

# On the minimum time control of affine control systems



Michael Orieux  
Under the supervision of  
J.-B. Caillau and J. Féjóz

PhD Defense, 27th of November, 2018

- 1 Introduction and motivation

- ① Introduction and motivation
- ② Structure of the optimal time extremal flow
  - ① For mechanical systems (and more)
  - ② Bounding the switchings for orbit transfers
  - ③ The general case of control-affine systems

- ① Introduction and motivation
- ② Structure of the optimal time extremal flow
  - ① For mechanical systems (and more)
  - ② Bounding the switchings for orbit transfers
  - ③ The general case of control-affine systems
- ③ Optimality of minimum time extremals

- ① Introduction and motivation
- ② Structure of the optimal time extremal flow
  - ① For mechanical systems (and more)
  - ② Bounding the switchings for orbit transfers
  - ③ The general case of control-affine systems
- ③ Optimality of minimum time extremals
- ④ Integrability of the minimum time Kepler problem

# INTRODUCTION AND MOTIVATION

$M$  a smooth manifold,  $U \subset \mathbb{R}^m$ .

Let  $f : M \times \mathbb{R}^m \rightarrow TM$  be a family of smooth vector fields on  $M$ .

$M$  a smooth manifold,  $U \subset \mathbb{R}^m$ .

Let  $f : M \times \mathbb{R}^m \rightarrow TM$  be a family of smooth vector fields on  $M$ .

**An optimal control problem** is given by

$$\begin{cases} \dot{x} = f(x, u), & u(t) \in U \\ x(0) = x_0, \\ x(t_f) = x_f, \\ C(u) = \int_0^{t_f} \varphi(x(t), u(t)) dt \rightarrow \min. \end{cases}$$

Where  $C : L^\infty([0, t_f], U) \rightarrow \mathbb{R}$  is the *cost function*.



$M$  a smooth manifold,  $U \subset \mathbb{R}^m$ .

Let  $f : M \times \mathbb{R}^m \rightarrow TM$  be a family of smooth vector fields on  $M$ .

**An optimal control problem** is given by

$$\begin{cases} \dot{x} = f(x, u), & u(t) \in U \\ x(0) = x_0, \\ x(t_f) = x_f, \\ C(u) = \int_0^{t_f} \varphi(x(t), u(t)) dt \rightarrow \min. \end{cases}$$

Where  $C : L^\infty([0, t_f], U) \rightarrow \mathbb{R}$  is the *cost function*.

$f$  continuous in  $u$ ,  $x_u$ , the solution associated with a control  $u \in L^\infty([0, t_f], U)$  uniquely well defined (Carathéodory).

- Existence of any globally optimal trajectory ?  
→ *Sufficient conditions*: **Filippov's theorem**.
- How can we find - characterize optimal trajectories ?  
→ *Necessary conditions*: **Pontrjagin's Maximum Principle**, extremals.
- Local/global optimality of our extremal trajectories ?  
→ *Second order conditions*: **conjugate** points, **symplectic methods** (Agrachev).
- Regularity of optimal trajectory ?  
→ Techniques from *dynamical systems*: Normal hyperbolicity, invariant manifolds, normal forms.
- Are extremals computable - define an integrable system ?  
→ *Galois differential theory* (but also, symbolic dynamics, Smale's horseshoe...)

# Minimum time affine control systems

Set  $\varphi = 1$ ,  $f(x, u) = F_0(x) + u_1 F_1(x) + u_2 F_2(x)$  and  $U = B$  Euclidean ball.

$$\left\{ \begin{array}{l} \dot{x}(t) = F_0(x(t)) + u_1(t)F_1(x(t)) + u_2(t)F_2(x(t)), \quad u_1^2 + u_2^2 \leq 1 \\ x(0) = x_0 \\ x(t_f) = x_f \\ t_f \rightarrow \min. \end{array} \right. \quad (1)$$

$F_i$  smooth,  $i = 0, 1, 2$ ,  $x_0, x_f \in M$  a 4 dimensional manifold (can be generalized to  $2n$  with  $n$  controls).

# Minimum time affine control systems

Set  $\varphi = 1$ ,  $f(x, u) = F_0(x) + u_1 F_1(x) + u_2 F_2(x)$  and  $U = B$  Euclidean ball.

$$\begin{cases} \dot{x}(t) = F_0(x(t)) + u_1(t)F_1(x(t)) + u_2(t)F_2(x(t)), & u_1^2 + u_2^2 \leq 1 \\ x(0) = x_0 \\ x(t_f) = x_f \\ t_f \rightarrow \min. \end{cases} \quad (1)$$

$F_i$  smooth,  $i = 0, 1, 2$ ,  $x_0, x_f \in M$  a 4 dimensional manifold (can be generalized to  $2n$  with  $n$  controls).

Structure of **Lie algebra**  $\mathbf{Lie}(F_0, F_1, F_2)$  is crucial.

**Example:** *Mechanical systems.*

$$\ddot{q} + \nabla V(q) = u,$$

$V$  a smooth potential.

