

# Uniform in time propagation of chaos for the 2D vortex model and other singular stochastic systems.

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# Idea

*In a system of  $N$  interacting particles, as  $N$  increases, two particles become more and more statistically independent.*

# Formal limit of SDE

$N$ -particle system on the torus  $\mathbb{T}^d$

$$dX_t^i = \sqrt{2}dB_t^i + \frac{1}{N} \sum_{j=1}^N K(X_t^i - X_t^j)dt.$$

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Limit as  $N$  tends to infinity ?

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$$dX_t^i = \sqrt{2}dB_t^i + K * \mu_t^N(X_t^i)dt,$$

$$\mu_t^N := \frac{1}{N} \sum_{i=1}^N \delta_{X_t^i}.$$

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Limit as  $N$  tends to infinity? Formally

$$\begin{cases} d\bar{X}_t = \sqrt{2}dB_t + K * \bar{\rho}_t(\bar{X}_t)dt, \\ \bar{\rho}_t = \text{Law}(\bar{X}_t). \end{cases}$$

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$$dX_t^i = \sqrt{2}dB_t^i + \frac{1}{N} \sum_{j=1}^N K(X_t^i - X_t^j)dt. \quad (\text{PS})$$

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# Liouville equations

For the particle system

$$dX_t^i = \sqrt{2}dB_t^i + \frac{1}{N} \sum_{j=1}^N K(X_t^i - X_t^j) dt$$

$\longleftrightarrow$

$$\partial_t \rho_t^N = - \sum_{i=1}^N \nabla_{x_i} \cdot \left( \left( \frac{1}{N} \sum_{j=1}^N K(x_i - x_j) \right) \rho_t^N \right) + \sum_{i=1}^N \Delta_{x_i} \rho_t^N.$$

For the non linear equation

$$\begin{cases} d\bar{X}_t = \sqrt{2}dB_t + K * \bar{\rho}_t(\bar{X}_t) dt, \\ \bar{\rho}_t = \text{Law}(\bar{X}_t). \end{cases} \quad \longleftrightarrow \quad \partial_t \bar{\rho}_t = -\nabla \cdot (\bar{\rho}_t (K * \bar{\rho}_t)) + \Delta \bar{\rho}_t.$$

# Main example : 2D vortex model

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On the assumptions

The Biot-Savart kernel, defined in  $\mathbb{R}^2$  by

$$K(x) = \frac{1}{2\pi} \frac{x^\perp}{|x|^2} = \frac{1}{2\pi} \left( -\frac{x_2}{|x|^2}, \frac{x_1}{|x|^2} \right).$$

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Consider the 2D incompressible Navier-Stokes system on  $u \in \mathbb{R}^2$

$$\begin{aligned} \partial_t u &= -u \cdot \nabla u - \nabla p + \Delta u \\ \nabla \cdot u &= 0, \end{aligned}$$

where  $p$  is the local pressure.

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$$\partial_t \omega = -\nabla \cdot ((K * \omega) \omega) + \Delta \omega.$$

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**Goal :** Obtain a limit " $\rho_t^N \rightarrow \bar{\rho}_t$ " as  $N$  tends to infinity for this Biot-Savart kernel.

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# Propagation of chaos

*In a system of  $N$  interacting particles, as  $N$  increases, two particles become more and more statistically independent.*

# Propagation of chaos

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On the assumptions

*In a system of  $N$  interacting particles, as  $N$  increases, two particles become more and more statistically independent.*

To quantify this "more and more", we compare the law of any subset of  $k$  particles within the  $N$  particles system to the law of  $k$  independent non-linear particles.

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To quantify this "more and more", we compare the law of any subset of  $k$  particles within the  $N$  particles system to the law of  $k$  independent non-linear particles.

