

Inclined co-orbitals in the three-body problem

From Marchal's family to P_{12}

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Introduction

The three-body problem: generalities

- Problem definition: 3 masses and gravity

$$\mathcal{H}(\mathbf{u}, \tilde{\mathbf{u}}) = \underbrace{\frac{\tilde{\mathbf{u}}_0^2}{2M} + \sum_{i=1}^2 \left(\frac{m_0 + m_i}{2m_0 m_i} \tilde{\mathbf{u}}_i^2 - \mathcal{G} \frac{m_0 m_i}{|\mathbf{u}_i - \mathbf{u}_0|} \right)}_{\mathcal{H}_0} + \underbrace{\frac{\tilde{\mathbf{u}}_1 \cdot \tilde{\mathbf{u}}_2}{m_0}}_{T_1} - \underbrace{\mathcal{G} \frac{m_1 m_2}{|\mathbf{u}_2 - \mathbf{u}_1|}}_{U_1}$$

$\underbrace{\hspace{15em}}_{\mathcal{H}_1}$

The three-body problem: generalities

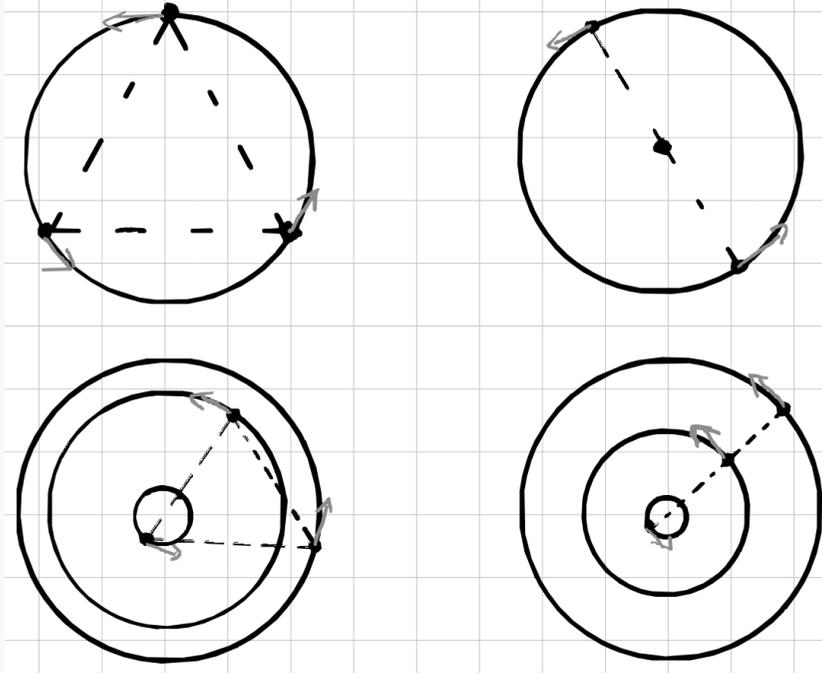


Figure 1: Lagrange (left) and Euler (right) periodic orbits.

The three-body problem: generalities

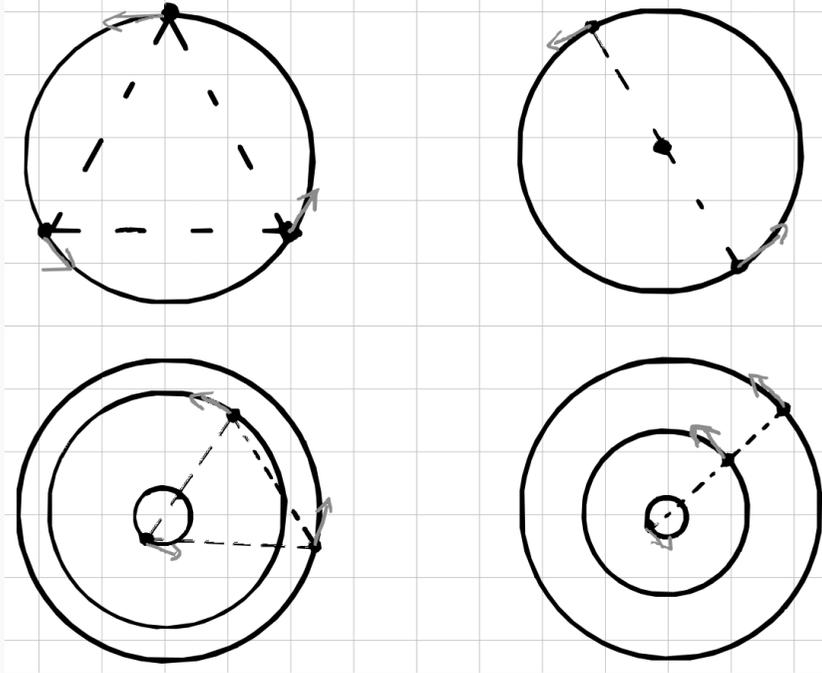
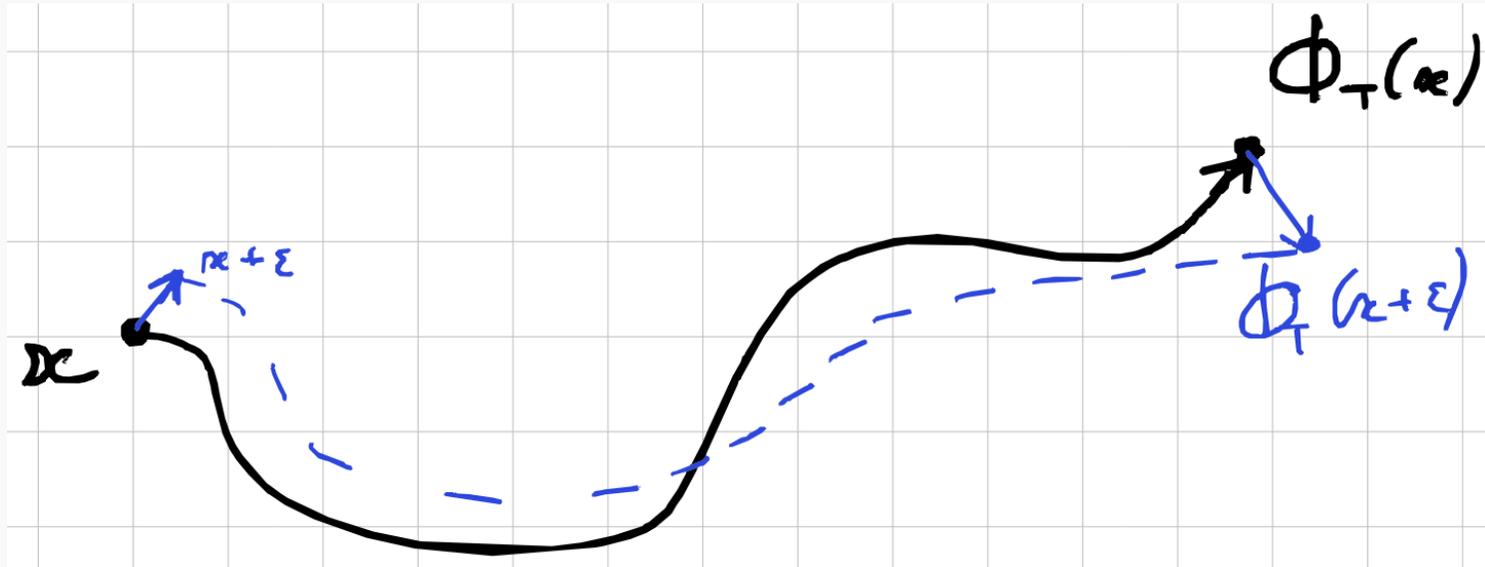


Figure 2: Lagrange (left) and Euler (right) periodic orbits.

Gascheau's criterium:

$$\frac{m_0 m_1 + m_0 m_2 + m_1 m_2}{(m_0 + m_1 + m_2)^2} > \frac{1}{27}$$

Some concepts



- \mathbf{x} = (positions, momenta)
- $\Phi_T : \mathbf{x} \mapsto$ evolution after a time T
- Eigenvalues of $D\Phi_T$: describe how $\Phi_T(\mathbf{x} + \boldsymbol{\epsilon})$ will behave

The eigenvalues of Lagrange for the planetary problem

Family	Ending
Homographics	Triple collisions

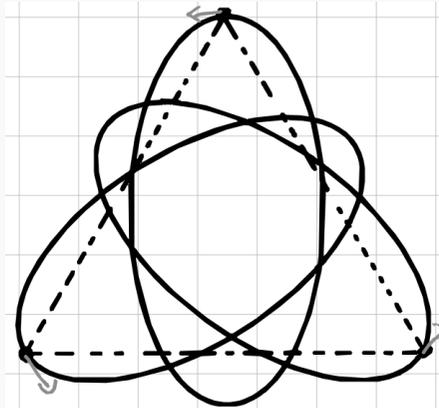


Figure 4: Homographic, anti-Lagrange, and tadpole/horseshoe orbits. Cf. (Robutel and Pousse, 2013).

