Real Geometric Quantization

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ICMAT, CSIC

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Francisco Presas Real Geometric Quantization

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- Classical example
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Bertram Kostant.



Nick Woodhouse.



Jędrzej Šniatycki.



Eva Miranda.



Mark Hamilton.



Romero Solha.

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 Phase space: Symplectic manifold (M, ω).



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Not unique: Planck's constant \hbar parameter.

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Geometric Quantization

 (M,ω)

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Geometric Quantization

(*M*,

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Hilbert space: $L^2(\Gamma(L)).$

•
$$k = \frac{1}{\hbar} \in \mathbb{Z}_+$$
. Different quantizations given by $L^{\otimes k}$,
curv $(L^{\otimes k}) = -ik\omega$.

• Space too big; i.e. $(T^*N, \omega = d\lambda_{\text{Liouville}}), L = \mathbb{C}$. We obtain $L^2\Gamma(T^*N; \mathbb{C})$, and however we expect $L^2\Gamma(N; \mathbb{C})$.

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Geometric Quantization

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- Space too big; i.e. (*T***N*, ω = dλ_{Liouville}), *L* = C. We obtain L²Γ(*T***N*; C), and however we expect L²Γ(*N*; C).
- Polarizations: A way of reducing the dimension of the Hilbert space. P ⊂ (T^CM = TM ⊗ C, ω_C) regular n-dimensional complex Lagrangian integrable distribution.
- $\mathcal{P}^{\mathbb{R}} \subset TM$ Lagrangian foliation $\Rightarrow \mathcal{P}^{\mathbb{R}} \otimes \mathbb{C}$ is a polarization. $M = T^*N$ then $\mathcal{P}^{\mathbb{R}} = \ker d\pi, \ \pi : T^*N \to N$ projection.

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 - $M = T^*N, \mathcal{P}^{\mathbb{R}} = \ker d\pi$ then $H^0(M; \mathcal{J}) \equiv \{\text{sections constant along leaves}\}$. We obtain $\mathcal{H} = L^2\Gamma(N)$.

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2 \mathcal{P}^J then $H^0(M; \mathcal{J}) \equiv \{\text{Holomorphic sections}\}$.

• *H* Hamiltonian, $s \in \mathcal{H}$.

$$\hat{H}s = -i\hbar \nabla_{X_H}s + Hs.$$

A la Kostant Quantization

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Theorem (Šniatycki)

$$(f_1, \ldots, f_n) : (M^{2n}, \omega) \to B^n \subset \mathbb{R}^n$$
 integrable system, then $L^2 H^0(M; \mathcal{J}) = \mathcal{Q}(M).$

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Computation Kit

 $0 \to \mathcal{J} \to \mathcal{S}^0_{\mathcal{P}} \xrightarrow{\nabla} \mathcal{S}^1_{\mathcal{P}} \xrightarrow{d^{\nabla}} \mathcal{S}^2_{\mathcal{P}} \xrightarrow{d^{\nabla}} \cdots \text{ exact fine resolution.}$

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- **3** $H^{j}(M; \mathcal{J}) = 0$ for j > n.

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- Exactness: Poincaré Lemma works for generic singular polarizations (Miranda, Solha).
- $I = H^{j}(M; \mathcal{J}) = 0 \text{ for } j > n.$
- Mayer-Vietoris works $M = M_1 \cup M_2$ (Miranda, P, Solha)

$$\cdots \to H^{j}(M; \mathcal{J}) \to H^{j}(M_{1}; \mathcal{J}) \oplus H^{j}(M_{2}; \mathcal{J}) \to$$
$$H^{j}(M_{1} \cap M_{2}; \mathcal{J}) \to H^{j+1}(M; \mathcal{J}) \to \cdots$$

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Künneth (Miranda, P):

$$H^{j}(M_{1} \times M_{2}; \mathcal{J}_{1} \boxtimes \mathcal{J}_{2}) \cong \bigoplus_{j=p+q} H^{p}(M_{1}; \mathcal{J}_{1}) \oplus H^{q}(M_{2}; \mathcal{J}_{2})$$

whenever M_1 admits a good cover, the geometric quantization associated to (M_2, \mathcal{J}_2) has finite dimension and M_2 is a submanifold of a compact manifold.

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Regular integrable system

$$\begin{split} I_j &= (-\varepsilon, \varepsilon), j = 1, 2.\\ \text{Computation 1: } \mathcal{Q}(I_1 \times I_2, \omega = \mathsf{d} x_1 \wedge \mathsf{d} x_2; \mathcal{P} = \partial_{x_2}). \end{split}$$

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Computation 2: $\mathcal{Q}(I_1 \times \mathbb{S}_2^1, \omega = \mathsf{d} x_1 \wedge \mathsf{d} \theta_2; \mathcal{P} = \partial_{\theta_1}).$

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• $H^0(I_1 \times \mathbb{S}^1_2; \mathcal{J}) = 0$ since BS leaves are isolated.