We denote, for any  $k \leq N$

$$\rho_t^{k,N}(x_1, \dots, x_k) = \int_{\mathbb{T}^{(N-k)d}} \rho_t^N(x_1, \dots, x_N) dx_{k+1} \dots dx_N$$
$$\bar{\rho}_t^k = \bar{\rho}_t^{\otimes k}$$

# (Rescaled) relative entropy

## Definition

Let  $\mu$  and  $\nu$  be two probability measures on  $\mathbb{T}^{dN}$ . We consider the rescaled relative entropy

$$\mathcal{H}_N(\nu, \mu) = \begin{cases} \frac{1}{N} \mathbb{E}_\mu \left( \frac{d\nu}{d\mu} \log \frac{d\nu}{d\mu} \right) & \text{if } \nu \ll \mu, \\ +\infty & \text{otherwise.} \end{cases}$$

# Results

## Theorem (adapted from Jabin-Wang ('18))

*Under some assumptions (satisfied by the Biot-Savart kernel) there are constants  $C_1$  and  $C_2$  such that for all  $N \in \mathbb{N}$ , all exchangeable probability density  $\rho_0^N$  and all  $t \geq 0$*

$$\mathcal{H}_N(\rho_t^N, \bar{\rho}_t^N) \leq e^{C_1 t} \left( \mathcal{H}_N(\rho_0^N, \bar{\rho}_0^N) + \frac{C_2}{N} \right)$$

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## Theorem (Guillin-LB-Monmarché ('21))

*Under some assumptions (satisfied by the Biot-Savart kernel) there are constants  $C_1$ ,  $C_2$  and  $C_3$  such that for all  $N \in \mathbb{N}$ , all exchangeable probability density  $\rho_0^N$  and all  $t \geq 0$*

$$\mathcal{H}_N(\rho_t^N, \bar{\rho}_t^N) \leq C_1 e^{-C_2 t} \mathcal{H}_N(\rho_0^N, \bar{\rho}_0^N) + \frac{C_3}{N}$$

# Various distances

For  $\mathbf{x} = (x_i)_{i \in [1, N]} \in \mathbb{T}^{dN}$ , we write  $\pi(\mathbf{x}) = \frac{1}{N} \sum_{i=1}^N \delta_{x_i}$  the associated empirical measure.

## Corollary

*Under some assumptions (satisfied by the Biot-Savart kernel), assuming moreover that  $\rho_0^N = \bar{\rho}_0^N$ , there is a constant  $C$  such that for all  $k \leq N \in \mathbb{N}$  and all  $t \geq 0$ ,*

$$\|\rho_t^{k, N} - \bar{\rho}_t^k\|_{L^1} + \mathcal{W}_2\left(\rho_t^{k, N}, \bar{\rho}_t^k\right) \leq C \left( \left\lfloor \frac{N}{k} \right\rfloor \right)^{-\frac{1}{2}}$$

and

$$\mathbb{E}_{\rho_t^N}(\mathcal{W}_2(\pi(\mathbf{X}), \bar{\rho}_t)) \leq C\alpha(N)$$

where  $\alpha(N) = N^{-1/2} \ln(1 + N)$  if  $d = 2$  and  $\alpha(N) = N^{-1/d}$  if  $d > 2$ .



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We write

$$\mathcal{H}_N(t) = \mathcal{H}_N(\rho_t^N, \bar{\rho}_t^N), \quad \mathcal{I}_N(t) = \frac{1}{N} \sum_i \int_{\mathbb{T}^{dN}} \rho_t^N \left| \nabla_{x_i} \log \frac{\rho_t^N}{\bar{\rho}_t^N} \right|^2 d\mathbf{X}^N.$$

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It has been shown, by Jabin-Wang, that

$$\begin{aligned} \frac{d}{dt} \mathcal{H}_N(t) &\leq -\mathcal{I}_N(t) \\ &\quad - \frac{1}{N^2} \sum_{i,j} \int_{\mathbb{T}^{dN}} \rho_t^N (K(x_i - x_j) - K * \rho(x_i)) \cdot \nabla_{x_i} \log \bar{\rho}_t^N d\mathbf{X}^N \\ &\quad - \frac{1}{N^2} \sum_{i,j} \int_{\mathbb{T}^{dN}} \rho_t^N (\operatorname{div} K(x_i - x_j) - \operatorname{div} K * \bar{\rho}_t(x_i)) d\mathbf{X}^N. \end{aligned}$$