## Motivating example: the controlled CR3BP

$$\ddot{q} + \nabla V_\mu(q) - 2i\dot{q} = u, \quad \|u\| \leq 1 \quad (2)$$

in the rotating frame,  $u$  being the control (thrust of the engine) and

$$V_\mu(q) = \frac{1}{2}|q|^2 + \frac{1-\mu}{|q+\mu|} + \frac{\mu}{|q-1+\mu|}, \quad \mu = \text{mass ratio.}$$

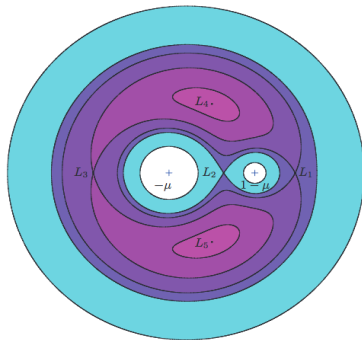


Figure: Hill's region and Lagrange points for the RC3BP

Minimizing the cost generates *singularities*.

This thesis focuses on understanding of the consequences of these singularities:

- On the extremal flow, irregular behavior, non-uniqueness,
- On the optimality of the trajectories,
- On the computability of the possible optimal trajectories: destruction of first integrals.

# STRUCTURE OF THE MINIMUM TIME EXTREMAL FLOW

## Definition (Pseudo-Hamiltonian)

$$\forall (x, p) \in T^*M, H(x, p, u) = \langle p, f(x, u) \rangle$$

Control affine dynamics:  $H(x, p, u) = H_0(x, p) + u_1 H_1(x, p) + u_2 H_2(x, p)$ ,  
 $H_i(x, p) = \langle p, F_i(x) \rangle$ ,  $i = 0, 1, 2$ .



## Definition (Pseudo-Hamiltonian)

$$\forall (x, p) \in T^*M, H(x, p, u) = \langle p, f(x, u) \rangle$$

Control affine dynamics:  $H(x, p, u) = H_0(x, p) + u_1 H_1(x, p) + u_2 H_2(x, p)$ ,  
 $H_i(x, p) = \langle p, F_i(x) \rangle$ ,  $i = 0, 1, 2$ .

## Theorem (Pontrjagin)

*If  $(x, u)$  is a minimum time trajectory then there exists an absolutely continuous Lipschitz curve  $p(t) \in T_{x(t)}M^* \setminus \{0\}$  s.t.*

-  $(x, p)$  is solution of :

$$\begin{cases} \dot{x} = \frac{\partial H}{\partial p}(x, p, u) \\ \dot{p} = -\frac{\partial H}{\partial x}(x, p, u). \end{cases} \quad (3)$$

-  $H(x(t), p(t), u(t)) = \max_{\tilde{u} \in U} H(x(t), p(t), \tilde{u})$ .

-  $H(x(t), p(t), u(t)) \geq 0$ .

$$\begin{cases} \dot{x}(t) = F_0(x(t)) + u_1(t)F_1(x(t)) + u_2(t)F_2(x(t)), & u_1^2 + u_2^2 \leq 1 \\ x(0) = x_0 \\ x(t_f) = x_f \\ t_f \rightarrow \min. \end{cases} \quad (4)$$

## Proposition

For system (4)  $\forall z = (x, p) \in T^*M$ ,  $H^{\max}(x, p) = H_0(x, p) + \sqrt{H_1^2(x, p) + H_2^2(x, p)}$

$u = \frac{1}{\sqrt{H_1^2 + H_2^2}}(H_1, H_2)$  : discontinuities of the control  $u$  are called **switchings**.

$$\begin{cases} \dot{x}(t) = F_0(x(t)) + u_1(t)F_1(x(t)) + u_2(t)F_2(x(t)), & u_1^2 + u_2^2 \leq 1 \\ x(0) = x_0 \\ x(t_f) = x_f \\ t_f \rightarrow \min. \end{cases} \quad (4)$$

## Proposition

For system (4)  $\forall z = (x, p) \in T^*M$ ,  $H^{\max}(x, p) = H_0(x, p) + \sqrt{H_1^2(x, p) + H_2^2(x, p)}$

$u = \frac{1}{\sqrt{H_1^2 + H_2^2}}(H_1, H_2)$  : discontinuities of the control  $u$  are called **switchings**.

## Definition (Singular locus / Switching surface.)

$\Sigma = \{z = (x, p) \in T^*M, H_1(x, p) = H_2(x, p) = 0\} = F_1^\perp \cap F_2^\perp$  smooth submanifold if  $F_1$  and  $F_2$  are linearly independent.

Set  $\bar{z} \in \Sigma$ ,  $\bar{x} = \pi(\bar{z})$ , where  $\pi : T^*M \rightarrow M$  the canonical projection.

Denote  $F_{ij} = [F_i, F_j]$ ,  $H_{ij} = \{H_i, H_j\}$ .

Assume :

$$(\mathcal{A}) : \det(F_1(\bar{x}), F_2(\bar{x}), F_{01}(\bar{x}), F_{02}(\bar{x})) \neq 0$$

Checked for mechanical systems.

→  $(\mathcal{A}) + F_0$  recurrent ( $\mu = 0$  or certain Hill's regions of the RC3BP)  $\Rightarrow$  Controllability.

Set  $\bar{z} \in \Sigma$ ,  $\bar{x} = \pi(\bar{z})$ , where  $\pi : T^*M \rightarrow M$  the canonical projection.

Denote  $F_{ij} = [F_i, F_j]$ ,  $H_{ij} = \{H_i, H_j\}$ .