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• Consider
$$I_1 \times \mathbb{S}_2^1 = U \cup V = (I_1 \times (0.4, 1.1)) \cup (I_1 \times (-0.1, 0.6)).$$



$$H^0(V)\oplus H^0(U)\hookrightarrow H^0(W_1)\oplus H^0(W_2)\twoheadrightarrow H^1(I_1 imes \mathbb{S}^1_2).$$

$$H^0(V)\oplus H^0(U) \hookrightarrow H^0(W_1)\oplus H^0(W_2)\twoheadrightarrow H^1(I_1 imes \mathbb{S}^1_2).$$

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 and
 $H^{0}(W_{2}) = \Gamma(I_{1} \times \{0.5\}; \mathbb{C})$. Take $f_{0} \in H^{0}(V)$ and
 $f_{1} \in H^{0}(U) = \Gamma(I_{1} \times \{0\}; \mathbb{C})$. The first map of the sequence is
given by

$$\left(\begin{array}{c} f_2\\ f_3\end{array}\right) = \left(\begin{array}{cc} 1 & -1\\ e^{i\theta \times} & e^{-i\theta \times}\end{array}\right) \left(\begin{array}{c} f_0\\ f_1\end{array}\right)$$

Apply Mayer-Vietoris and computation 1 to obtain

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Thus

$$H^1(I_1 imes \mathbb{S}^1_2) = egin{cases} 0 & ext{if non BS,} \ \mathbb{C} & ext{if there is one BS,} \end{cases}$$

Computation 3: $\mathcal{Q}(I^k \times \mathbb{T}^k; \mathbb{T}^k)$.

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Computation 4:

$$\mathcal{Q}(M^{2n}_{Tor,Reg}; \mathcal{P}(Torus)) = \bigoplus_{j=1}^{n} H^{j}(M; \mathcal{J}) = \mathbb{C}^{b}, \ b = \#BS.$$

Consider an integrable system $(f_1, \ldots, f_n) : M^{2n} \to \mathbb{R}^n$ with, possibly, singular fibers given by non-degenerate critical points.

• Regular:
$$f_j = x_j$$

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• Elliptic:
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- Focus-focus: $\begin{cases} h_j = x_j y_j + x_{j+1} y_{j+1} \\ h_{j+1} = x_j y_{j+1} x_{j+1} y_j \end{cases}$

Consider an integrable system $(f_1, \ldots, f_n) : M^{2n} \to \mathbb{R}^n$ with, possibly, singular fibers given by non-degenerate critical points. We have a local model $B_1 \times \cdots \times B_p$, in some Darboux local coordinates $(x_1, y_1, \ldots, x_n, y_n)$, where each B_l iconforms one of the following canonical models (Miranda):

- Regular: $f_j = x_j$
- Elliptic: $f_j = x_j^2 + y_j^2$
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Geometric quantization of these models gives us all the possible combinations.

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- Focus-Focus fiber: (Miranda, P, Solha) $\begin{cases}
 0 & \text{if the singular fiber is not BS,} \\
 \bigoplus_{1 \le j \le k} C^{\infty}(\mathbb{R}; \mathbb{C}) & \text{if the singular fiber is BS.} \\
 \text{Where, } k = \#\{\text{singular points in the fiber}\} & \text{if the fiber is compact and } k = \#\{\text{singular points in the fiber}\} - 1 & \text{in other case.} \end{cases}$

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- **③** Focus-Focus \times (Whatever) fiber. Apply Künneth and that's it.

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- Focus-Focus × (Whatever) fiber. Apply Künneth and that's it.
 Lego land provides any answer.



Figure: Delzant polytopes of $\mathbb{C}P^2$, $\mathbb{C}P^2 \# 3\overline{\mathbb{C}P}^2$ and $\mathbb{C}P^2 \# 9\overline{\mathbb{C}P}^2$.



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Figure: K3 surface as a singular fiber bundle over the sphere.



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Figure: K3 surface as a singular fiber bundle over the sphere.

Thus, we obtain a K3 with up to 24 BS focus-focus fibers.

$$\mathcal{Q}(\mathcal{K}3)\cong\mathbb{C}^{14}\oplus\bigoplus_{j=1}^{24}C^{\infty}(\mathbb{R};\mathbb{C}).$$

Other examples





Figure: Nodal trade on $\mathbb{S}^2 \times \mathbb{S}^2$.

Other examples





Figure: Nodal trade on $\mathbb{S}^2 \times \mathbb{S}^2$.

 $\mathcal{Q}(\mathbb{S}^2 \times \mathbb{S}^2) \cong C^\infty(\mathbb{R}; \mathbb{C}).$

Other examples

O Spin-spin system.



Figure: Nodal trade on $\mathbb{S}^2 \times \mathbb{S}^2$.

 $\mathcal{Q}(\mathbb{S}^2 \times \mathbb{S}^2) \cong \mathcal{C}^\infty(\mathbb{R};\mathbb{C}).$

2 Spherical pendulum and the spin-oscillator system.



Thanks for listening!

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