# Assumptions ?

$$\mathbf{Goal : } K(x) = \frac{1}{2\pi} \frac{x^\perp}{|x|^2} = \frac{1}{2\pi} \left( -\frac{x_2}{|x|^2}, \frac{x_1}{|x|^2} \right)$$

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- $\bar{\rho} \in C^\infty(\mathbb{R}^+ \times \mathbb{T}^d)$  and there is  $\lambda > 1$ , s.t  $\frac{1}{\lambda} \leq \bar{\rho} \leq \lambda$

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# Step one : Time evolution of the relative entropy

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## Step two : Integration by part

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On the assumptions

We are left with

$$\begin{aligned} \frac{d}{dt} \mathcal{H}_N(t) &\leq - \mathcal{I}_N(t) \\ &\quad - \frac{1}{N^2} \sum_{i,j} \int_{\mathbb{T}^{dN}} \rho_t^N (K(x_i - x_j) - K * \rho(x_i)) \cdot \nabla_{x_i} \log \bar{\rho}_t^N d\mathbf{X}^N. \end{aligned}$$

**Idea :** Use the regularity of  $\bar{\rho}$  to deal with the singularity of  $K$

## Step two : Integration by part

We are left with

$$\frac{d}{dt} \mathcal{H}_N(t) \leq - \mathcal{I}_N(t) - \frac{1}{N^2} \sum_{i,j} \int_{\mathbb{T}^{dN}} \rho_t^N (K(x_i - x_j) - K * \rho(x_i)) \cdot \nabla_{x_i} \log \bar{\rho}_t^N d\mathbf{X}^N.$$

**Idea :** Use the regularity of  $\bar{\rho}$  to deal with the singularity of  $K$

**Remark :** Notice that, for the Biot-Savart kernel on the whole space  $\mathbb{R}^2$

$$\tilde{K}(x) = \frac{1}{2\pi} \frac{x^\perp}{|x|^2},$$

we have  $\tilde{K} = \nabla \cdot \tilde{V}$  with

$$\tilde{V}(x) = \frac{1}{2\pi} \begin{pmatrix} -\arctan\left(\frac{x_1}{x_2}\right) & 0 \\ 0 & \arctan\left(\frac{x_2}{x_1}\right) \end{pmatrix}.$$

# Assumptions ?

$$\text{Goal : } K(x) = \frac{1}{2\pi} \frac{x^\perp}{|x|^2} = \frac{1}{2\pi} \left( -\frac{x_2}{|x|^2}, \frac{x_1}{|x|^2} \right)$$

## Justifying the calculations

- There is  $\lambda > 1$  such that  $\bar{\rho}_0 \in C_\lambda^\infty(\mathbb{T}^d)$   
 $\implies \bar{\rho} \in C_\lambda^\infty(\mathbb{R}^+ \times \mathbb{T}^d)$  (Ben-Artzi '94)
- $\rho^N \in C_\lambda^\infty(\mathbb{R}^+ \times \mathbb{T}^{Nd})$  (???)

## Dealing with the terms

- In the sense of distributions,  $\nabla \cdot K = 0$ .
- There is a matrix field  $V \in L^\infty$  such that  $K = \nabla \cdot V$ , i.e for  $1 \leq \alpha \leq d$ ,  $K_\alpha = \sum_{\beta=1}^d \partial_\beta V_{\alpha,\beta}$  (Phuc-Torres '08).

