C \|\nu^{-1/2}\widehat{h}(t,k)\|_{L_{\xi}^{2}}^{2}.$$

The proper linear combination leads to the desired estimate. Then,

$$\|\widehat{u}(t,k)\|_{L_{\xi}^{2}}^{2} \leq Ce^{-\frac{\lambda|k|^{2}}{1+|k|^{2}}t}\|\widehat{u_{0}}(k)\|_{L_{\xi}^{2}}^{2} + C\int_{0}^{t}e^{-\frac{\lambda|k|^{2}}{1+|k|^{2}}(t-s)}\|\nu^{-1/2}\widehat{h}(s,k)\|_{L_{\xi}^{2}}^{2}ds,$$

for any t > 0 and  $k \in \mathbb{R}^n$ .



# Rate of convergence in the nonlinear case

Write the nonlinear equation as

$$u(t) = e^{Bt}u_0 + \int_0^t e^{B(t-s)}G(s)ds,$$

where the source term G is denoted by

$$G = \Gamma(u, u) - U * b^{u} \cdot \nabla_{\varepsilon} u.$$

Energy-spectrum method + Weighted estimates: Use

$$\frac{d}{dt}\mathcal{E}(u(t)) + \lambda \mathcal{E}(u(t)) \le C||u(t)||^2.$$

#### Lemma

If  $||u_0||_{Z_1}$  is bounded, then

$$||u(t)||^{2} \leq C(\mathcal{E}(u_{0}) + ||u_{0}||_{Z_{1}}^{2})(1+t)^{-\frac{n}{2}} + C\int_{0}^{t} (1+t-s)^{-\frac{n}{2}} \mathcal{E}(u(s))[\mathcal{E}(u(s)) + ||\xi\{\mathbf{I} - \mathbf{P}\}u(s)||^{2}]ds + C\left[\int_{0}^{t} (1+t-s)^{-\frac{n}{4}} \mathcal{E}(u(s))ds\right]^{2}.$$

# Thanks for your attention!