The eigenvalues of Lagrange for the planetary problem

Family	Ending
Homographics	Triple collisions
Anti-Lagrange	Triple collisions (unproven)

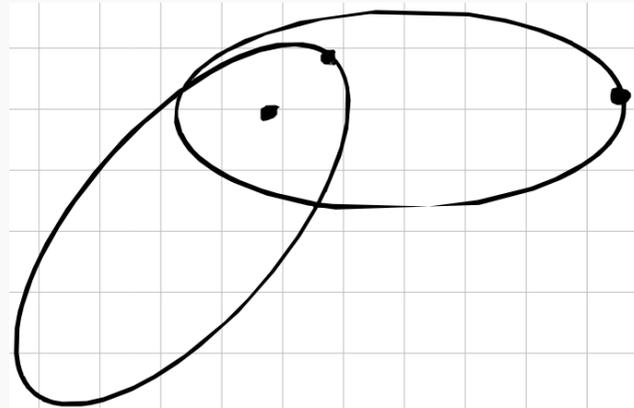
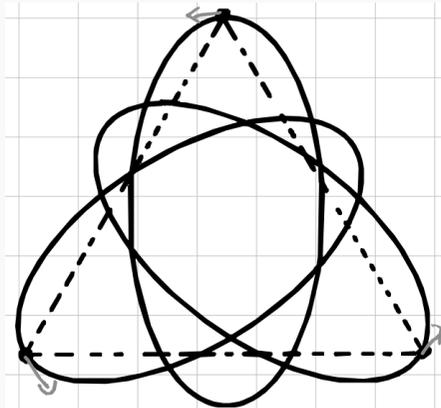


Figure 5: Homographic, anti-Lagrange, and tadpole/horseshoe orbits. Cf. (Robutel and Pousse, 2013).

The eigenvalues of Lagrange for the planetary problem

Family	Ending
Homographics	Triple collisions
Anti-Lagrange	Triple collisions (unproven)
Tadpole	Horseshoe orbits

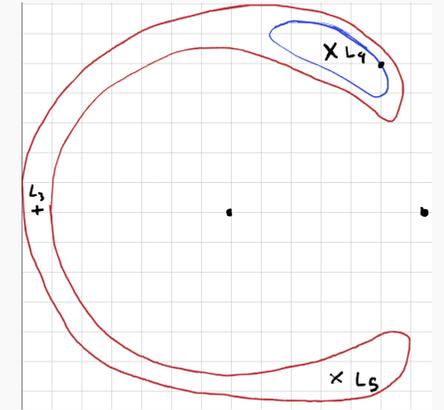
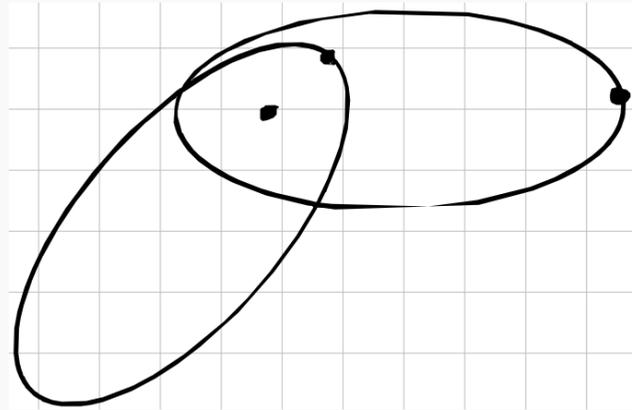
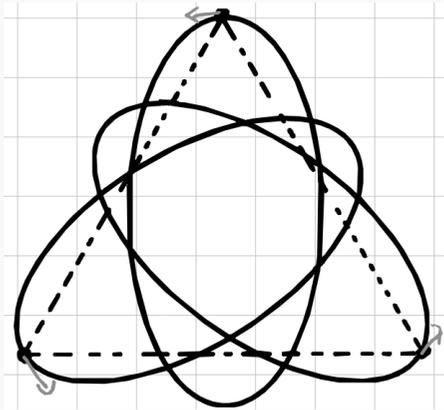


Figure 6: Homographic, anti-Lagrange, and tadpole/horseshoe orbits. Cf. (Robutel and Pousse, 2013).

The eigenvalues of Lagrange for the planetary problem

Family	Ending
Homographics	Triple collisions
Anti-Lagrange	Triple collisions (unproven)
Tadpole	Horseshoe orbits
Vertical	Unknown (depending on the masses)

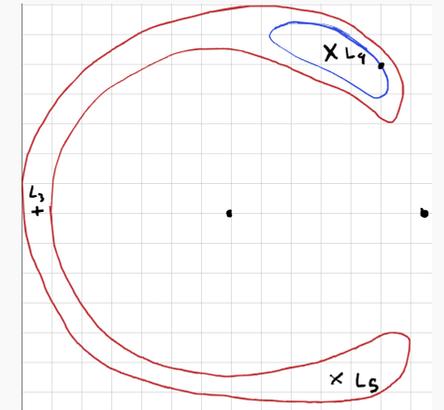
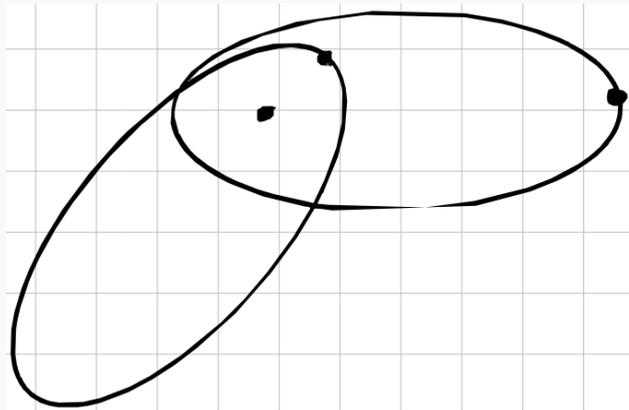
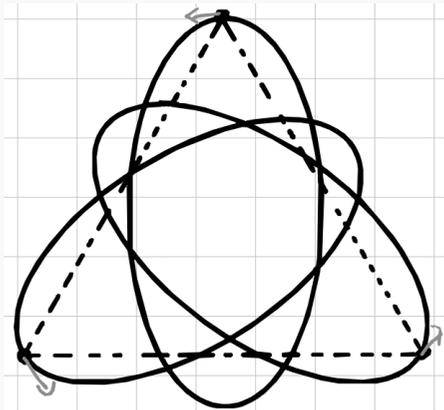


Figure 7: Homographic, anti-Lagrange, and tadpole/horseshoe orbits. Cf. (Robutel and Pousse, 2013).