## Step two : Integration by part

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For all  $t \geq 0$ ,

$$\frac{d}{dt} \mathcal{H}_N(t) \leq A_N(t) + \frac{1}{2} B_N(t) - \frac{1}{2} \mathcal{I}_N(t),$$

with

$$A_N(t) := \frac{1}{N^2} \sum_{i,j} \int_{\mathbb{T}^{dN}} \rho_t^N (V(x_i - x_j) - V * \bar{\rho}(x_i)) : \frac{\nabla_{x_i}^2 \bar{\rho}_t^N}{\bar{\rho}_t^N} d\mathbf{X}^N$$

$$B_N(t) := \frac{1}{N} \sum_i \int_{\mathbb{T}^{dN}} \rho_t^N \frac{|\nabla_{x_i} \bar{\rho}_t^N|^2}{|\bar{\rho}_t^N|^2} \left| \frac{1}{N} \sum_j V(x_i - x_j) - V * \bar{\rho}(x_i) \right|^2 d\mathbf{X}^N.$$

# Step three : Change of reference measure and large deviation estimates

## Lemma

For two probability densities  $\mu$  and  $\nu$  on a set  $\Omega$ , and any  $\Phi \in L^\infty(\Omega)$ ,  
 $\eta > 0$  and  $N \in \mathbb{N}$ ,

$$\mathbb{E}^\mu \Phi \leq \eta \mathcal{H}_N(\mu, \nu) + \frac{\eta}{N} \log \mathbb{E}^\nu e^{N\Phi/\eta}.$$

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## Theorem (Jabin-Wang '18)

Consider any probability measure  $\mu$  on  $\mathbb{T}^d$ ,  $\epsilon > 0$  and a scalar function  $\psi \in L^\infty(\mathbb{T}^d \times \mathbb{T}^d)$  with  $\|\psi\|_{L^\infty} < \frac{1}{2\epsilon}$  and such that for all  $z \in \mathbb{T}^d$ ,  $\int_{\mathbb{T}^d} \psi(z, x) \mu(dx) = 0$ . Then there exists a constant  $C$  such that

$$\int_{\mathbb{T}^{dN}} \exp\left(\frac{1}{N} \sum_{j_1, j_2=1}^N \psi(x_1, x_{j_1}) \psi(x_1, x_{j_2})\right) \mu^{\otimes N} d\mathbf{X}^N \leq C,$$

where  $C$  depends on

$$\alpha = (\epsilon \|\psi\|_{L^\infty})^4 < 1, \quad \beta = \left(\sqrt{2\epsilon} \|\psi\|_{L^\infty}\right)^4 < 1.$$

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On the assumptions

## Theorem (Jabin-Wang '18)

Consider any probability measure  $\mu$  on  $\mathbb{T}^d$  and  $\phi \in L^\infty(\mathbb{T}^d \times \mathbb{T}^d)$  with

$$\gamma := \left(1600^2 + 36e^4\right) \left(\sup_{\rho \geq 1} \frac{\|\sup_z |\phi(\cdot, z)|\|_{L^\rho(\mu)}}{\rho}\right)^2 < 1.$$

Assume that  $\phi$  satisfies the following cancellations

$$\forall z \in \mathbb{T}^d, \quad \int_{\mathbb{T}^d} \phi(x, z) \mu(dx) = 0 = \int_{\mathbb{T}^d} \phi(z, x) \mu(dx).$$

Then, for all  $N \in \mathbb{N}$ ,

$$\int_{\mathbb{T}^{dN}} \exp\left(\frac{1}{N} \sum_{i,j=1}^N \phi(x_i, x_j)\right) \mu^{\otimes N} d\mathbf{X}^N \leq \frac{2}{1-\gamma} < \infty.$$

# Conclusion

For all  $t \geq 0$ ,

$$\frac{d}{dt} \mathcal{H}_N(t) \leq C \left( \mathcal{H}_N(t) + \frac{1}{N} \right) - \frac{1}{2} \mathcal{I}_N(t),$$

with

$$C = \hat{C}_1 \|\nabla^2 \bar{\rho}_t\|_{L^\infty} \|V\|_{L^\infty} \lambda + \hat{C}_2 \|V\|_{L^\infty}^2 \lambda^2 d^2 \|\nabla \bar{\rho}_t\|_{L^\infty}^2$$

where  $\hat{C}_1, \hat{C}_2$  are universal constants.