Assume :

$$(\mathcal{A}) : \det(F_1(\bar{x}), F_2(\bar{x}), F_{01}(\bar{x}), F_{02}(\bar{x})) \neq 0$$

Checked for mechanical systems.

$\rightarrow (\mathcal{A}) + F_0$  recurrent ( $\mu = 0$  or certain Hill's regions of the RC3BP)  $\Rightarrow$  Controllability.

## Proposition

*Any system of the form  $\ddot{q} + g(q, \dot{q}) = u$  verifies  $(\mathcal{A})$ .*

In mechanical systems we also have  $(\mathcal{B}) : F_{12} = 0 \Rightarrow H_{12} = 0$ .

Using  $(\mathcal{A})$ :  $(x, p) \mapsto (x, H_1, H_2, H_{01}, H_{02})$  is a change of coordinates.

→ **Polar blow up**:  $(H_1, H_2) = (\rho \cos \theta, \rho \sin \theta)$ ,

Using  $(\mathcal{A})$ :  $(x, p) \mapsto (x, H_1, H_2, H_{01}, H_{02})$  is a change of coordinates.

→ **Polar blow up**:  $(H_1, H_2) = (\rho \cos \theta, \rho \sin \theta)$ ,

$$(\Upsilon) : \begin{cases} \dot{\rho} = -\sin \theta \\ \dot{\theta} = \frac{1}{\rho}(1 + \alpha - \cos \theta + f(\rho, \theta, \xi)) \\ \dot{\xi} = h(\rho, \theta, \xi), \end{cases} \quad (5)$$

$\xi = (x, H_{01}, H_{02})$ .

- (i)  $f, h$  smooth functions,  $h$  has values in  $\mathbb{R}^k$ ;
- (ii)  $f(\bar{z}) = 0$ ,  $f$  has a nice local behavior around  $\bar{z}$ .

Using  $(\mathcal{A})$ :  $(x, p) \mapsto (x, H_1, H_2, H_{01}, H_{02})$  is a change of coordinates.

→ **Polar blow up**:  $(H_1, H_2) = (\rho \cos \theta, \rho \sin \theta)$ ,

$$(\Upsilon) : \begin{cases} \dot{\rho} = -\sin \theta \\ \dot{\theta} = \frac{1}{\rho}(1 + \alpha - \cos \theta + f(\rho, \theta, \xi)) \\ \dot{\xi} = h(\rho, \theta, \xi), \end{cases} \quad (5)$$

$\xi = (x, H_{01}, H_{02})$ .

(i)  $f, h$  smooth functions,  $h$  has values in  $\mathbb{R}^k$ ;

(ii)  $f(\bar{z}) = 0$ ,  $f$  has a nice local behavior around  $\bar{z}$ .

## Definition (Partition of the singular locus)

$\Sigma = \Sigma_0 \cup \Sigma_- \cup \Sigma_+$  with :

$\Sigma_- = \{\alpha < 0\}$ ,  $\Sigma_+ = \{\alpha > 0\}$ ,  $\Sigma_0 = \{\alpha = 0\}$ .



# THE CASE $\Sigma_-$ : MECHANICAL SYSTEMS

## Theorem (Caillau, Fejoz, O.)

*There exists unique solution for system (1) in a neighborhood  $O_{\bar{z}}$  of  $\bar{z}$ , and there is at most one switch on  $O_{\bar{z}}$ .*

- *If  $\bar{z} \in \Sigma_-$ : The local extremal flow  $z : (t, z_0) \in [0, t_f] \times O_{\bar{z}} \mapsto z(t, z_0) \in M$  is piecewise smooth, and smooth on each strata :*

$$O_{\bar{z}} = S_0 \sqcup S^s \cup S^u \cup \Sigma$$

- *where  $S^s$  (resp.  $S^u$ ) is the codimension one submanifold of initial conditions leading to the switching surface (resp. in negative times),*
- *$S_0 = O_{\bar{z}} \setminus (S^s \cup S^u \cup \Sigma)$ .*

Constraint on the control  $\rightarrow$  No singular flow inside  $\Sigma_-$ . Extend previous results from Agrachev and Biolo (2016) in a close context.

Regularize (Y) by rescaling the time  $dt = \rho d\tau$ .

$$(Z) : \begin{cases} \rho' = -\rho \sin \theta \\ \theta' = 1 + \alpha - \cos \theta + O(\rho; \xi) \\ \xi' = \rho h(\rho, \theta, \xi) \end{cases} \quad (6)$$

Regularize (Y) by rescaling the time  $dt = \rho d\tau$ .

$$(Z) : \begin{cases} \rho' = -\rho \sin \theta \\ \theta' = 1 + \alpha - \cos \theta + O(\rho; \xi) \\ \xi' = \rho h(\rho, \theta, \xi) \end{cases} \quad (6)$$

Then, in  $\{\rho = 0\}$ , for each  $\xi$ , two parabolic equilibria  $\theta_{\pm}$ .

The manifolds  $N_{\pm} = \{(0, \theta_{\pm}, \bar{\xi})\}$  are **normally hyperbolic** invariant submanifolds of *equilibria*, with stable and unstable manifolds of **dimension one**.

→ Finite initial time  $t_f$ : *existence* and *uniqueness*.