The vertical family – equal masses: P_{12}

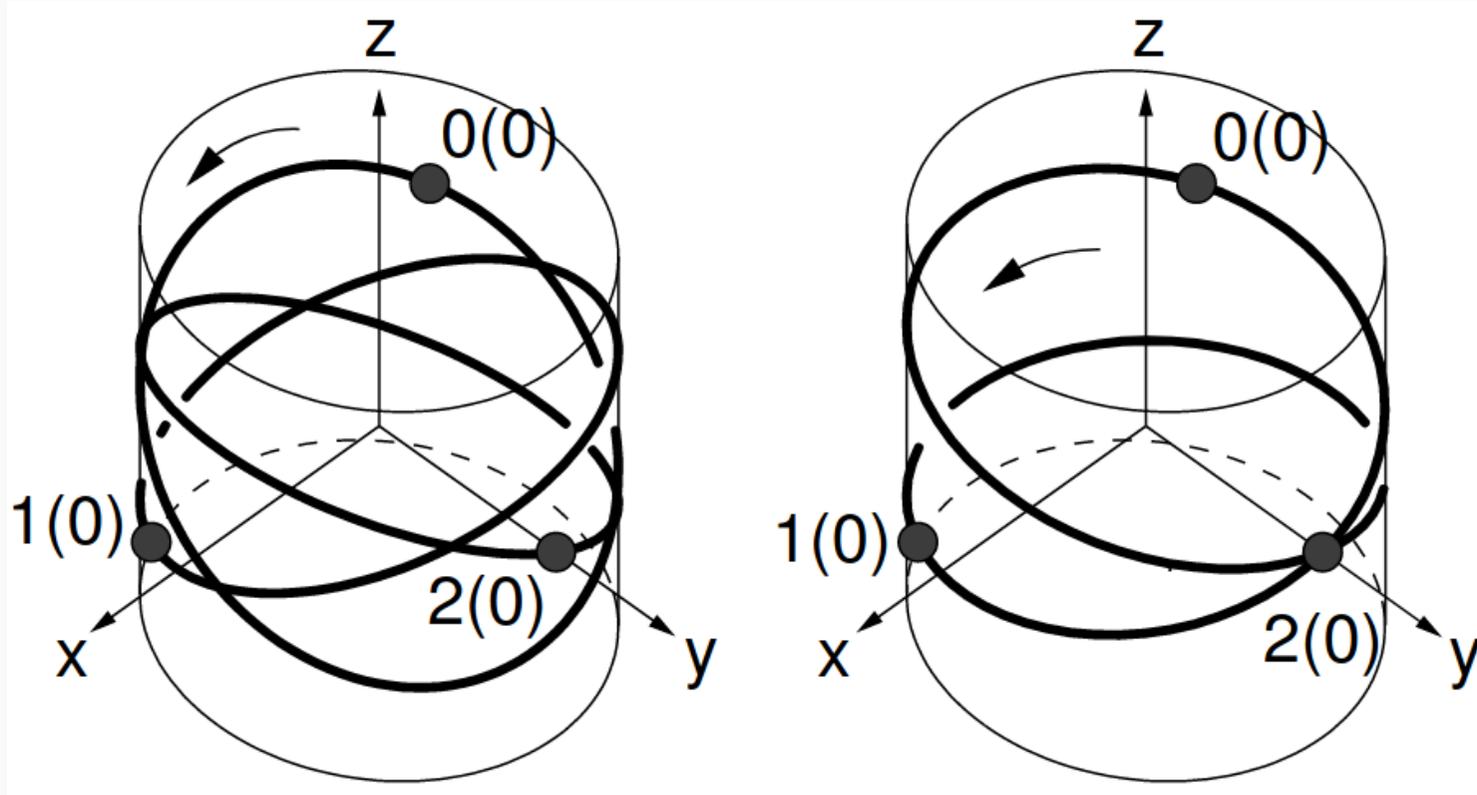


Figure 8: Figure 1 of (Chenciner and Féjóz, 2008). Also studied in (Chenciner, Féjóz and Montgomery, 2005; Chenciner and Féjóz, 2008; Calleja *et al.*, 2024)...

The vertical family – restricted problem: Marchal 2009

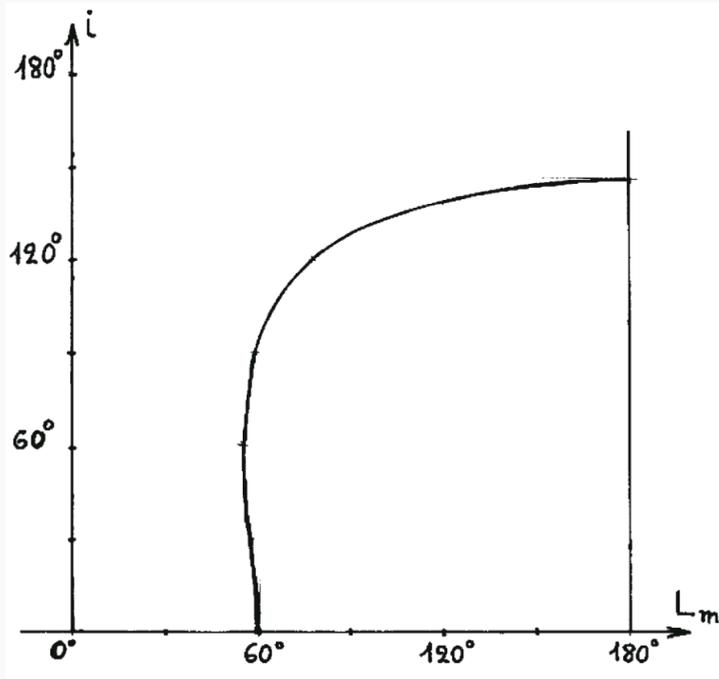


Figure 9: Fig. 3 of (Marchal, 2009) showing the vertical family in the restricted problem, with masses $(1 - \varepsilon, \varepsilon, 0)$.

The vertical family – restricted problem: Marchal 2009

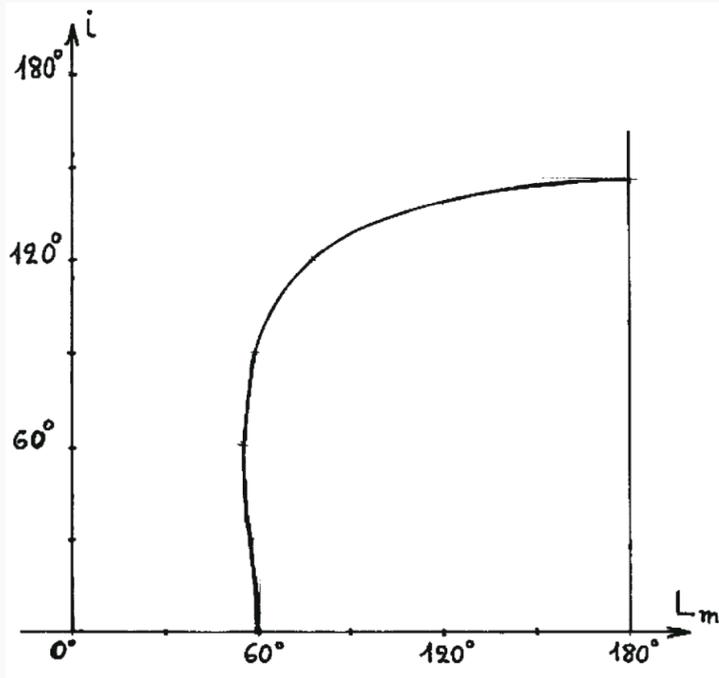


Figure 10: Fig. 3 of (Marchal, 2009) showing the vertical family in the restricted problem, with masses $(1 - \varepsilon, \varepsilon, 0)$.

“These periodic orbits belong to a three parameter family of similar periodic orbits of the general three body problem, two parameters being the mass ratios of the three bodies and the third a generalization of the inclination i . [...] The one parameter family P_{12} of Marchal (2000) is another limit case of this three parameter family: the case of three equal masses.”

Agenda

1. Introduction
2. How to chase a quasi-periodic solution
3. Marchal's family for small masses
4. From Marchal to P_{12} : varying the masses
5. Further work

How to chase a quasi-periodic solution

Rephrasing the problem

**18
dims**

Start: 18 dimensional-problem, 3 positions+velocities for 3 bodies

Rephrasing the problem

**12
dims**

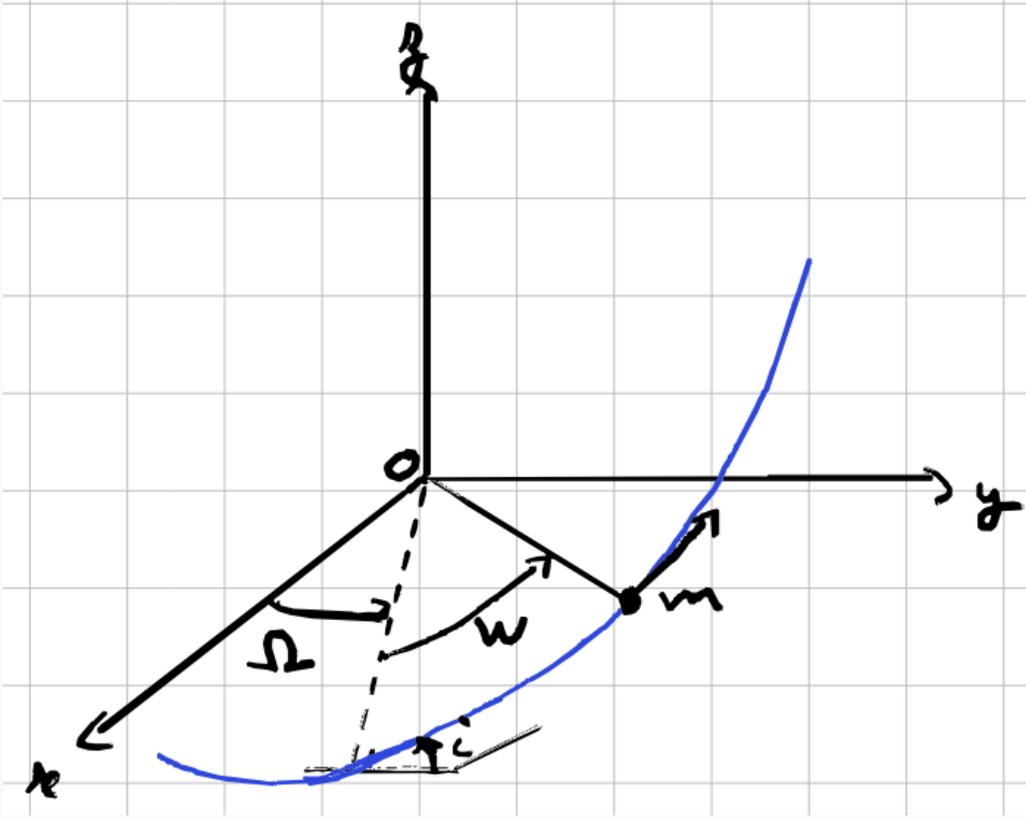
Start: 18 dimensional-problem, 3 positions+velocities for 3 bodies

Fix one body at the origin: “lose one body” → 12 dimensions.

