# Stochastic Hamilton-Jacobi equations and their long time behaviour

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1 Fully nonlinear SPDE

2 Qualitative results : x-independent case



Qualitative results : x-dependent case

Fully nonlinear SPDE	Qualitative results : x-independent case	Quantitative results in 1d	Qualitative results : x-dependent case
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# Outline

## 1 Fully nonlinear SPDE

2 Qualitative results : x-independent case

3 Quantitative results in 1d

4 Qualitative results : x-dependent case

# Fully nonlinear (2nd order parabolic) Stochastic PDEs

Evolution equations of the type

$$\begin{cases} u(0,\cdot) = u_0, \\ du = F(x, Du, D^2u)dt + H(x, Du) \circ dB_t, \end{cases}$$

where  $u : [0, T] \times \mathbb{R}^N \times \Omega \to \mathbb{R}$ .

Introduced by Lions-Souganidis ('98).

Here B is a random path (typically Brownian motion) in  $\mathbb{R}^d$ , and F is degenerate elliptic i.e. nondecreasing in  $D^2u$ .

## Motivation : movement of interfaces

Interface  $\Gamma$  moving with normal velocity :

$$V = f(x, \vec{n}, D\vec{n}) + \sum_{i} g_i(x, \vec{n}) \dot{B}^i$$

If we represent  $\Gamma$  as a level set  $\{u = 0\}$ , we can write a PDE for u:

$$du = F(x, Du, D^2u)dt + \sum_i H_i(x, Du) \circ dB_t^i$$

with

$$F(x, p, A) = -|p|f\left(x, -\frac{p}{|p|}, -Q_p(A)\right)$$
$$H_i = -|p|g_i\left(x, -\frac{p}{|p|}\right).$$

(Then F elliptic  $\Leftrightarrow$  monotonicity of dynamics)

## Stochastic mean curvature flow

Guiding example : mean curvature flow + stochastic perturbation

$$V = -div_{\Gamma}\vec{n}(x) + \dot{B}(t)$$

Level set PDE

$$\partial_t u = \operatorname{Tr}\left(D^2 u \left(I - \frac{Du \otimes Du}{|Du|^2}\right)\right) + |Du| \circ \dot{B}(t)$$

 $\rightarrow$  quasilinear, parabolic, degenerate PDE + non-linear multiplicative noise

Deterministic dynamics have been studied by many authors. Level set PDE : Evans-Spruck ('91), Chen-Giga-Goto ('91), Barles, Soner, Souganidis,...

The stochastic dynamics can be obtained as limit of bistable reaction-diffusion with additive noise

$$\partial_t v^{\varepsilon} = \Delta v^{\varepsilon} - \frac{1}{\varepsilon} F'(v^{\varepsilon}) + \xi^{\varepsilon}(t)$$

Funaki ('98), Weber ('10), Lions-Souganidis ('00,'20),...

## Solution theory

Pathwise approach :  $\xi = B(\omega)$  is fixed, i.e. consider

 $\partial_t u = F(x, u, Du, D^2u) + H_i(x, u, Du)\dot{\xi}^i(t) \text{ on } (0, T] \times \mathbb{R}^N, \quad u(0, \cdot) = u_0$ 

for fixed  $u_0 \in BUC(\mathbb{R}^N)$ ,  $\xi \in C^0([0, T]; \mathbb{R}^d)$   $(\xi(0) = 0)$ .

- Classical case : ξ ∈ C<sup>1</sup>. No classical solution in general. But Crandall-Lions theory of viscosity solutions give existence/uniqueness.
- Stochastic case : need to extend (continuously) the map  $\xi \mapsto u$  from  $C^1$  to  $C^0$

## Irreversibility

Nonlinearity of the equation  $\rightarrow$  even for smooth  $u_0$ , "shocks" (discontinuities in Du) typically appear in finite time.

Need some PDE theory for weak solutions, the solution is NOT just an SDE solution in infinite dimensions.

For example :

• Finite-dimensional ODE with 1-dimensional noise

$$\dot{X}(t) = V(X)\dot{\xi}(t).$$

The solution is simply given by  $X(t) = \phi_V(X_0, \xi(t))$ .

• For stochastic Hamilton-Jacobi

$$\partial_t u = H(x, Du)\dot{\xi},$$

this is not so simple !