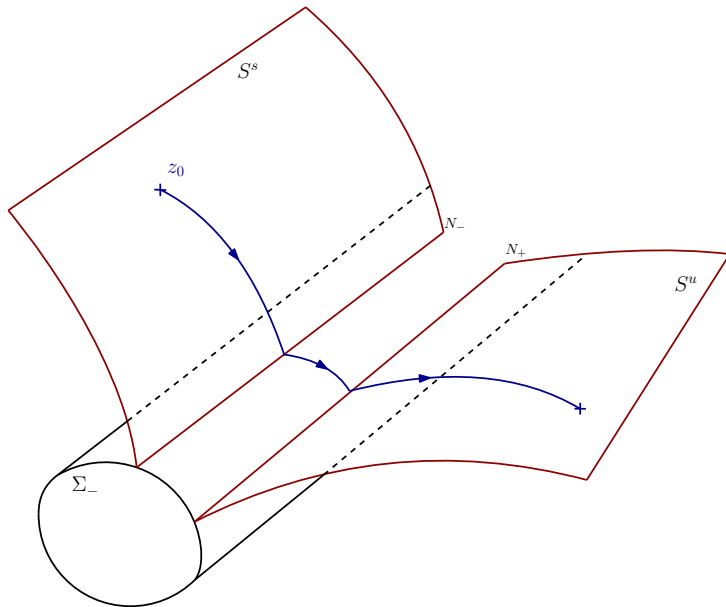


Figure: Blown-up phase portrait around  $\Sigma_-$

Set

$$S^s = \cup_{z \in N_-} W^s(z)$$

and

$$S^u = \cup_{z \in N_+} W^u(z).$$

Eigenvalues on  $N_{\pm}$  are null, and in the direction of  $\rho$ , the Jacobian has a strictly negative eigenvalue: the spectral gap is infinite  $\rightarrow S^s$  ( $S^u$ ) is a  $C^{\infty}$ -smooth foliation (Hirsch, Pugh, Shub).

The flow is smooth restricted to  $S^s$  and  $S_0$ : the regularized vector field is smooth, consequence of the dominated convergence theorem.

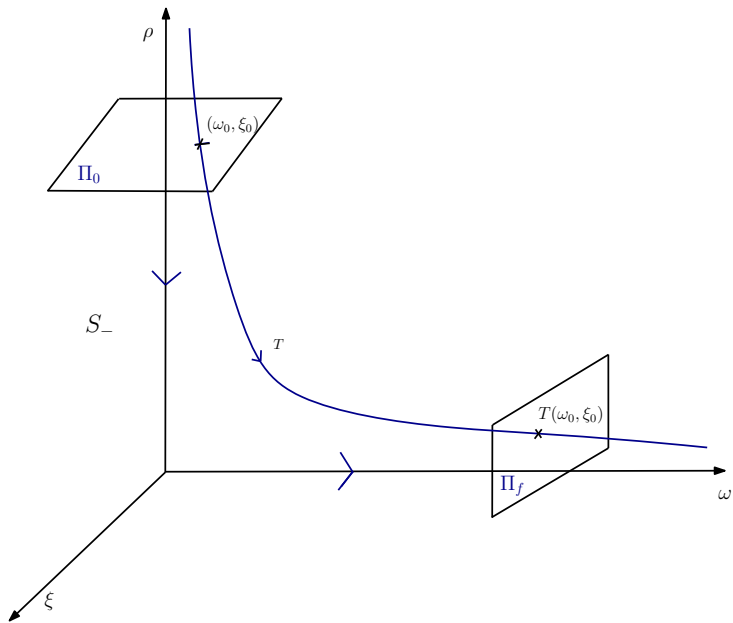


Figure: The regular-singular transition

$\Pi_0 \subset \{\rho = \rho_0\}$ ,  $\Pi_f \subset \{\omega = \omega_f\}$ .

Theorem (Caillau, Fejoz, O., Roussarie)

*Let  $T : \Pi_0 \rightarrow \Pi_f$  be the Poincaré mapping between the two sections.*

*Then,  $T$  is smooth in  $(\omega_0 \ln \omega_0, \omega_0, \xi_0)$ , that*

$$T(\omega_0, \xi_0) = (R(\omega_0 \ln \omega_0, \omega_0, \xi_0), X(\omega_0 \ln \omega_0, \omega_0, \xi_0)).$$

*$R, X$  smooth.*

Idea : Straightening via a normal form + Blow up



# BOUNDING THE NUMBER OF SWITCHINGS FOR ORBIT TRANSFERS

Previous theorems apply to the CR3BP but not to the **non-autonomous** Restricted Elliptic Three Body Problem:

$$\ddot{q} + \nabla V_{\mu}(t, q) = u,$$

with  $V_{\mu}(t, q) = \frac{1-\mu}{\|q-q^1(t)\|} + \frac{\mu}{\|q-q^2(t)\|}$ .  $q^1, q^2$  position vectors of the two primaries.

Previous theorems apply to the CR3BP but not to the **non-autonomous** Restricted Elliptic Three Body Problem:

$$\ddot{q} + \nabla V_\mu(t, q) = u,$$

with  $V_\mu(t, q) = \frac{1-\mu}{\|q-q^1(t)\|} + \frac{\mu}{\|q-q^2(t)\|}$ .  $q^1, q^2$  position vectors of the two primaries.