- Canonical heliocentric coordinates
 - Heliocentric positions, barycentric momenta

Rephrasing the problem

**12
dims**



Hill coordinates (Laskar, 2017)

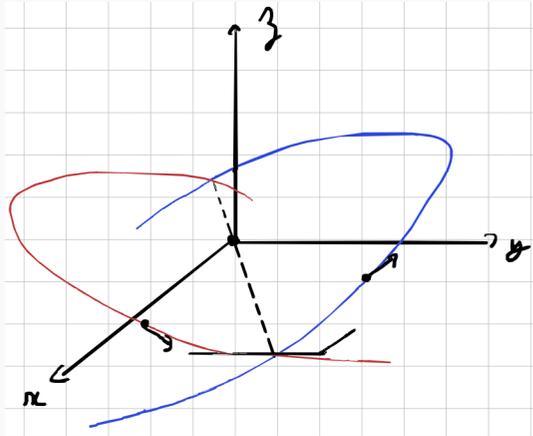
- r_i radii, \dot{r}_i radial velocities
- w_i true anomalies, G_i angular momenta
- Ω_i node longitudes, $G_i \cos(i_i)$ vertical angular momenta
- C total angular momentum

Reduction of the nodes

$$\mathcal{H}\left(\Omega_1 - \Omega_2, \cancel{\Omega_1 + \Omega_2}, \underbrace{i_1 + i_2}_J, \cancel{i_1 - i_2} \dots\right)$$

**12
dims**

Reduction of the nodes



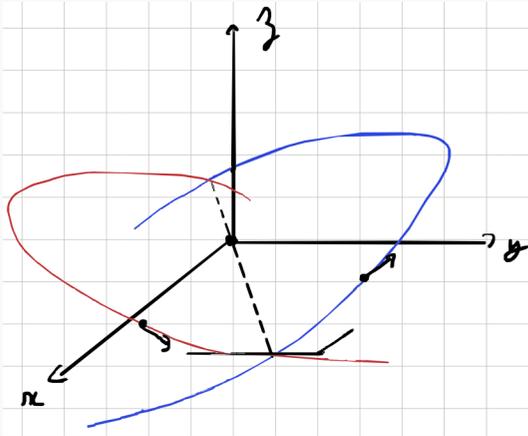
$$\mathcal{H}(\Omega_1 - \Omega_2, \cancel{\Omega_1 + \Omega_2}, \underbrace{i_1 + i_2}_J, \cancel{i_1 - i_2} \dots)$$

**10
dims**

Reduction of the nodes

9 dims

$$\mathcal{H}(\Omega_1 - \Omega_2, \cancel{\Omega_1 + \Omega_2}, \underbrace{i_1 + i_2}_J, \cancel{i_1 - i_2} \dots)$$



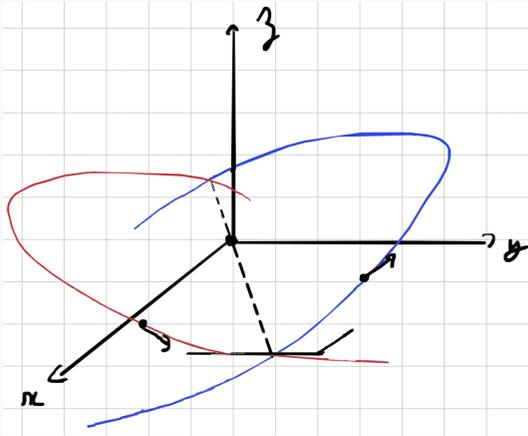
But $\Omega_1 - \Omega_2 = \pi$

We can gain 3 dimensions ! (Robutel, 1995; Mastroianni and Efthymiopoulos, 2019)

Reduction of the nodes

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But $\Omega_1 - \Omega_2 = \pi$

We can gain 3 dimensions ! (Robutel, 1995; Mastroianni and Efthymiopoulos, 2019)

Draconitic frame

- 9 dimensions : $(r_i, w_i, \tilde{r}_i, G_i)_{i=1,2}, C$

The numerical method

The problem :

- Fixed : masses, period (space-time invariance)
- 9-dimensional search space : 8 coordinates and C
- We are looking for **periodic orbits**... in a rotating frame

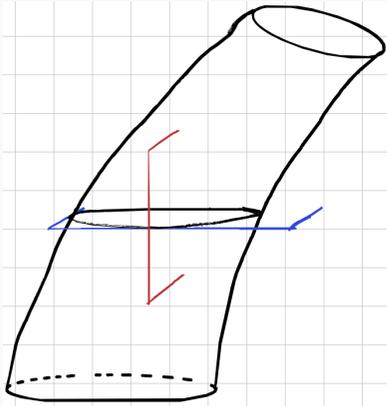
$$\Phi_T(\mathbf{x}) - \mathbf{x} = \mathbf{0} \rightarrow \text{Newton-Raphson algorithm}$$

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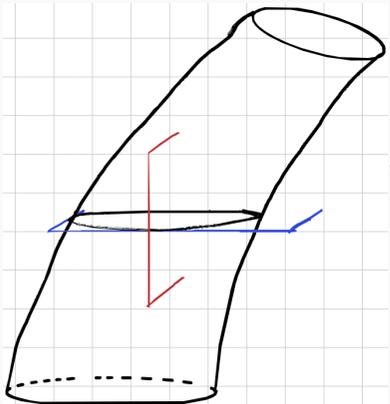
- Lack of local unicity: add sections
 - One for the family, one for the orbit

The numerical method

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- Fixed : masses, period (space-time invariance)
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$$\begin{pmatrix} \Phi_T(\mathbf{x}) - \mathbf{x} \\ \sigma(\mathbf{x}) \end{pmatrix} = \mathbf{0} \rightarrow \text{Newton-Raphson algorithm}$$



- Lack of local unicity: add sections
 - One for the family, one for the orbit

Marchal's family for small masses

General properties

- Two small masses (10^{-3}) and one big mass ($1 - 2 * 10^{-3}$)

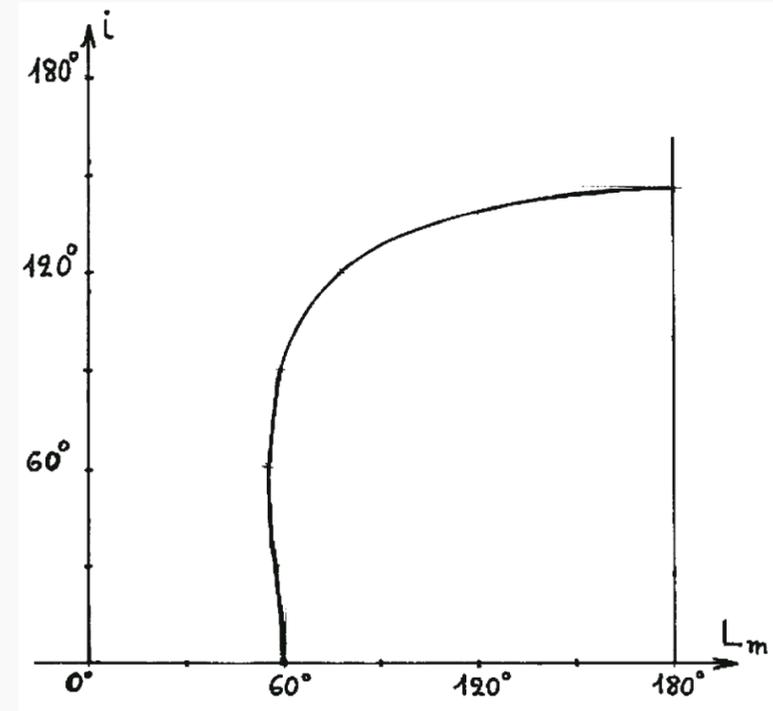
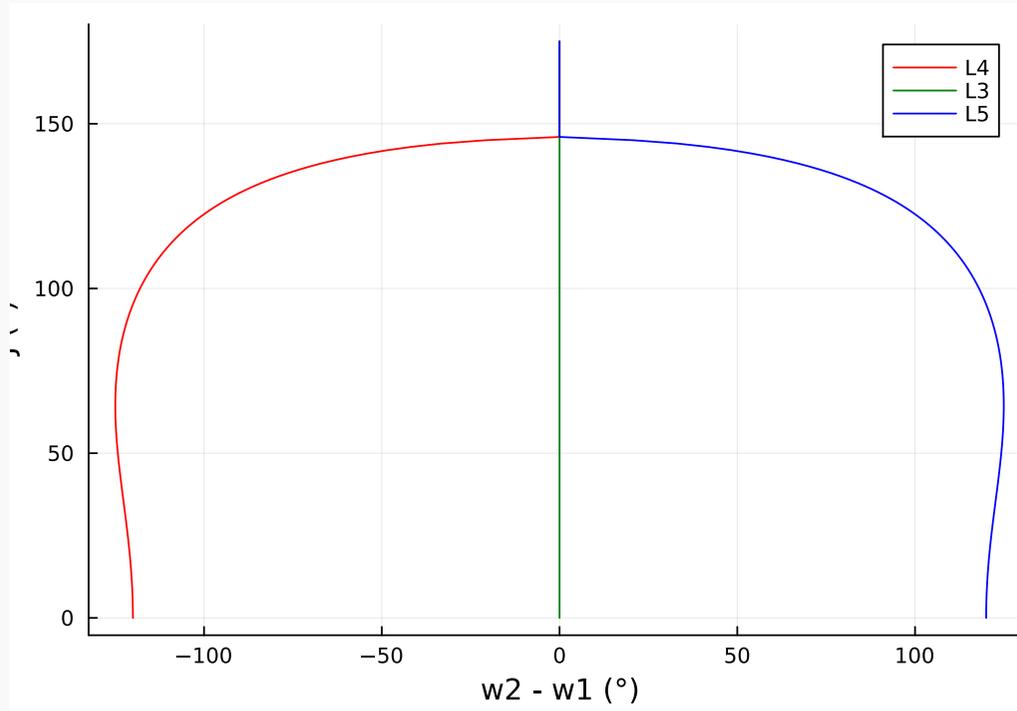