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Two goals :

- A logarithmic Sobolev inequality for  $\bar{\rho}^N : \mathcal{H}_N(t) \leq C\mathcal{I}_N(t)$

# Step four : Uniform bounds and logarithmic Sobolev inequality

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Two goals :

- A logarithmic Sobolev inequality for  $\bar{\rho}^N : \mathcal{H}_N(t) \leq C\mathcal{I}_N(t)$
- Uniform in time bounds on  $\|\nabla \bar{\rho}_t\|_{L^\infty}$  and  $\|\nabla^2 \bar{\rho}_t\|_{L^\infty}$

# A logarithmic Sobolev inequality

## Lemma (Tensorization)

*If  $\nu$  is a probability measure on  $\mathbb{T}^d$  satisfying a LSI with constant  $C_\nu^{LS}$ , then for all  $N \geq 0$ ,  $\nu^{\otimes N}$  satisfies a LSI with constant  $C_\nu^{LS}$*

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# A logarithmic Sobolev inequality

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## Lemma (Perturbation)

If  $\nu$  is a probability measure on  $\mathbb{T}^d$  satisfying a LSI with constant  $C_\nu^{LS}$ , and  $\mu$  is a probability measure with density  $h$  with respect to  $\nu$  such that, for some constant  $\lambda > 0$ ,  $\frac{1}{\lambda} \leq h \leq \lambda$ , then  $\mu$  satisfies a LSI with constant  $C_\mu^{LS} = \lambda^2 C_\nu^{LS}$ .

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## Lemma (LSI for the uniform distribution)

The uniform distribution  $u$  on  $\mathbb{T}^d$  satisfies a LSI with constant  $\frac{1}{8\pi^2}$ .

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## Lemma (LSI for the uniform distribution)

The uniform distribution  $u$  on  $\mathbb{T}^d$  satisfies a LSI with constant  $\frac{1}{8\pi^2}$ .

For all  $N \in \mathbb{N}$ ,  $t \geq 0$  and all probability density  $\mu_N \in C_{>0}^\infty(\mathbb{T}^{dN})$ ,

$$\mathcal{H}_N(\mu_N, \bar{\rho}_t^N) \leq \frac{\lambda^2}{8\pi^2} \frac{1}{N} \sum_{i=1}^N \int_{\mathbb{T}^d} \mu_N \left| \nabla_{x_i} \log \frac{\mu_N}{\bar{\rho}_t^N} \right|^2 d\mathbf{X}^N$$

# Uniform in time bounds on the derivatives

## Lemma

*For all  $n \geq 1$  and  $\alpha_1, \dots, \alpha_n \in \llbracket 1, d \rrbracket$ , there exist  $C_n^u, C_n^\infty > 0$  such that for all  $t \geq 0$ ,*

$$\|\partial_{\alpha_1, \dots, \alpha_n} \bar{\rho}_t\|_{L^\infty} \leq C_n^u \quad \text{and} \quad \int_0^t \|\partial_{\alpha_1, \dots, \alpha_n} \bar{\rho}_s\|_{L^\infty}^2 ds \leq C_n^\infty$$

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Thanks to Morrey's inequality and Sobolev embeddings, it is sufficient to prove such bounds in the Sobolev space  $H^m$  for all  $m$ , i.e in  $L^2$

# Uniform in time bounds on the derivatives-2

By induction on the order of the derivative

$$\frac{1}{2} \frac{d}{dt} \|\bar{\rho}_t\|_{L^2}^2 + \|\nabla \bar{\rho}_t\|_{L^2}^2 = 0,$$