## Definition

*Define  $\delta = \inf_{[0, t_f]} |q(t)|$ ,  $\delta_1 = \inf_{[0, t_f]} |q(t) - q^1(t)|$ ,  $\delta_2 = \inf_{[0, t_f]} |q(t) - q^2(t)|$ . This quantities represents the distance to the collisions in the two body, and restricted three-body problems respectively. Finally note  $\delta_{12}(\mu) = \frac{\delta_1 \delta_2}{((1-\mu)\delta_2^3 + \mu\delta_1^3)^{1/3}}$ .*

Upper bound on the number of switchings  $\rightarrow$  bounds the number of heteroclinic connections between  $S^s$  and  $S^u$ .

## Theorem (Morse)

Let  $a, b \in \mathbb{R}$ , with  $b > a$ . Consider the two linear second order equations

$$z'' + P(t)z = 0, \quad (7)$$

$$z'' + Q(t)z = 0, \quad (8)$$

with  $P(t), Q(t) \in S_n(\mathbb{R})$  s. t.  $Q(t) - P(t) \geq 0$ , there exists  $\bar{t}$  with  $Q(\bar{t}) - P(\bar{t}) > 0$ . If (7) has a non trivial solution  $y$ ,  $y(a) = y(b) = 0$ , then (8) has a non trivial solution which vanishes in  $a$  and  $c < b$ .

## Theorem (Morse)

Let  $a, b \in \mathbb{R}$ , with  $b > a$ . Consider the two linear second order equations

$$z'' + P(t)z = 0, \quad (7)$$

$$z'' + Q(t)z = 0, \quad (8)$$

with  $P(t), Q(t) \in S_n(\mathbb{R})$  s. t.  $Q(t) - P(t) \geq 0$ , there exists  $\bar{t}$  with  $Q(\bar{t}) - P(\bar{t}) > 0$ . If (7) has a non trivial solution  $y$ ,  $y(a) = y(b) = 0$ , then (8) has a non trivial solution which vanishes in  $a$  and  $c < b$ .

## Proposition (Caillau, Fejoz, O.)

- Keplerian case, the maximum amount of switchings is  $N_0 = \left\lceil \frac{t_f}{\pi\delta^{3/2}} \right\rceil$  on  $[0, t_f]$ .
- Controlled Elliptic Three-Body Problem with a mass ratio  $\mu$ , On a time interval  $[0, t_f]$  the maximum amount of such singularities is  $N_\mu = \left\lceil \frac{t_f}{\pi\delta_{12}(\mu)^{3/2}} \right\rceil$ .

The controlled ER3BP dynamics:

$$H(q, v, p_q, p_v) = p_q \cdot v - p_v \cdot \nabla V_\mu(t, q) + p_v \cdot u,$$

and

$$H^{\max}(q, v, p_q, p_v) = p_q \cdot v - p_v \cdot \nabla V_\mu(t, q) + \|p_v\|.$$

Linear equation in  $p_v$ :

$$\ddot{p}_v + \nabla_q^2 V_\mu(t, q) p_v = 0. \quad (9)$$

Compare  $V_\mu(t, q)$  with a well-chosen matrix.

$$A_t(q) = \begin{pmatrix} 1 + \frac{1-\mu}{|q-q^1(t)|^3} + \frac{\mu}{|q-q^2(t)|^3} & 0 \\ 0 & \frac{1-\mu}{|q-q^1(t)|^3} + \frac{\mu}{|q-q^2(t)|^3} \end{pmatrix} \in \mathcal{S}_2(\mathbb{R}).$$

# THE GENERAL AFFINE CASE

$$(\alpha = 0) \Rightarrow \begin{cases} \dot{\rho} = -\rho \sin \theta \\ \dot{\theta} = 1 - \cos \theta + O(\rho; \xi) \\ \dot{\xi} = \rho \tilde{h}(\rho, \theta, \xi) \end{cases} \quad (10)$$

$\alpha_0$  smooth on  $\mathbb{R}^k$ . Set  $\tilde{h}_1(0, 0, 0) = c$ .

## Theorem (O., Roussarie)

*Let  $\bar{z}$  be in  $\Sigma_0$ . If  $c > 0$ , there exist extremals passing through  $\bar{z}$ , these extremals are connected to the singular flow in  $\Sigma_0$ .*

Idea: Dimensional reduction + quasi-homogeneous blow up + dynamical study in the plane.



Under generic hypothesis,  $\exists$  coordinates  $\tilde{\xi} = (\zeta, \tilde{\xi}_2, \dots, \tilde{\xi}_k)$  s.t.

$$\begin{cases} \rho' = -\rho\theta + O(\rho\theta^3) \\ \theta' = \zeta + \theta^2/2 + O(\rho + |\theta|^4) \\ \zeta' = c\rho + \rho O(\rho + |\theta| + |\tilde{\xi}|). \end{cases} \quad (11)$$

Under generic hypothesis,  $\exists$  coordinates  $\tilde{\xi} = (\zeta, \tilde{\xi}_2, \dots, \tilde{\xi}_k)$  s.t.

$$\begin{cases} \rho' = -\rho\theta + O(\rho\theta^3) \\ \theta' = \zeta + \theta^2/2 + O(\rho + |\theta|^4) \\ \zeta' = c\rho + \rho O(\rho + |\theta| + |\tilde{\xi}|). \end{cases} \quad (11)$$

**Blow up.** **Nilpotent equilibrium**  $(\rho, s, \zeta) = (0, 0, 0)$ , we will use a specific blow-up:

$$\begin{cases} \rho = R^3 \bar{\rho} \\ \theta = R\bar{\theta} \\ \zeta = R^2 \bar{\zeta} \end{cases}$$

with  $(\bar{\rho}, \bar{s}, \bar{\zeta}) \in S_+^2$  the *hemisphere*  $\rho \geq 0$ ,  $R \in \mathbb{R}_+$ .