Figure 19: Lagrange and Euler families

Stability

- Linearly stable until $\sim 60^\circ$

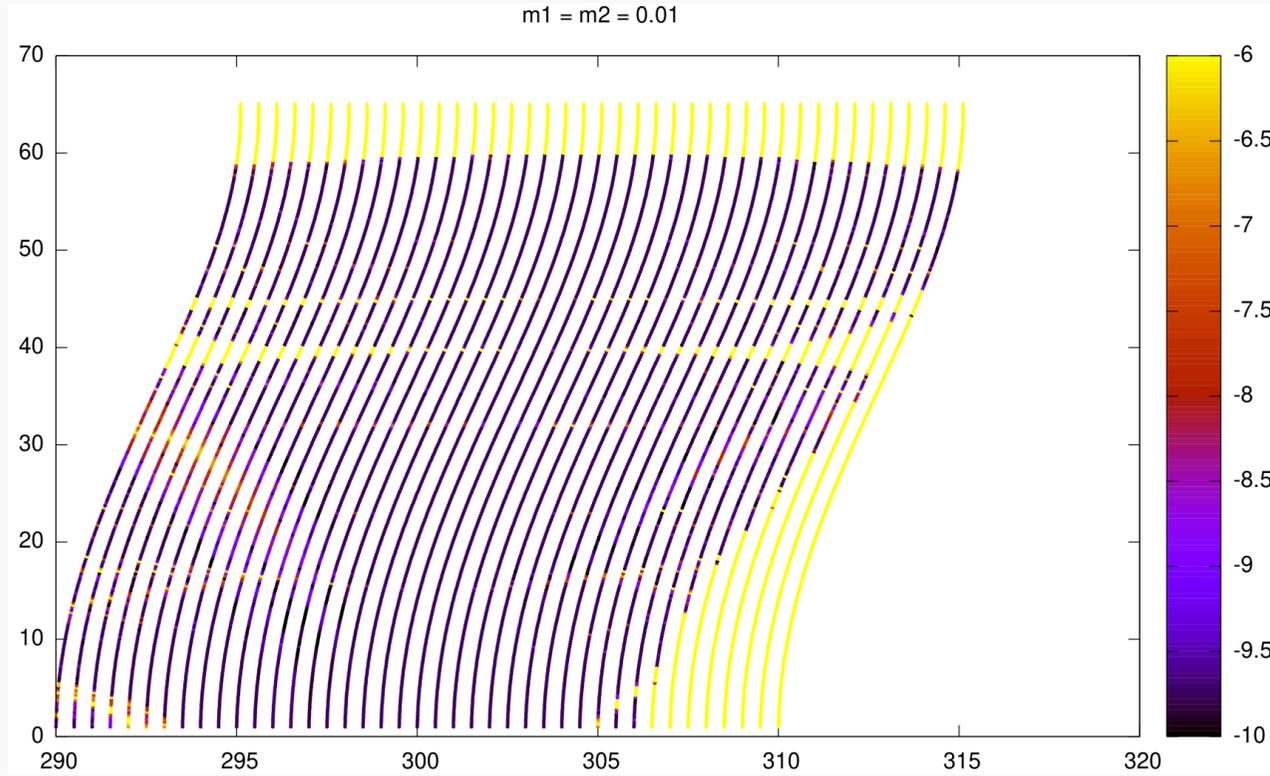


Figure 20: Full stability around the limit for $\varepsilon_m = 0.01$.

$$\Phi_{\frac{T}{2}}(\mathbf{x}) = -\mathbf{x}$$

- Italian symetry
 - Cut integration time in half
 - Eliminate parasitic families

Symetries

$$\Phi_{\frac{T}{2}}(\mathbf{x}) = -\mathbf{x}$$

- Italian symetry
 - Cut integration time in half
 - Eliminate parasitic families

$$R_y(\pi)\mathbf{x} = \sigma_{1\leftrightarrow 2}(\mathbf{x})$$

- Second symetry
 - Fixes position on the orbit
 - Adds three more sections.

From Marchal to P_{12} : varying
the masses

Stability: over Gascheau's limit

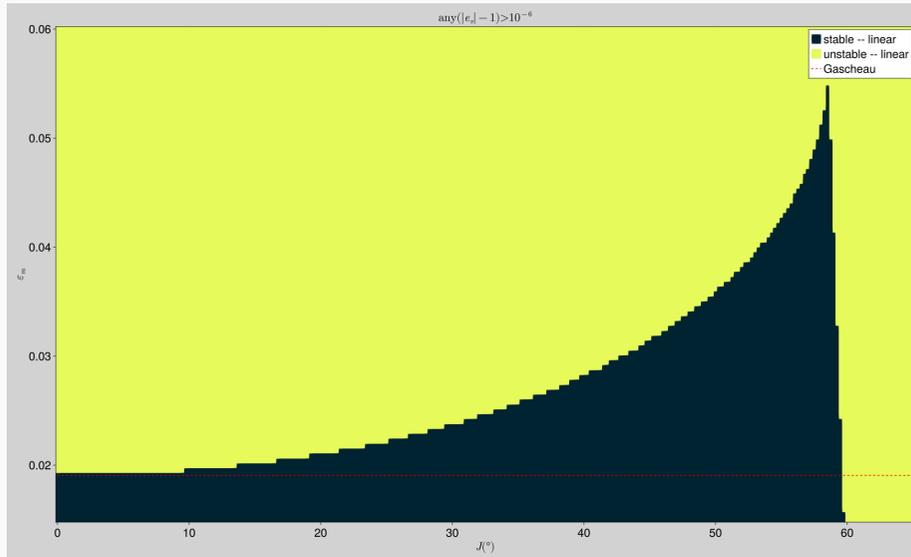


Figure 21: Linear and full stability along the family for various masses.