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$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\partial_{\alpha_1, \alpha_2} \bar{\rho}_t\|_{L^2}^2 + \frac{1}{2} \sum_{\alpha_3} \|\partial_{\alpha_1, \alpha_2, \alpha_3} \bar{\rho}_t\|_{L^2}^2 &\leq \|V\|_{L^\infty}^2 \|\partial_{\alpha_1} \nabla \bar{\rho}_t\|_{L^2}^2 \|\nabla \bar{\rho}_t\|_{L^2}^2 \\ &\quad + \|K\|_{L^1}^2 \|\bar{\rho}_t\|_{L^\infty}^2 \|\partial_{\alpha_1} \nabla \bar{\rho}_t\|_{L^2}^2, \end{aligned}$$

# Uniform in time bounds on the derivatives-2

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etc

# Assumptions ?

$$\text{Goal : } K(x) = \frac{1}{2\pi} \frac{x^\perp}{|x|^2} = \frac{1}{2\pi} \left( -\frac{x_2}{|x|^2}, \frac{x_1}{|x|^2} \right)$$

## Justifying the calculations

- There is  $\lambda > 1$  such that  $\bar{\rho}_0 \in C_\lambda^\infty(\mathbb{T}^d)$   
 $\implies \bar{\rho} \in C_\lambda^\infty(\mathbb{R}^+ \times \mathbb{T}^d)$  (Ben-Artzi '94)
- $\rho^N \in C_\lambda^\infty(\mathbb{R}^+ \times \mathbb{T}^{Nd})$  (???)

## Dealing with the terms

- In the sense of distributions,  $\nabla \cdot K = 0$ .
- There is a matrix field  $V \in L^\infty$  such that  $K = \nabla \cdot V$ , i.e for  $1 \leq \alpha \leq d$ ,  $K_\alpha = \sum_{\beta=1}^d \partial_\beta V_{\alpha,\beta}$  (Phuc-Torres '08).

## Uniformity in time

- For all  $n \geq 1$ ,  $C_n^0 := \|\nabla^n \bar{\rho}_0\|_{L^\infty} < \infty$
- $\|K\|_{L^1} < \infty$  (also used to show regularity).

## Step five : Conclusion

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On the assumptions

There are constants  $C_1, C_2^\infty, C_3 > 0$  and a function  $t \mapsto C_2(t) > 0$  with  $\int_0^t C_2(s) ds \leq C_2^\infty$  for all  $t \geq 0$  such that for all  $t \geq 0$

$$\frac{d}{dt} \mathcal{H}_N(t) \leq -(C_1 - C_2(t)) \mathcal{H}_N(t) + \frac{C_3}{N}.$$

Multiplying by  $\exp(C_1 t - \int_0^t C_2(s) ds)$  and integrating in time we get

$$\begin{aligned} \mathcal{H}_N(t) &\leq e^{-C_1 t + \int_0^t C_2(s) ds} \mathcal{H}_N(0) + \frac{C_3}{N} \int_0^t e^{C_1(s-t) + \int_s^t C_2(u) du} ds \\ &\leq e^{C_2^\infty - C_1 t} \mathcal{H}_N(t) + \frac{C_3}{C_1 N} e^{C_2^\infty}, \end{aligned}$$

which concludes.

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On the assumptions

$$\text{On } \rho^N \in \mathcal{C}_\lambda^\infty(\mathbb{R}^+ \times \mathbb{T}^{Nd})$$

Everything works for regularized kernels  $K^\epsilon$ , and the final result is independent of  $\epsilon$ .

# Assumptions

## On the initial condition

- There is  $\lambda > 1$  such that  $\bar{\rho}_0 \in \mathcal{C}_\lambda^\infty(\mathbb{T}^d)$
- For all  $n \geq 1$ ,  $C_n^0 := \|\nabla^n \bar{\rho}_0\|_{L^\infty} < \infty$

## On the potential $K$

- $\|K\|_{L^1} < \infty$ .
- In the sense of distributions,  $\nabla \cdot K = 0$ ,
- There is a matrix field  $V \in L^\infty$  such that  $K = \nabla \cdot V$ , i.e for  $1 \leq \alpha \leq d$ ,  $K_\alpha = \sum_{\beta=1}^d \partial_\beta V_{\alpha,\beta}$ .