**The chart (i).** Interior of  $S_+^2$ : One hyperbolic equilibrium  $m_0$ , unstable node restricted to  $S_+^2$ .

**The chart (ii).** 4 semi-hyperbolic equilibria  $\in \partial S_+^2 = \{R = 0, \bar{\rho} = 0\} \cong \mathbb{S}^1$ .

–  $-\pi/2$ : **unstable** node in restriction to  $S_+^2$

–  $\pi/2$ : **stable** node in restriction to  $S_+^2$

–  $\omega_0 \in ]\pi/2, \pi[$  (**unstable**) and  $-\omega_0$  (**stable**) nodes restricted to  $\mathbb{S}^1$ .

$\pm\omega_0$ : hyperbolic restricted to  $S_+^2$ , with lines of zeros in the plane  $\bar{\rho} = 0$ .

The stable and unstable manifolds of these equilibria can be connected using

## Theorem (Poincaré-Bendixson)

*Let  $X$  be a vector field in the plane, any maximal solution of  $\dot{x} = X(x)$  contained in a compact set, is either converging to an equilibrium point or a limit cycle.*

No periodic orbit: choosing a transverse domain containing  $m_0$  and using Poincaré-Hopf formula.

→ stable and unstable manifolds can be connected in a **unique way**: the stable manifold from  $\omega_0$  is connected to  $m_0$ .

→ Finite initial time.

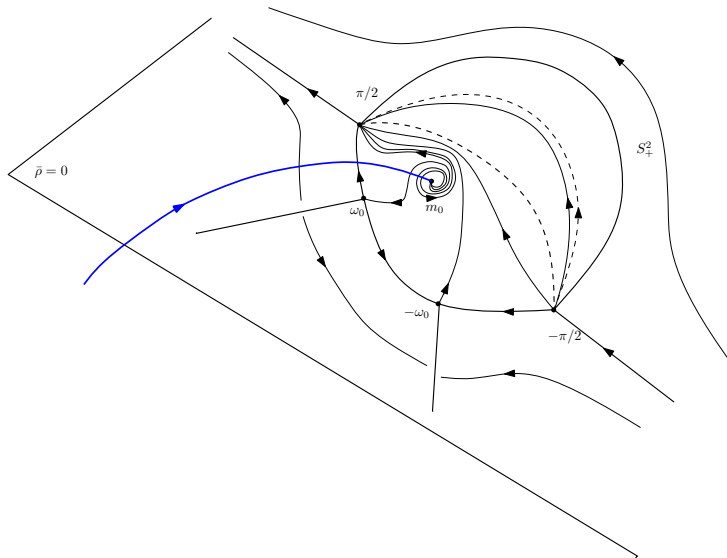


Figure: Phase portrait around the nilpotent equilibrium

## Theorem (O., Roussarie)

*In a neighborhood  $O_{\bar{z}}$  of a point  $\bar{z} \in \Sigma_0$ , the flow is well defined, continuous, and piecewise smooth. More precisely, there exists a stratification:*

$$O_{\bar{z}} = S_0 \cup S^s \cup S^u \cup S_0^s$$

*where*

- $S_0^s$  is the submanifold of codimension 2 of initial conditions leading to  $\Sigma_0$ ,*
- $S^s$  is the submanifold of codimension 1 of initial conditions leading to  $\Sigma_-$ ,*
- $S_0 = O_{\bar{z}} \setminus (S_1^0 \cup S_1^1)$ .*

*The extremal flow is smooth on each stratum.*

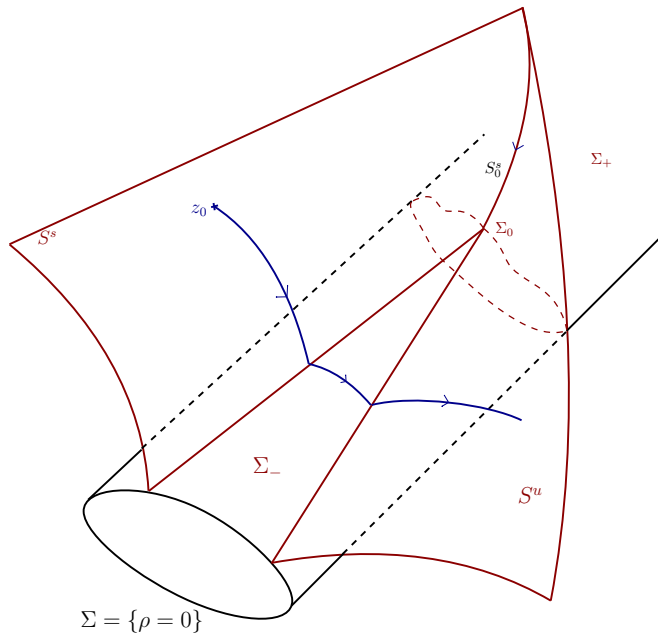


Figure: Stratification of the extremal flow around  $\Sigma_0$

## The case $\Sigma_+$ .

### Proposition

*No switching at  $\Sigma_+$ . In a neighborhood of a point  $\bar{z}$  in  $\Sigma_+$ , there is no switch, and the extremal flow is smooth, i.e.,  $\Sigma_+$  is never crossed. In other words,  $\rho$  does not vanish in (5).*

### The singular flow (flow inside $\Sigma$ ):

- There exists a smooth singular flow inside  $\Sigma_+$ , which cannot be optimal (Goh condition).
- There cannot exist any admissible singular extremal in  $\Sigma_-$ .
- There exists a singular flow in  $\Sigma_0$ .