Stability: over Gascheau's limit

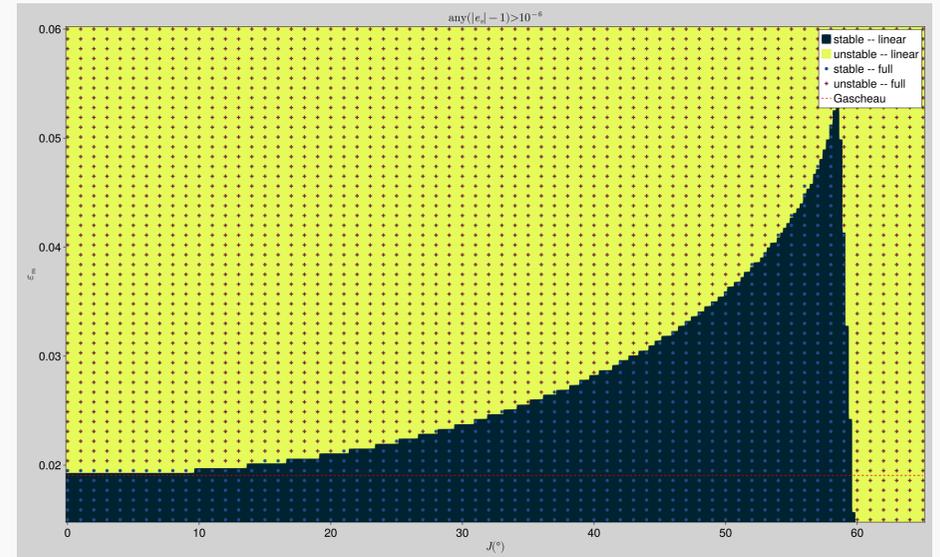
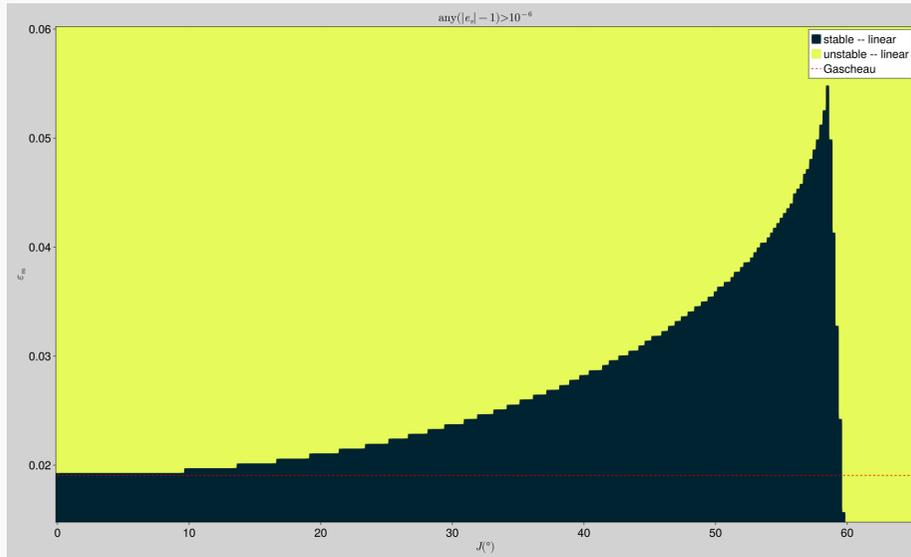


Figure 22: Linear and full stability along the family for various masses.

Stability: over Gascheau's limit

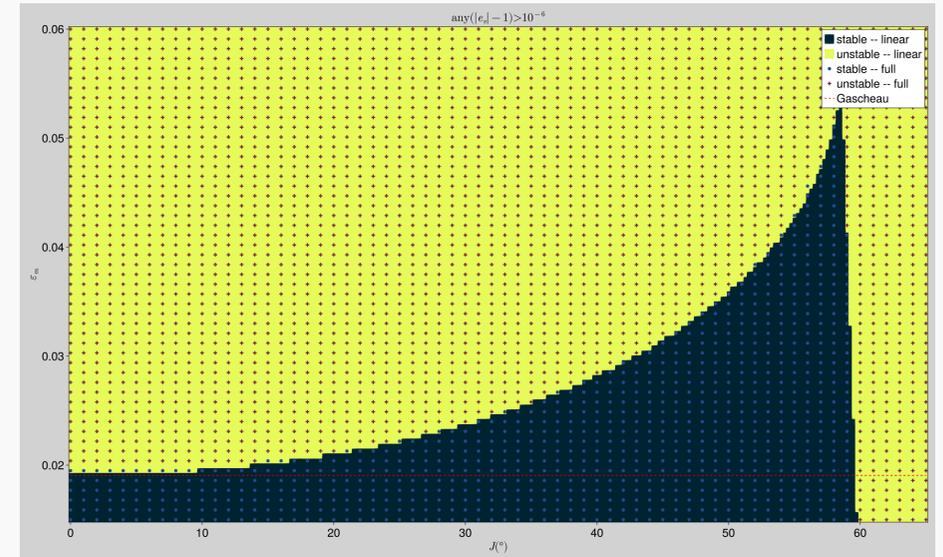
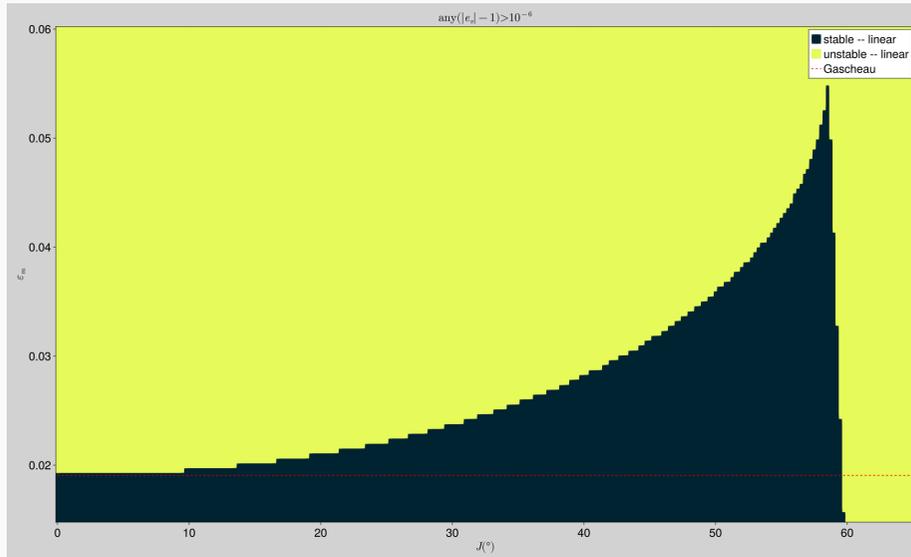


Figure 23: Linear and full stability along the family for various masses.

Similar results obtained for elliptic Lagrange (Roberts, 2002)

Approaching the equal masses

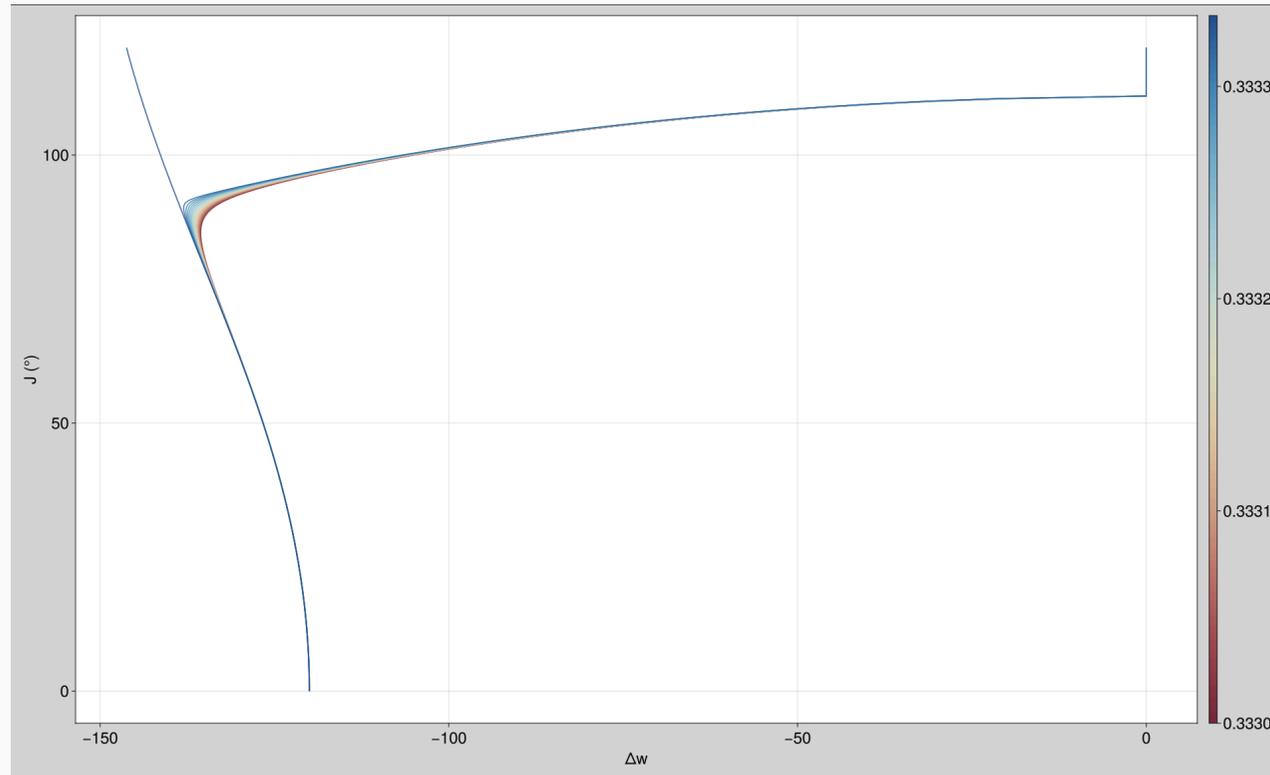


Figure 24: Start of Marchal family for $\varepsilon_m \in \left[0.3330, \frac{1}{3}\right]$ in the $(\Delta w, J)$ plane.