(Indeed : the (deterministic) equation

$$\partial_s U = H(x, DU), \ U(0, \cdot) = u_0$$

typically does not have a  $\mathit{C}^1$  solution on  $(-\infty,\infty)\times\mathbb{R}^{\mathit{N}}.)$ 

## Irreversibility

• Simple example :

$$\partial_t u = |Du| \dot{\xi}(t), \quad u(0,x) = |x|.$$

Then the (sub-)level sets  $\Gamma_R(t) = \{u(t,\cdot) \leqslant R\}$  are given by

$$\Gamma_R(t) = egin{cases} B_0(R-\xi(t)), & t\leqslant au\ \emptyset, & t> au \end{cases}$$

where  $\tau = \inf\{t > 0, R - \xi(t) < 0\}.$ 

## Some well-posedness results

$$\partial_t u = F(x, u, Du, D^2u) + H_i(x, Du)\dot{\xi}^i(t) \text{ on } (0, T] \times \mathbb{R}^N, \quad u(0, \cdot) = u_0$$

- Lions-Souganidis (98-00) : x-independent case. H = H(Du) difference of convex functions (d-dimensional noise), F = F(Du, D<sup>2</sup>u).
- Semilinear case (H = H(x, u, Du) linear in Du) : global transformations, cf. rough path approaches by Friz and coauthors ('10 '13)
- d = 1, H = H(x, Du) strictly convex in Du. Friz-G.-Lions-Souganidis ('17), Seeger ('18), Lions-Souganidis.
- Neumann boundary conditions : G.-Seeger ('21).

Proofs based on maximum principle arguments

$$\frac{d}{dt}\max_{x,y} \left\{ u(t,x) - v(t,y) - \Phi(t,x,y) \right\} \leqslant 0,$$

for suitably chosen test-function  $\Phi$ .

## Rest of the talk

Given a solution u to an equation of the form

$$du = F(x, Du, D^2u)dt + H(x, Du) \circ dB(t), t > 0, x \in \mathbb{T}^N$$

what can we say about  $u(t, \cdot)$  as  $t \to \infty$ ?

No general theory, but we present results in three regimes :

- Convergence to equilibrium in x-independent equations
- Quantitative rate of convergence in dimension 1 ( $x \in \mathbb{T}$ )
- Convergence to equilibrium in x dependent equations (when  $F \equiv 0$ )

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## Basic idea

Consider

$$du = F(D^2u, Du)dt + H(Du) \circ dB_t.$$

Up to translating the solution, assume  $F(0,0) = H(0) = 0 \rightarrow \text{constants}$  are solutions.

- Say we know that solutions to the deterministic equation  $\partial_t v = F(D^2 v, Dv)$  converge to constants.
- *B* Brownian motion : almost surely, *B* will be very small on large intervals.
- By continuity, this means that on these intervals, *u* will behave like *v*, and be close to a constant.

Conclusion :  $\lim_{t\to\infty} u(t, \cdot) = (const.).$ 

(Very classical idea, cf e.g. Dirr-Souganidis '03)

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## Long time behaviour for a stochastic mean curvature flow

Let  $D = D' \times \mathbb{R}$ , where D is a bounded,  $C^2$ , convex domain in  $\mathbb{R}^{N-1}$ . Fix  $\Gamma_0$  closed subset of D, with for some a < b

$$D' imes (-\infty, a] \subset \Gamma_0 \subset D' imes (-\infty, b]$$

and consider  $\Gamma(t)$  its evolution under stochastic mean curvature flow

$$V = -div_{\partial\Gamma(t)}\vec{n}(x) + \dot{B}(t)$$

with  $\Gamma(t) = \Gamma_0$ , and right angle boundary condition

$$n_{\Gamma(t)} \cdot n_D = 0$$
 on  $\partial \Gamma(t) \cap \partial D$ .

#### Theorem (G. Seeger '21+)

Almost surely, for some  $c \in [a, b]$ ,

$$\lim_{t\to\infty} dist_H\left(\Gamma(t),\bar{\Gamma}^c(t)\right)=0,$$

where

$$\bar{\Gamma}^{c}(t) = D \times (-\infty, c + B(t)].$$

## Ideas of proof

We consider the level set formulation

$$\begin{cases} du = \operatorname{Tr}\left(D^2 u \left(I - \frac{D u \otimes D u}{|D u|^2}\right)\right) dt + |D u| d\xi(t) \\ u(0, x) = u_0(x), \qquad \partial_n u = 0 \text{ on } (0, +\infty) \times \partial D \end{cases}$$

Then we use the scheme of proof described above (comparison + related compactness estimates) in the following steps :

- convergence of deterministic dynamics : Giga Ohnuma Sato ('99),  $u \rightarrow v = v(x_n)$ .
- stochastic term  $\rightarrow v \nearrow$  in  $x_n$ .
- convergence in Hausdorff distance requires a Borel-Cantelli argument to rule out "macroscopic holes" for large time.