# Jump on the control

$$r := \sqrt{H_{01}^2 + H_{02}^2}$$

## Remark

*In the case  $\bar{z} \in \Sigma_-$ , the jump on the control at a switching time  $\bar{t}$  is:*

$$u(\bar{t}_\pm) = (\cos \theta_\pm, \sin \theta_\pm) = \frac{1}{r^2} (-H_{02}H_{12} \pm H_{01}\sqrt{r^2 - H_{12}^2}, H_{01}H_{12} \pm H_{02}\sqrt{r^2 - H_{12}^2}).$$

**(B)**  $\Rightarrow$   $\pi$ -singularities for mechanical systems.

# Jump on the control

$$r := \sqrt{H_{01}^2 + H_{02}^2}$$

## Remark

*In the case  $\bar{z} \in \Sigma_-$ , the jump on the control at a switching time  $\bar{t}$  is:*

$$u(\bar{t}_{\pm}) = (\cos \theta_{\pm}, \sin \theta_{\pm}) = \frac{1}{r^2} (-H_{02}H_{12} \pm H_{01}\sqrt{r^2 - H_{12}^2}, H_{01}H_{12} \pm H_{02}\sqrt{r^2 - H_{12}^2}).$$

**(B)**  $\Rightarrow$   $\pi$ -singularities for mechanical systems.

## Proposition

*In  $\Sigma_0$ . Consider the extremal  $z(t)$  entering the singular locus in  $z(\bar{t}) = \bar{z} \in \Sigma_0$ ,*

- If  $H_{12}(\bar{z}) = r(\bar{z})$ , the extremal control is continuous on  $[0, t_f]$ ,*
- If  $H_{12} = -r(\bar{z})$ , the extremal control has a  $\pi$ -singularity at the switching time  $\bar{t}$ .*

# OPTIMALITY OF MINIMUM TIME EXTREMALS

## Definition (exponential map)

We call exponential mapping from  $x_0$ , the map

$$\exp_{x_0} : (t, p_0) \in [0, t_f] \times T_{\bar{x}_0}^* M \cap S^S \rightarrow \pi(z(t, x_0, p_0)) = x(t, x_0, p_0) \in M$$

Assumption:  $T_{x_0}^* M \pitchfork S^S$ , then  $T_{x_0}^* M \cap S^S$  is a smooth submanifold of dimension 3.

## Theorem (O.)

Denote  $M(t) := d \exp_{\bar{x}_0}(t, \bar{p}_0)$ . If

- (i) The reference extremal is normal,
  - (ii)  $\det M(t) \neq 0$  for all  $t \in ]0, \bar{t}[ \cup ]\bar{t}, \bar{t}_f]$  and  $\det M(\bar{t}_-) \det M(\bar{t}_+) \neq 0$ ,
- then the reference trajectory is a  $\mathcal{C}^0$ -local minimizer among all trajectories with same endpoints.

$\exists$  Lagrangian submanifold  $\mathcal{L}$  transverse to  $T_{x_0}^*M$ , s. t.  $\mathcal{S}_0 = \mathcal{L} \cap S^s$  is a smooth submanifold of dimension 3  $\rightarrow$  *Regularity* on  $\mathcal{S}_0$ .

$\exists$  Lagrangian submanifold  $\mathcal{L}$  transverse to  $T_{x_0}^*M$ , s. t.  $\mathcal{S}_0 = \mathcal{L} \cap S^s$  is a smooth submanifold of dimension 3  $\rightarrow$  *Regularity* on  $\mathcal{S}_0$ .

The canonical projection  $\pi$  is a **homeomorphism** on

$$\mathcal{S}_1 = \{z(t, z_0), (t, z_0) \in [0, \bar{t}(z_0)] \times \mathcal{S}_0\} \quad (12)$$

onto its image. The same holds for

$$\mathcal{S}_2 = \{z(t, z_0), (t, z_0) \in [\bar{t}(z_0), t_f] \times \mathcal{S}_0\} \quad (13)$$

and  $\mathcal{S}_1 \cup \mathcal{S}_2$  (extremals cut  $\Sigma$  transversally).

$\exists$  Lagrangian submanifold  $\mathcal{L}$  transverse to  $T_{x_0}^*M$ , s. t.  $\mathcal{S}_0 = \mathcal{L} \cap S^s$  is a smooth submanifold of dimension 3  $\rightarrow$  *Regularity* on  $\mathcal{S}_0$ .