Approaching the equal masses

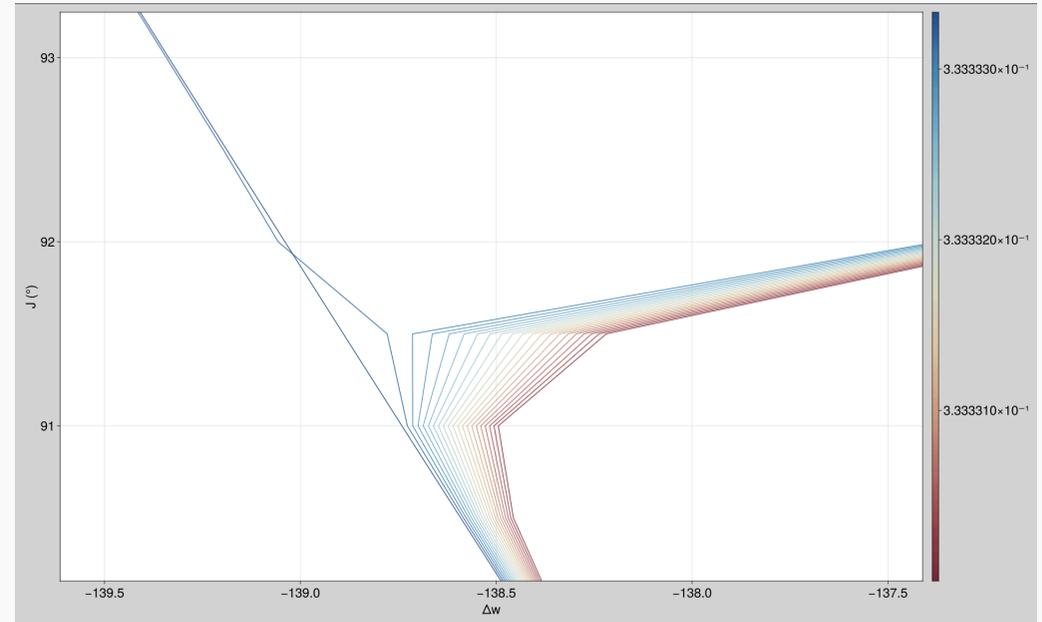
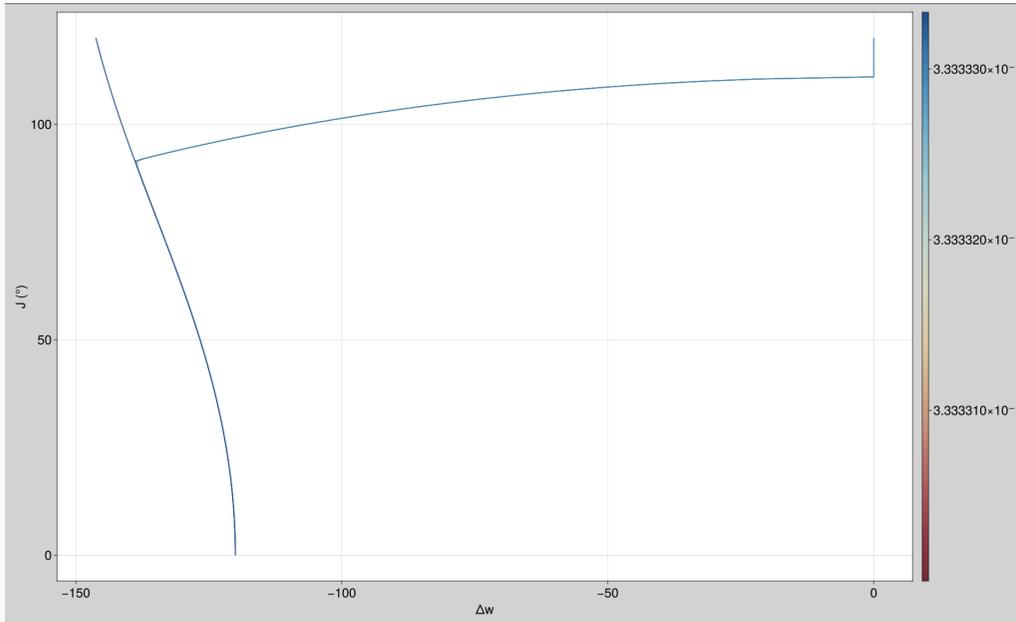


Figure 25: Zooming on the bifurcation for $\varepsilon_m \in \left[0.333330; \frac{1}{3}\right]$.

P_{12} : the bifurcations

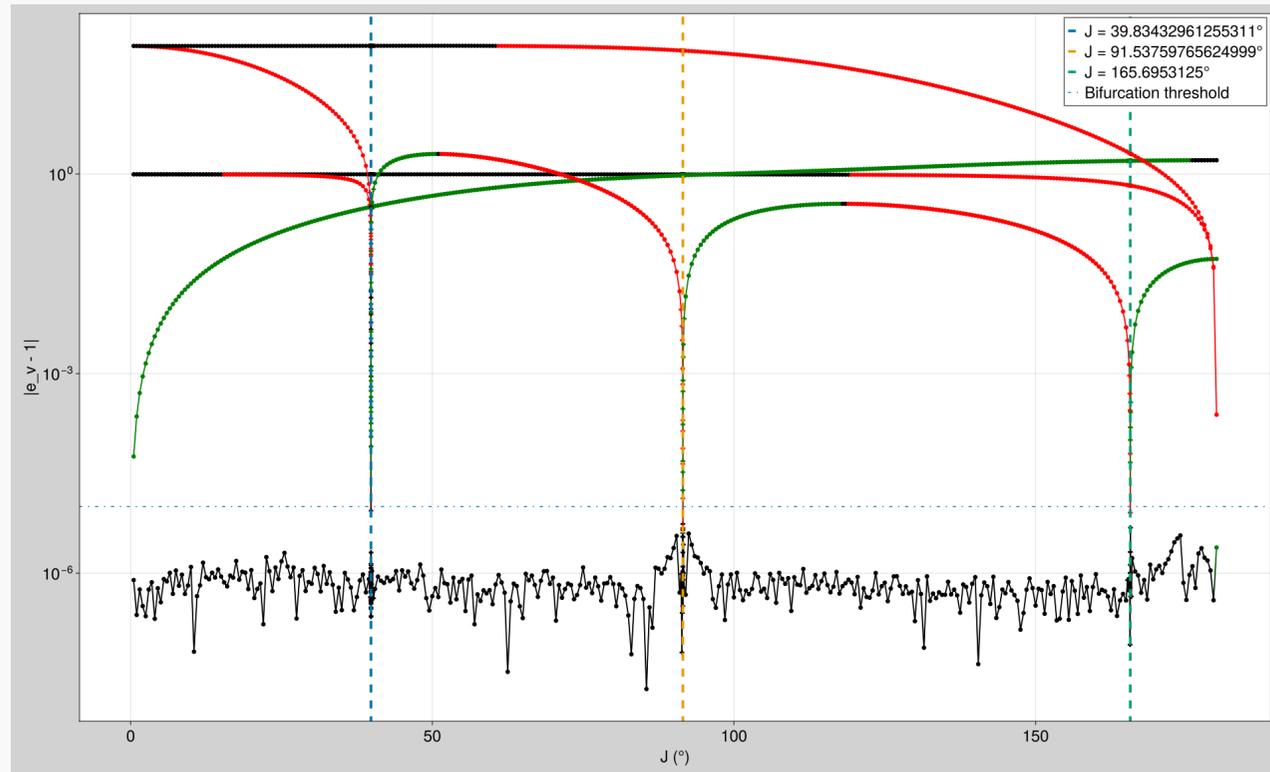


Figure 26: Distance to 1 of monodromy matrix eigenvalues along P_{12} .