(Similar results were obtained by different methods in the case of a periodic graph by Es-Sahrir & von Renesse ('12) and Dabrock-Hofmanova-Röger ('21))

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## Stochastic mean curvature in 1D

For illustration purposes, consider the stochastic mean curvature flow evolution of a periodic  $1D\ {\rm graph}$  :

$$\partial_t h = \frac{\partial_{xx} h}{1 + (\partial_x h)^2} + \sqrt{1 + (\partial_x h)^2} \circ \dot{B}(t), \text{ on } [0, \infty) \times \mathbb{T},$$
  
 $h(0, \cdot) = h_0 \in C(\mathbb{T}),$ 

where B is a scalar Brownian motion.

From Section 2 we expect (know) that

$$\operatorname{osc}(u, t \cdot) := \max_{\mathbb{T}} u(t, \cdot) - \min_{\mathbb{T}} u(t, \cdot) \to_{t \to \infty} 0.$$

What about quantitative statements ?

#### Theorem (G. Gess Lions Souganidis ('22+))

There exists c > 0 and a  $\mathbb{P}$ -a.s. finite random variable  $C(\omega)$ , such that, for any h solution to the previous equation, then,

 $\forall t \ge 0, \ \operatorname{osc}(h(t, \cdot)) \leqslant C(\omega) \operatorname{osc}(h_0) e^{-ct}.$ 

Remark : this is an improvement over both the deterministic ( $B \equiv 0$ ) and the purely stochastic case :

• due to the degeneracy as  $\partial_x h \to \infty$ , there exist solutions  $h^R$  to  $\partial_t h = \frac{\partial_\infty h}{1 + (\partial_x h)^2}$  s.t.

$$\operatorname{osc}(h^R(0,\cdot))=R, \quad \operatorname{osc}(h^R(t,\cdot))\geqslant rac{R}{2} ext{ for all } t\leqslant cR.$$

• for  $\partial_t h = \sqrt{1 + (\partial_x h)^2} \circ \dot{B}(t)$ ,  $\operatorname{osc}(h(t, \cdot)) \approx_{t \to \infty} \frac{C}{\max_{[0,t]} B - \min_{[0,t]} B} \approx \frac{C}{t^{1/2}}$ .

• Stochastic structure of B not so important, would also work for B(t) = t,  $C \sin(t)$  or fractional BM.

# Sketch of proof

Rewrite the equation as

$$\partial_t u = \partial_x F(u_x) + H(u_x)\dot{B}(t)$$

(with  $F(r) = \arctan(r)$  and  $H(r) = \sqrt{1 + r^2} - 1$ ). The proof proceeds in 3 steps :

The stochastic term guarantees that

$$\operatorname{osc}(u(\tau_1,\cdot))\leqslant C, \quad ext{for some } au_1\leqslant \ln(\operatorname{osc}(u_0))+C(\omega).$$

Proof : This follows from

$$\frac{d}{dt}\int_{\mathbb{T}} u = \int_{\mathbb{T}} \partial_x(F(u_x)) + \left(\int_{\mathbb{T}} H(u_x)\right) \dot{B}(t)$$

(where  $H(p) = \sqrt{1+p^2} - 1$ ), which leads to

$$\operatorname{osc}(u(t+h)) \leqslant \int_{\mathbb{T}} H(u_{\mathsf{x}}(t+h,\cdot)) \leqslant C rac{\operatorname{osc}(u(t))}{\operatorname{osc}(B)_{t,t+h}}$$

(where  $H(p) = \sqrt{1 + p^2} - 1$ ), and then using that [0, T] can be subdivided in T/C intervals where osc(B) is of order C', this gives  $osc(u(T)) \leq C(\omega)e^{-cT}osc(u_0)$ .