The canonical projection  $\pi$  is a **homeomorphism** on

$$\mathcal{S}_1 = \{z(t, z_0), (t, z_0) \in [0, \bar{t}(z_0)] \times \mathcal{S}_0\} \quad (12)$$

onto its image. The same holds for

$$\mathcal{S}_2 = \{z(t, z_0), (t, z_0) \in [\bar{t}(z_0), t_f] \times \mathcal{S}_0\} \quad (13)$$

and  $\mathcal{S}_1 \cup \mathcal{S}_2$  (extremals cut  $\Sigma$  transversally).

$\rightarrow$  Cost comparison using the Liouville form  $\lambda = p dx$  exact on  $\mathcal{S}_i$ : Extremals are **locally optimal**.

INTEGRABILITY OF MINIMUM TIME  
HAMILTONIAN IN THE KEPLER PROBLEM



## Definition (Liouville integrability)

*Let  $H$  be a smooth function on a  $2n$ -dimensional symplectic manifold. The associated Hamiltonian system is integrable iff there exists  $n$  independent first integrals (constant of motion) in involution.*

Classical reduction of the two body problem  $\mu = 0$ ,

$$\ddot{q} + \frac{q}{\|q\|^3} = \varepsilon u.$$

Uncontrolled two body problem ( $\varepsilon = 0$ ) is well known to be integrable.

## Definition (Liouville integrability)

*Let  $H$  be a smooth function on a  $2n$ -dimensional symplectic manifold. The associated Hamiltonian system is integrable iff there exists  $n$  independent first integrals (constant of motion) in involution.*

Classical reduction of the two body problem  $\mu = 0$ ,

$$\ddot{q} + \frac{q}{\|q\|^3} = \varepsilon u.$$

Uncontrolled two body problem ( $\varepsilon = 0$ ) is well known to be integrable.

Kepler problem with a constant force :  $u = \text{cst}$ , also integrable (Charlier and Saint Germain).

## Definition (Liouville integrability)

*Let  $H$  be a smooth function on a  $2n$ -dimensional symplectic manifold. The associated Hamiltonian system is integrable iff there exists  $n$  independent first integrals (constant of motion) in involution.*

Classical reduction of the two body problem  $\mu = 0$ ,

$$\ddot{q} + \frac{q}{\|q\|^3} = \varepsilon u.$$

Uncontrolled two body problem ( $\varepsilon = 0$ ) is well known to be integrable.

Kepler problem with a constant force :  $u = \text{cst}$ , also integrable (Charlier and Saint Germain).

Three body problem is **not integrable** (Poincaré).

$$\begin{cases} \ddot{\mathbf{q}} + \frac{\mathbf{q}}{\|\mathbf{q}\|^3} = \mathbf{u}, \|\mathbf{u}\| \leq 1, \\ (\mathbf{q}(0), \mathbf{v}(0)) = (\mathbf{q}_0, \mathbf{v}_0), \\ (\mathbf{q}(t_f), \mathbf{v}(t_f)) = (\mathbf{q}_f, \mathbf{v}_f) \\ t_f \rightarrow \min. \end{cases} \quad (14)$$

$$\begin{cases} \ddot{q} + \frac{q}{\|q\|^3} = u, \|u\| \leq 1, \\ (q(0), v(0)) = (q_0, v_0), \\ (q(t_f), v(t_f)) = (q_f, v_f) \\ t_f \rightarrow \min. \end{cases} \quad (14)$$

Maximized Hamiltonian:

$$H(q, v, p_q, p_v) = p_q \cdot v - \frac{p_v \cdot q}{\|q\|^3} + \|p_v\|. \quad (15)$$

→ **Liouville integrability** of  $H$  ?

## Theorem (Morales-Ramis)

*Let us consider a Hamiltonian  $\mathcal{H}$  analytic on a complex analytic symplectic manifold and a particular solution  $\Gamma$  not reduced to a point. If  $\mathcal{H}$  is integrable in the Liouville sense with meromorphic first integrals, then the first order variational equation near  $\Gamma$  has a virtually Abelian Galois group over the base field of meromorphic functions on  $\Gamma$ .*

## Theorem (Morales-Ramis)

*Let us consider a Hamiltonian  $\mathcal{H}$  analytic on a complex analytic symplectic manifold and a particular solution  $\Gamma$  not reduced to a point. If  $\mathcal{H}$  is integrable in the Liouville sense with meromorphic first integrals, then the first order variational equation near  $\Gamma$  has a virtually Abelian Galois group over the base field of meromorphic functions on  $\Gamma$ .*

## Theorem (Caillau, Combot, Fejoz, O.)

*The minimum time Kepler problem is not meromorphically Liouville integrable on  $\mathcal{M}$ .*

Invariant manifold of collisions

$$S = \{q_2 = v_2 = p_{q_2} = p_{v_2} = 0\} \cap \mathcal{M}$$

- $H$  is integrable on  $S$ : collision trajectory  $\Gamma(t)$ .
- Compute the Normal Variational Equation along  $\Gamma(t)$ .
- Its Galois group contains the group of a hypergeometric equation: Contains  $SL_2(\mathbb{C})$ , not even solvable.
- The variational equation is Fuchsian: Non-integrability in the class of meromorphic functions (Schlesinger's density theorem).



# PERSPECTIVES

- Find a generic condition to treat the general case.
- Numerical experimentation for the extremal.
- Real non integrability for the minimum time Kepler problem.
- KAM theory: non geodesic convexity in the Kepler configuration?

Thank you for  
your attention !