Marchal and P_{12}

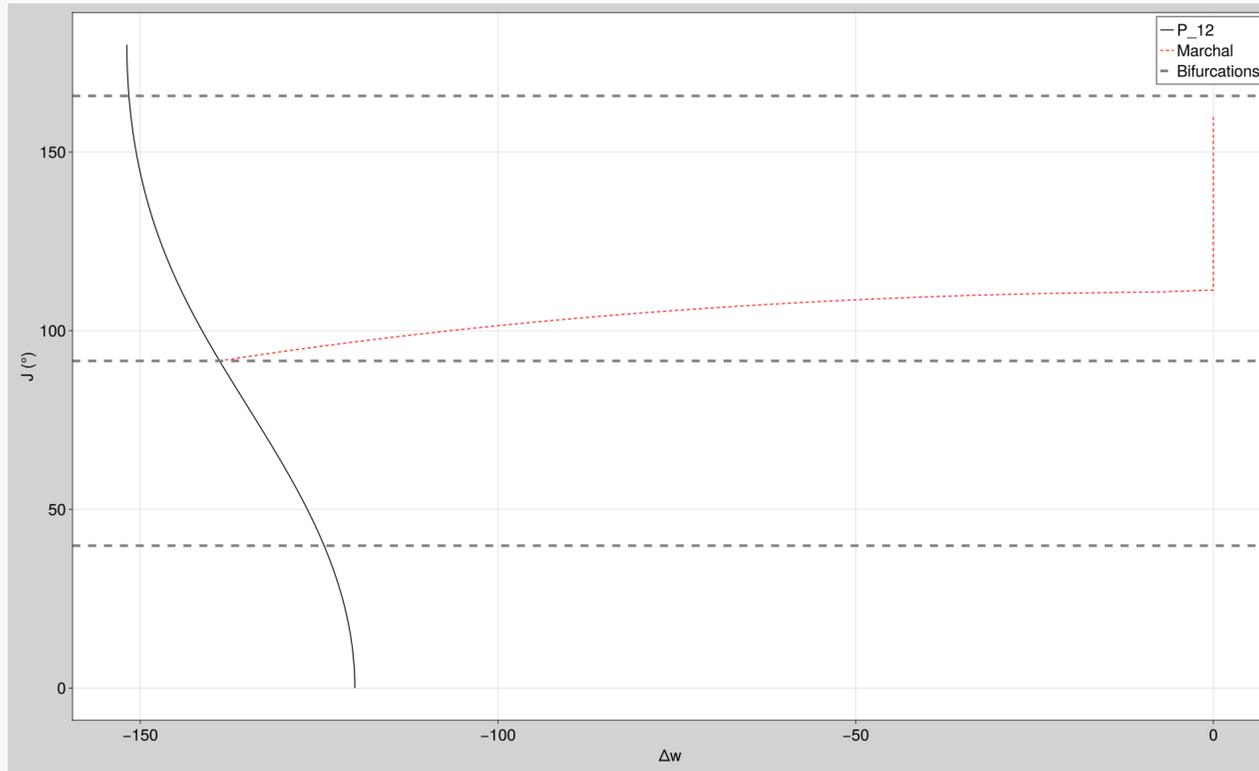


Figure 27: Marchal and P_{12} for equal masses.

P_{12} : the bifurcations



Figure 28: The first bifurcation was followed by (Calleja *et al.*, 2021) to a rotating Eight. The second is the separation with Marchal's family. The third one appears to be new.

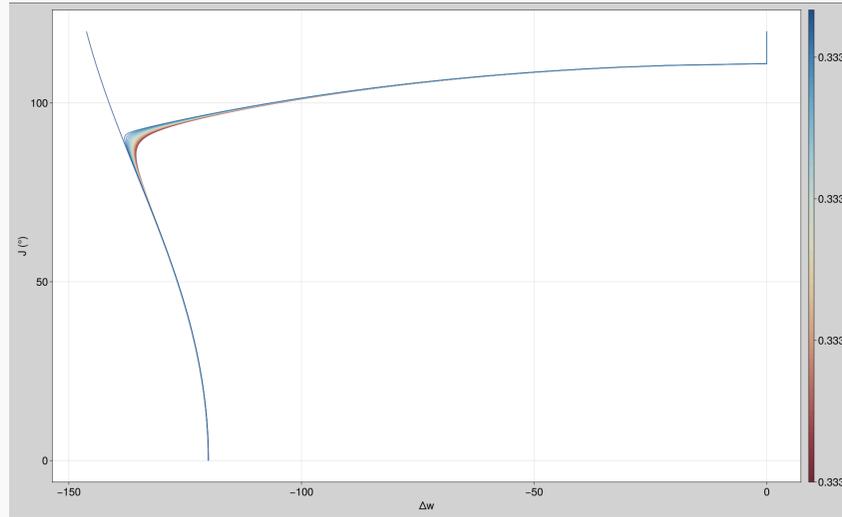
Further work

Digging into the bifurcations



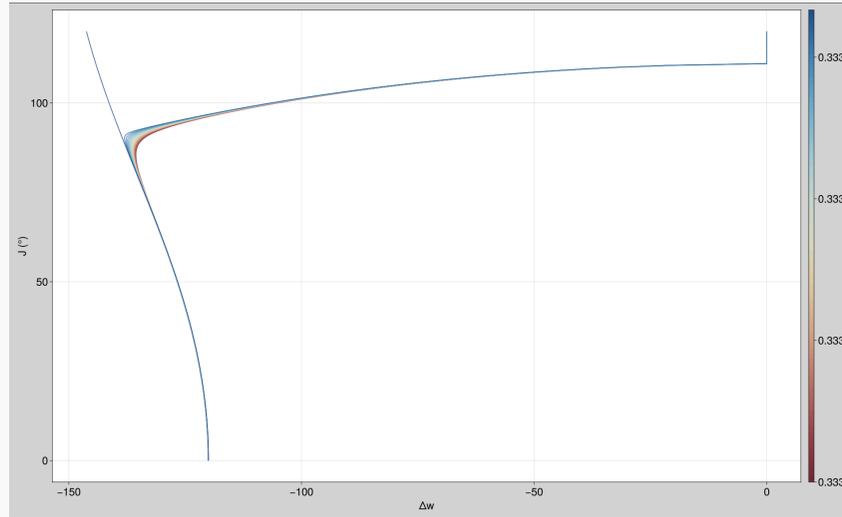
- Understand which bifurcations are crossings with families
- Which of these preserve symmetries

The P_{12} branch for unequal masses



Does it exist for unequal masses? Where does it go?

The P_{12} branch for unequal masses

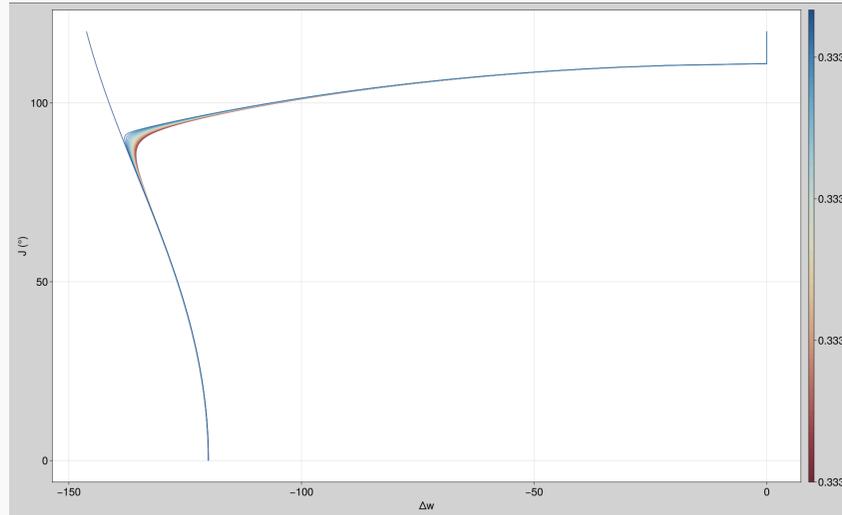


Does it exist for unequal masses? Where does it go?

(Doedel *et al.*, 2003) continues the Eight for masses $(1 - 2m, m, m)$, $m > \frac{1}{3}$

- Wild conjecture: could the P_{12} branch end at the Doedel family?

The P_{12} branch for unequal masses



Does it exist for unequal masses? Where does it go?

(Doedel *et al.*, 2003) continues the Eight for masses $(1 - 2m, m, m)$, $m > \frac{1}{3}$

- Wild conjecture: could the P_{12} branch end at the Doedel family?

What happens to the other bifurcations at unequal masses?

The Euler family at equal masses

- Currently failing (probably for numerical reasons, not mathematical ones)

Conclusion

We...

Confirmed part of Marchal's conjecture by linking Marchal's family in the restricted case to the P_{12} family in the equal masses case

- ... Up to a point: bifurcation

Conclusion

We...

Confirmed part of Marchal's conjecture by linking Marchal's family in the restricted case to the P_{12} family in the equal masses case

- ... Up to a point: bifurcation

Developed numerical tools to find periodic orbits

- Not necessarily in the three-body problem!

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