3 At a time 
$$au_2 = au_1 + \mathcal{C}'(\omega)$$
,

$$\|u_x(\tau_2)\|_{L^{\infty}} \leq C.$$

Proof : A similar argument to the previous step gives

$$\|G(u_x(t+h,\cdot))\|_{\infty} \leqslant \frac{\operatorname{osc}(u(t,\cdot))}{\sup_{[u,v]\subset [t,t+h], \ B(u)=\min_{[u,v]}B}\int_u^v (B(r)-B(u))dr},$$

where

$$G(r) = \int_0^r (F'H'')(u) du (\sim_{r\to 0} cr)$$

Indeed :  $N(t) = \int_{\mathbb{T}} H(u_x(\cdot, t)) dx$  satisfies

$$\frac{dN(r)}{dr} \leqslant -\int_{\mathbb{T}} (F'H'')(u_x(x,r))(u_{xx}(x,r))^2 dx \leqslant - \|G(u(\cdot,r)\|_{\infty})$$

and we have

$$\int u(\cdot,t) - \int u(\cdot,s) \ge - \int_s^t (B(u) - B(s)) \frac{dN(u)}{du} du$$

Once the gradient is bounded, we obtain

$$\|u_x(\tau_2+h)\|_{L^2} \leq e^{-ch}\|u_x(\tau_2)\|_{L^2}.$$

Proof : This follows from the uniform ellipticity of the deterministic part :

$$\begin{split} \frac{d}{dt} \int_{\mathbb{T}} u_x^2 &= -\int_{\mathbb{T}} F'(u_x) (u_{xx})^2 - \int_{\mathbb{T}} \underbrace{\left( H''(u_x) u_x u_{xx} \right)}_{\partial_x(\ldots)} \dot{B}(t) \\ &\leqslant - c(\|u_x\|_{\infty}) \int_{\mathbb{T}} u_x^2. \end{split}$$

(Remark : more estimates / examples can be found in the paper).

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In this section, we consider solutions to

$$du = H(x, Du) \circ dB(t)$$
, on  $(0, +\infty) \times \mathbb{T}^d$ 

where H is **convex** in Du.

Corresponds for instance to motion of a (periodic) graph where velocity is given by

$$V(t,z) = a(x)\dot{B}(t)$$

(where  $z = (x, z_n)$  and  $a(\cdot) \ge 0$ ), in which case

$$H(x,p) = a(x)\sqrt{1+|p|^2}.$$

## Main result

### Theorem (G. Gess Lions Souganidis ('22+))

Let H = H(x, Du) be convex in Du and (for instance) strictly convex in Du. Then, there exists  $c_H$  s.t., for any solution u to

$$du + H(x, Du) \circ dB(t) = 0$$
 on  $(0, \infty) \times \mathbb{T}^d$ ,

it holds that

$$u(t,x) = c_H B(t) + \psi(t,x) + o_{t\to\infty}(1).$$

Here  $\psi$  is a statistically stationary solution to

$$\partial_t \psi + (H(x, D\psi) - c_H) \circ dB(t) = 0,$$

satisfying

$$\phi^-(x) \leqslant \psi(t,x) \leqslant \phi^+(x) ext{ for all } t \in \mathbb{R} ext{ and } x \in \mathbb{T}^d.$$

with

$$\pm (c + H(x, D\phi^{\pm})) = 0$$
 in  $\mathbb{T}^d$  (in viscosity sense).

• The proof is based on repeated application of the inequality

 $S_{H}(\delta)S_{-H}(\delta) \leq Id \leq S_{H}(\delta)S_{-H}(\delta)$ 

which crucially relies on convexity of H. (cf. Barron et al. '98, GGLS '20).

The long-time behaviour of u(t, ·) is therefore given by :
(i) the macroscopic position (cB(t))
(ii) the selected limiting upper/profile (φ<sup>+</sup> / φ<sup>-</sup>)
(iii) the actual position of the associated stationary solution ψ. In fact, these are roughly independent in the t → ∞ limit :

$$\left(\frac{B(t)}{\sqrt{t}}, u(t, \cdot) - cB(t)\right) \xrightarrow[t\to\infty]{} \left(Z, \psi^{\Phi^+(B), B'}(0, \cdot)\right).$$

where  $Z \sim \mathcal{N}(0, 1)$ , *B*, *B'* are independent.

## Conclusion : some open questions

- Good understanding of x-dependent equations
- Well-posedness / selection by noise for ill-defined mean curvature flow ?
- Regularity theory ?

$$\partial_t u = \lambda \Delta u + H(Du) \circ \dot{B}(t) ? \Rightarrow ? u(t, \cdot) \in C^1$$

(scaling critical !)